

Simplified Dark Matter models with a long-lived co-annihilation partner

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based on arXiv:1702.00750

in collaboration with
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Introduction

- All evidence for the existence of Dark Matter is purely gravitational
- The particle physics nature of DM is unknown
— is our main ‘new physics’ challenge
[pretty much main particle physics challenge at present]
- The current measured relic density of DM from cosmic microwave background $\Omega_{\text{DM}} h^2 = 0.1198 \pm 0.0026$

DM searches:

- Direct Detection (DD) experiments
- Indirect Detection (ID) experiments
- Collider searches LHC (& future colliders 13-100 TeV)

Motivation

- A standard signature to search for dark matter at colliders is the mono-X (or multi-jets) plus missing energy.
- These searches are being exploited and interpreted in terms of **simplified dark matter models with mediators**.

Dark Matter + mediator + Standard Model particles

- [A growing number of the analyses are also dedicated to the direct search of the mediators which can decay back to the SM.]
- **We consider instead an alternative DM scenario** characterised by **simplified models without mediators**.

Dark Matter + co-annihilation partner + Standard Model particles

- Our dark sector includes a co-annihilation partner (CAP) particle instead of a mediator (in addition to the cosmologically stable DM).

Simplified Models without Mediators

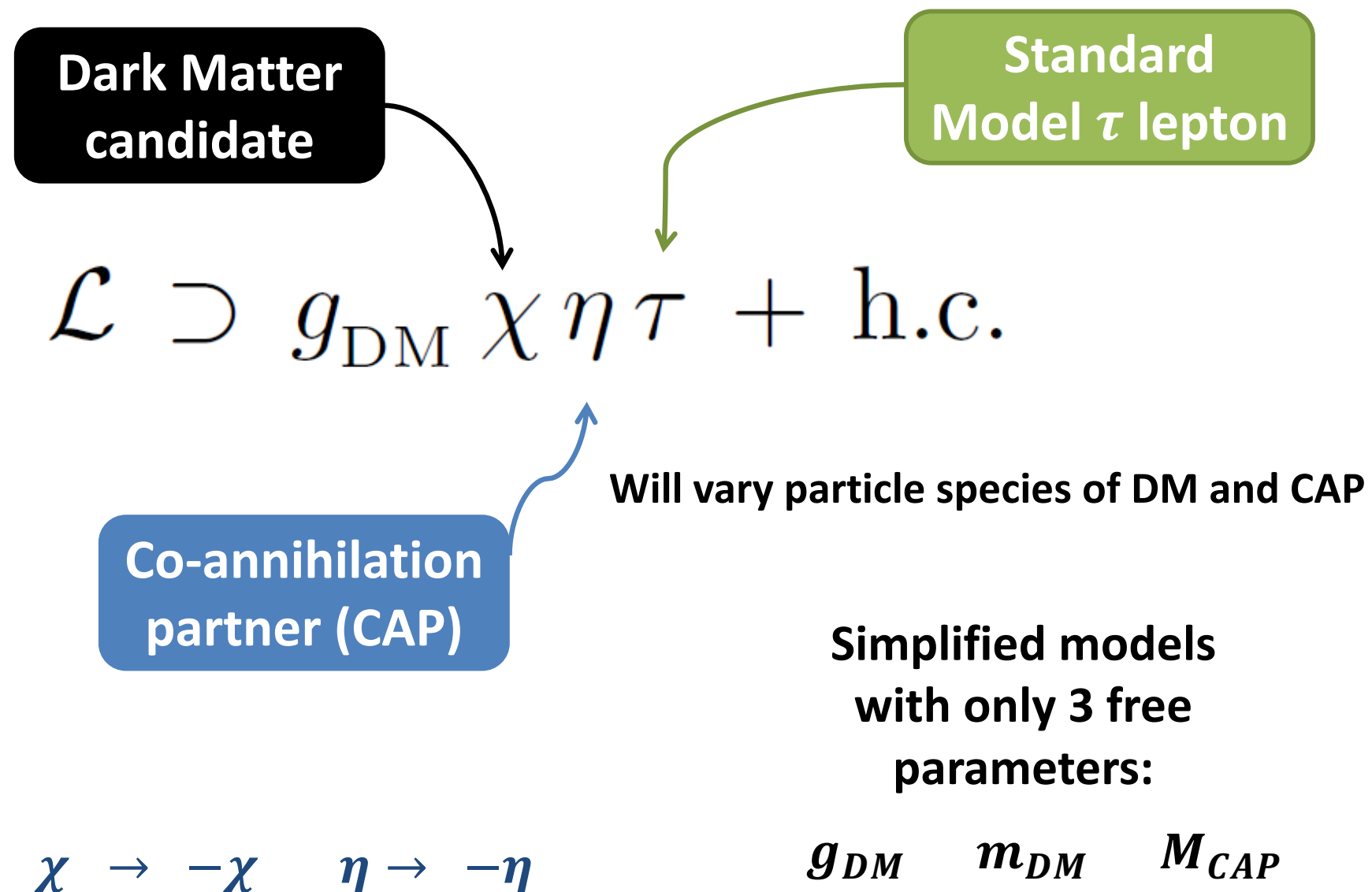
$$\boxed{\text{DM} + \text{CAP}} + \boxed{\text{SM particles}}$$

- Co-annihilation partner (CAP) has an important role in lowering the relic density
- The signal we study for collider searches is the pair-production of CAPs that then ultimately decay into cosmologically stable dark matter
- Focus on the scenario where:
 1. CAP is an SU(3)-singlet, charged under the EW sector
 2. DM and CAP are nearly mass-degenerate ($\Delta M < \tau \text{ mass} \sim 1 \text{ GeV}$)

=> the CAP is long-lived
- Compared to simplified DM models that rely on signals with missing energy at colliders, in these models the crucial collider signature to look for are tracks of long-lived electrically charged particles

- We now introduce a set of simplified models which are distinct from and complimentary to the standard mediator-based simplified DM models set:

All models involve the 3-point interactions:

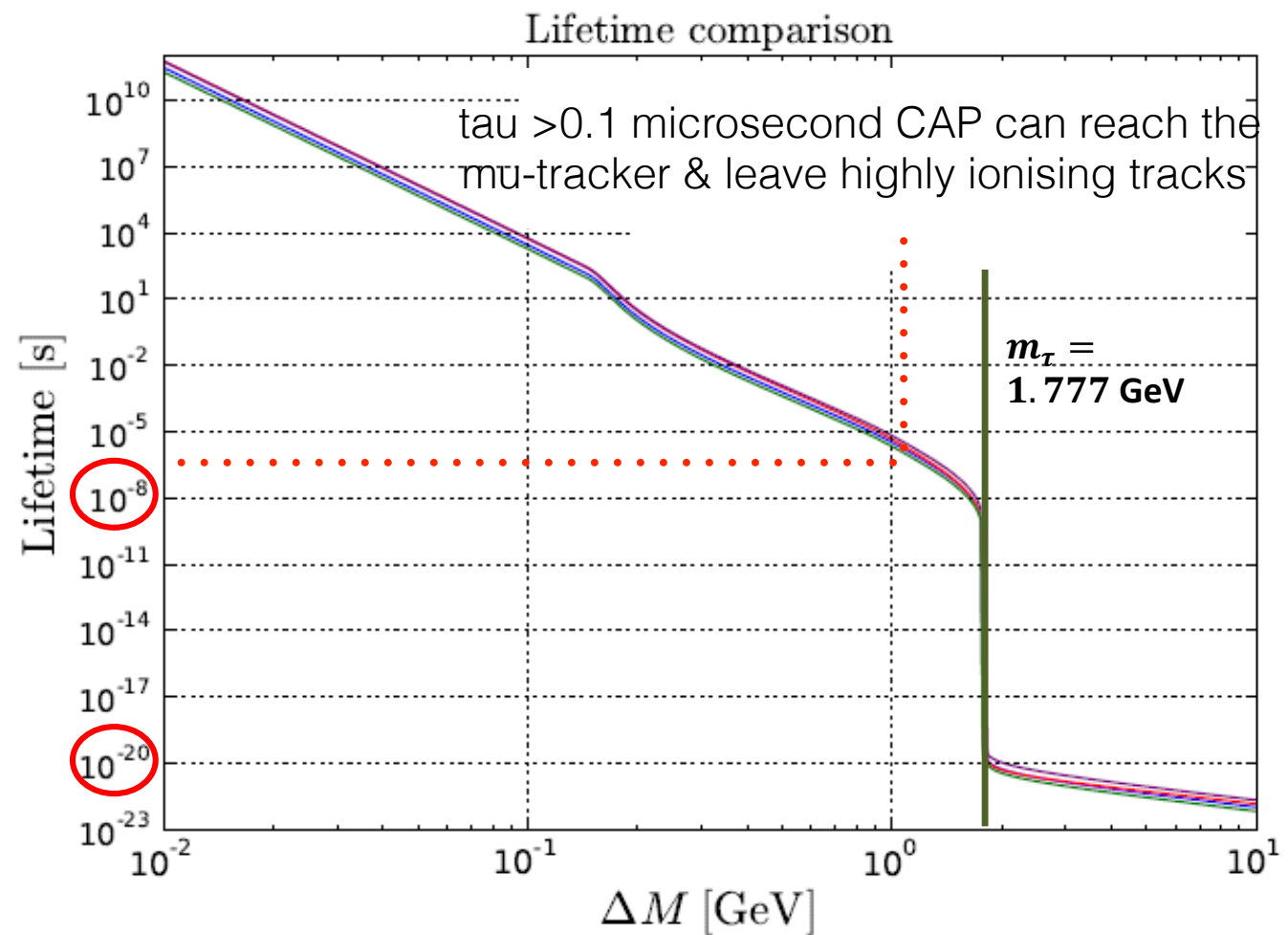
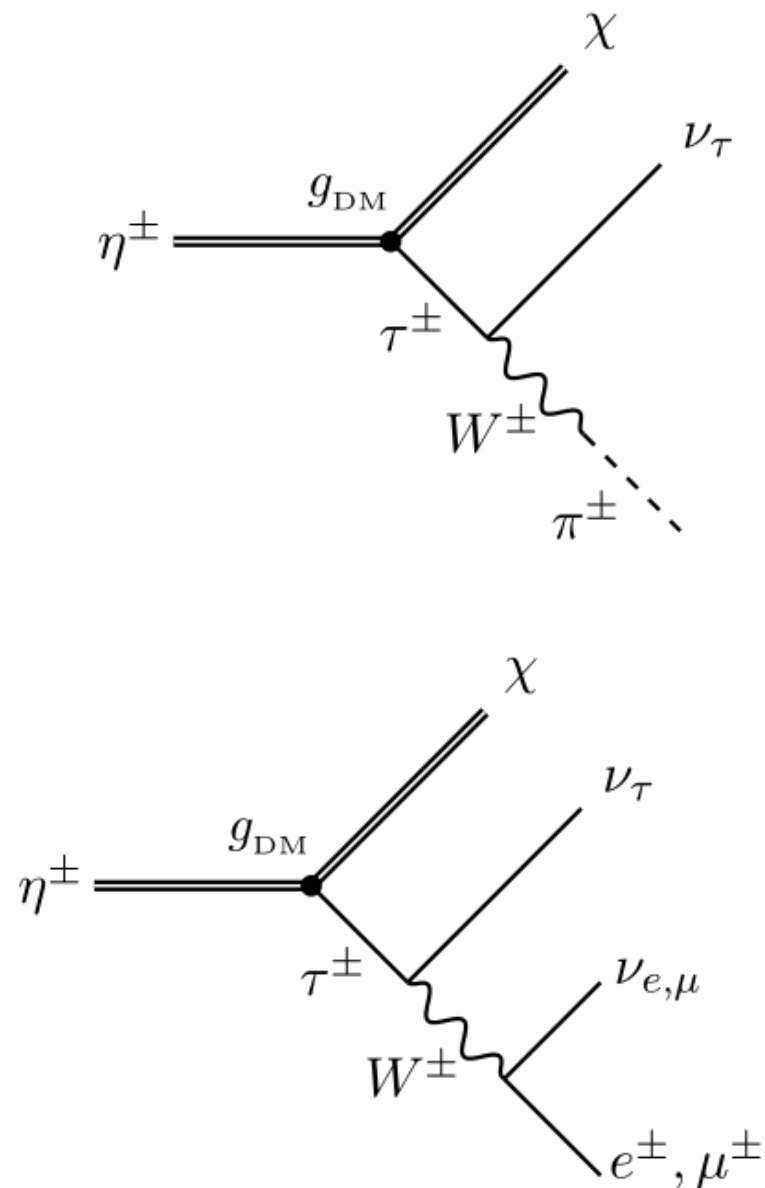


- These simplified modes can be used by ATLAS & CMS to interpret searches for long-lived charged particles to explore the new range of DM scenarios

- We now introduce a set of simplified models which are distinct from and complimentary to the standard mediator-based simplified DM models set:

+ Long-lived charged co-annihilation partners

If $\Delta M < m_\tau$ only 3-body and 4-body decays open:



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- We now introduce a set of simplified models which are distinct from and complimentary to the standard mediator-based simplified DM models set:

Fermion DM

Model-1a			
Component	Field	Charge	Interaction (5.1)
DM	Majorana fermion (χ)	$Y = 0$	$\phi^*(\chi\tau_R) + \text{h.c.}$
CAP	Complex scalar (ϕ)	$Y = -1$	

Model-1b			
Component	Field	Charge	Interaction (5.4)-(5.5)
DM	Majorana fermion (χ)	$Q = 0$	$\phi^*(\chi\tau_R) + \phi^*(\chi\tau_L) + \text{h.c.}$
CAP	Complex scalar (ϕ)	$Q = -1$	

Scalar DM

Model-2			
Component	Field	Charge	Interaction (5.6)
DM	Real scalar (S)	$Y = 0$	$S(\bar{\Psi}P_R\tau) + \text{h.c.}$
CAP	Dirac fermion (Ψ)	$Y = -1$	

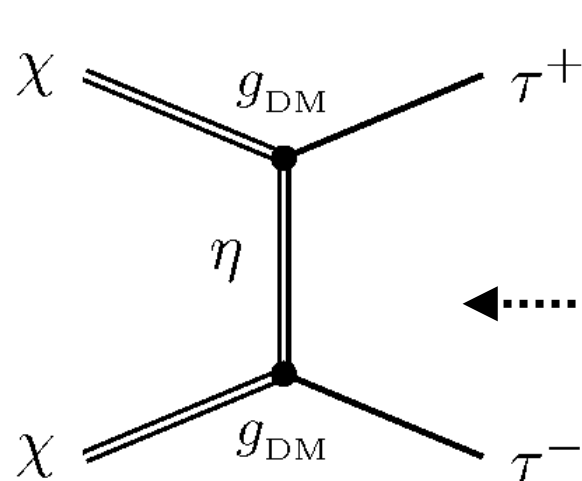
Vector DM

Model-3			
Component	Field	Charge	Interaction (5.7)
DM	Vector (V_μ)	$Y = 0$	$V_\mu(\bar{\Psi}\gamma^\mu P_R\tau) + \text{h.c.}$
CAP	Dirac fermion (Ψ)	$Y = -1$	

- These simplified modes can be used by ATLAS & CMS to interpret searches for long-lived charged particles to explore the new range of DM scenarios

Relic Density from annihilation & co-annihilation

Dark matter annihilation into pair of tau's



$$(\sigma v)_{\text{ann}}^{\text{s-wave}} = \frac{g_R^4 m_\tau^2}{32\pi m_\chi^4 (1+r^2)^2}$$

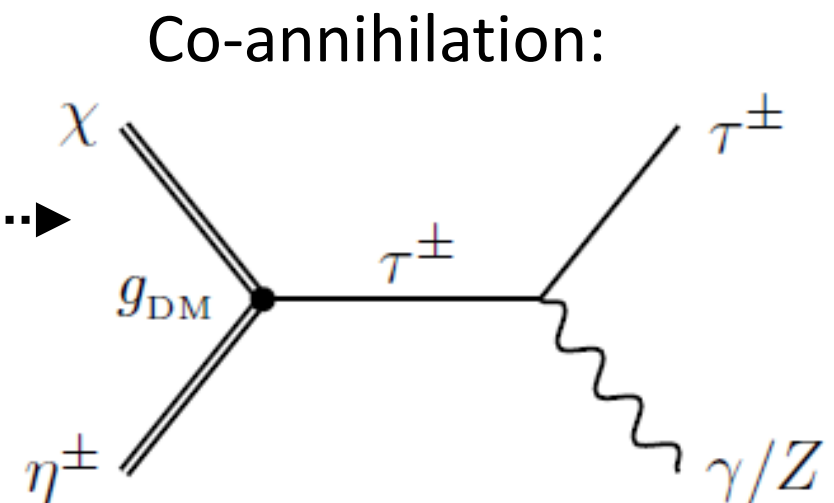
$\propto m_\tau$
Chiral suppression

- Overproduces dark matter (Unless large couplings)
- We need a mechanism to reduce the DM relic density

Freeze-out temperature $T_F \sim m_{DM}/25$

Boltzmann factor $\exp\left(-\frac{\Delta M}{T}\right) \longrightarrow \Delta M \sim m_{DM}/25$

We need **mass splitting of 4% of m_{DM}**



Relic Density from annihilation & co-annihilation

DM
 χ

CAP
 ϕ

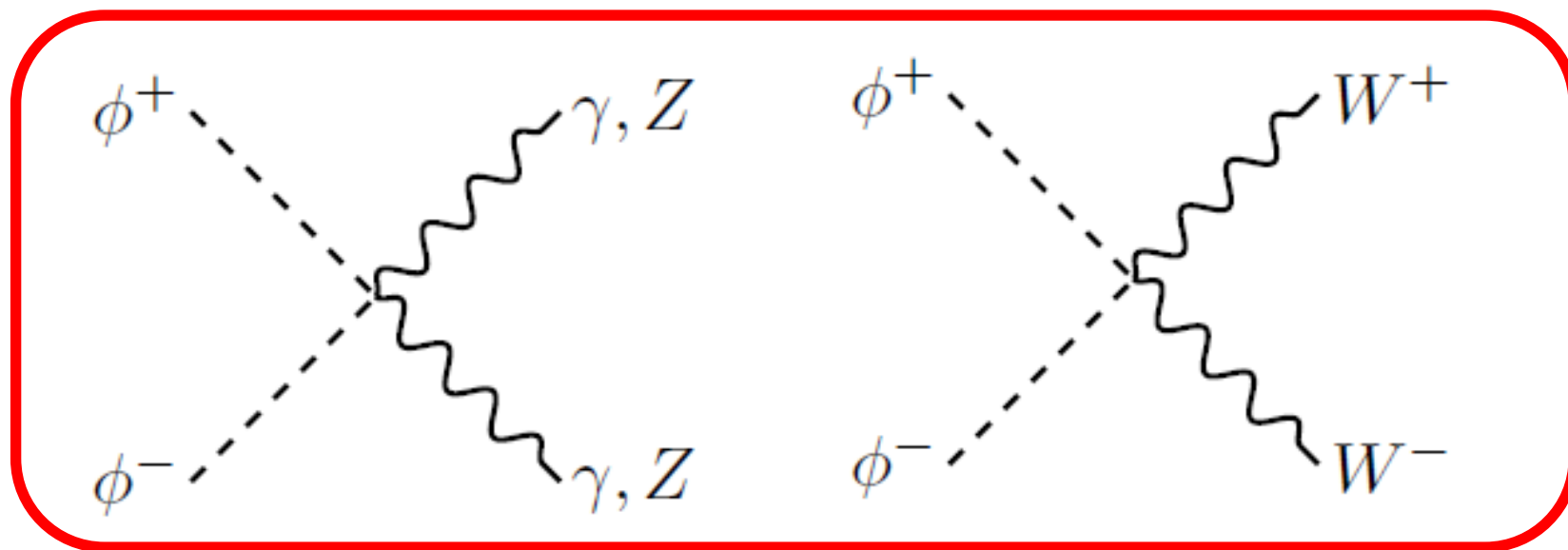
Using Majorana DM and scalar CAP as an example

dependence on g_{DM} is weak if $g_{\text{DM}} \ll 1$

$\langle \sigma_{\text{eff}} v \rangle$ dominated by $\phi^+ \phi^- \rightarrow \text{SM particles}$

Dominant contributions come from:

CAP-CAP



SM-SM

and that is why CAPs are important for cosmology

Relic Density from annihilation & co-annihilation

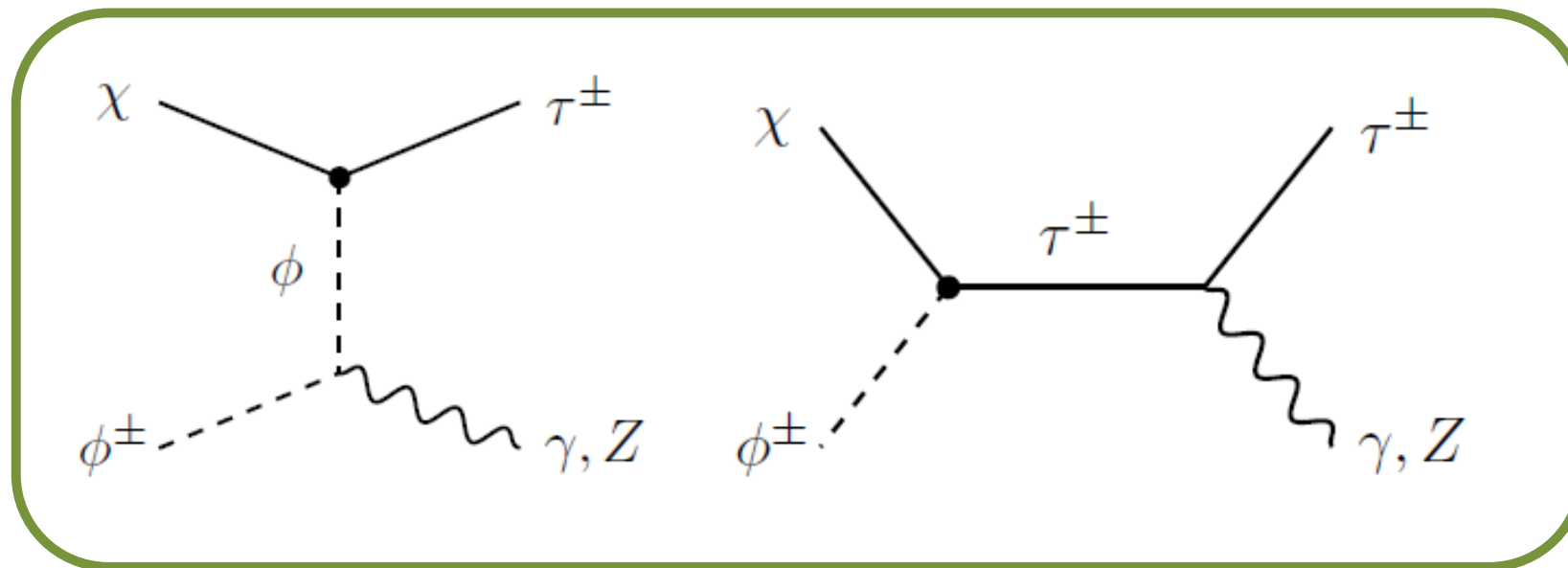
$$g_{\text{DM}} \gtrsim 0.3$$

$$\phi^\pm \chi \rightarrow \text{SM particles}$$

Becomes relevant

$$g'^2 g_{\text{DM}}^2$$

CAP-DM



SM-SM

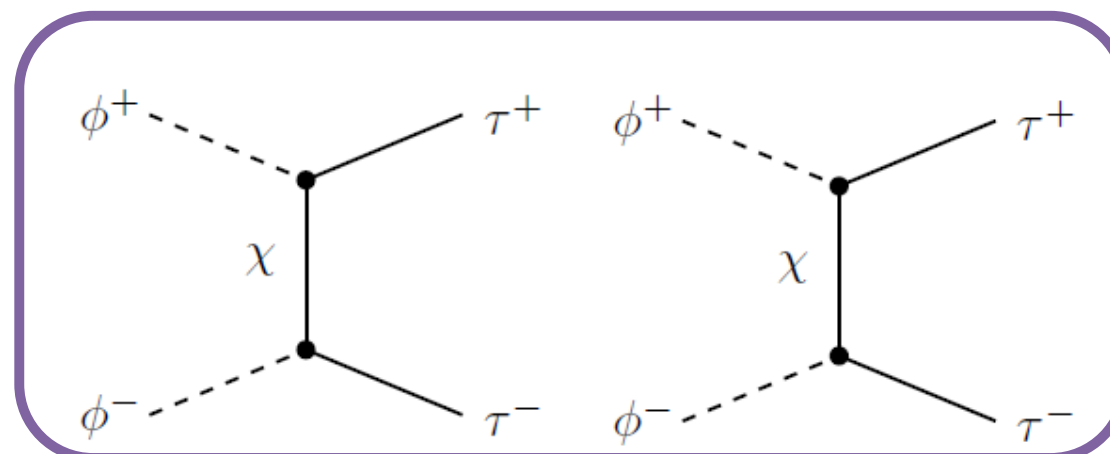
$$g_{\text{DM}} \gtrsim 1.0$$

$$\phi^+ \phi^+ \rightarrow \tau^+ \tau^+$$

Relevant

$$g_{\text{DM}}^4$$

CAP-CAP



SM-SM

Relic Density from annihilation & co-annihilation

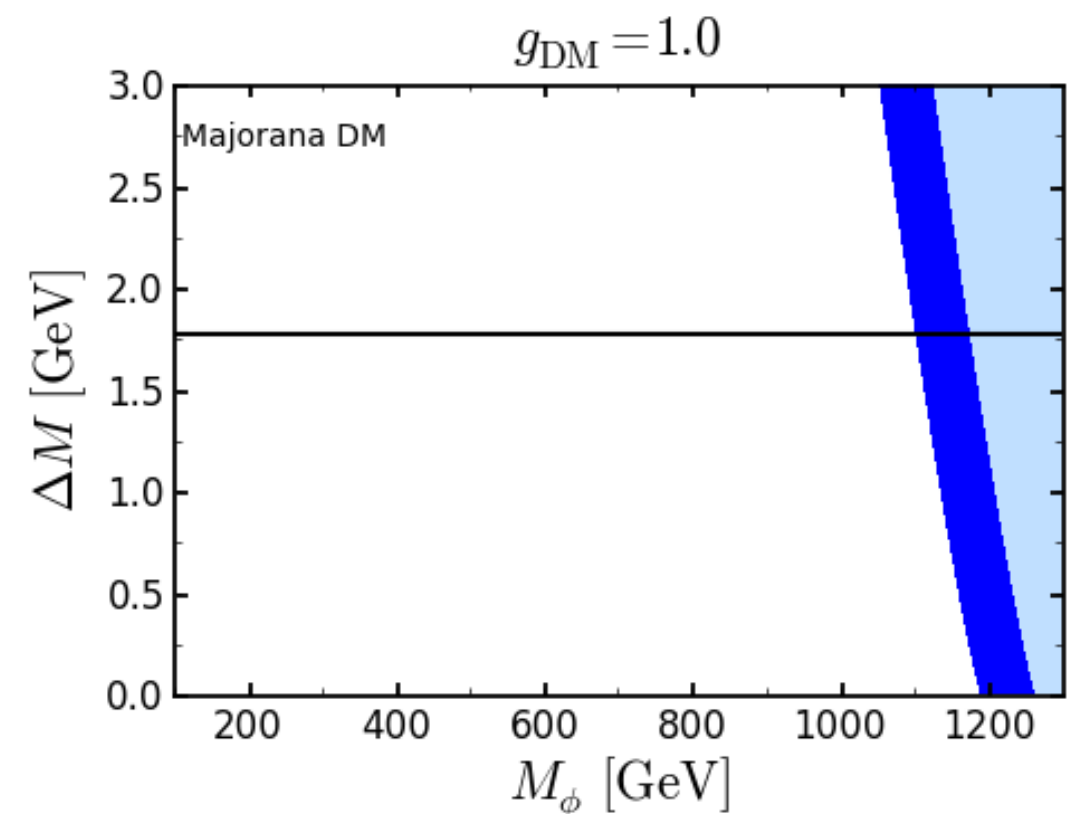
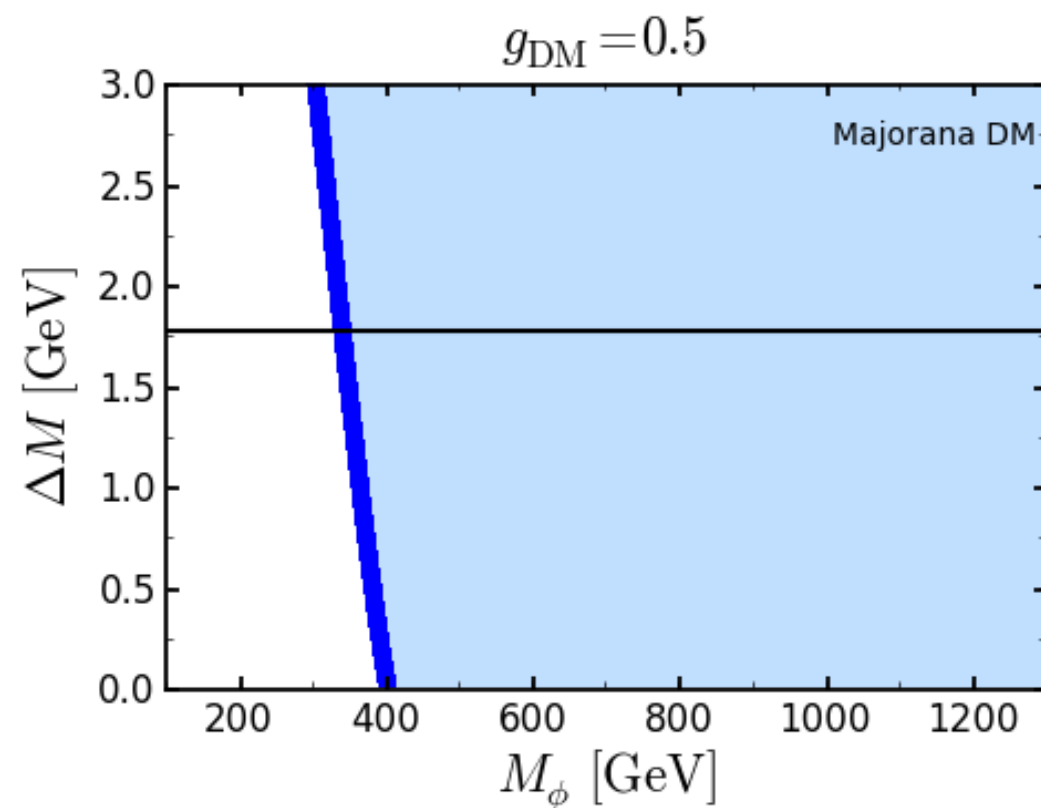
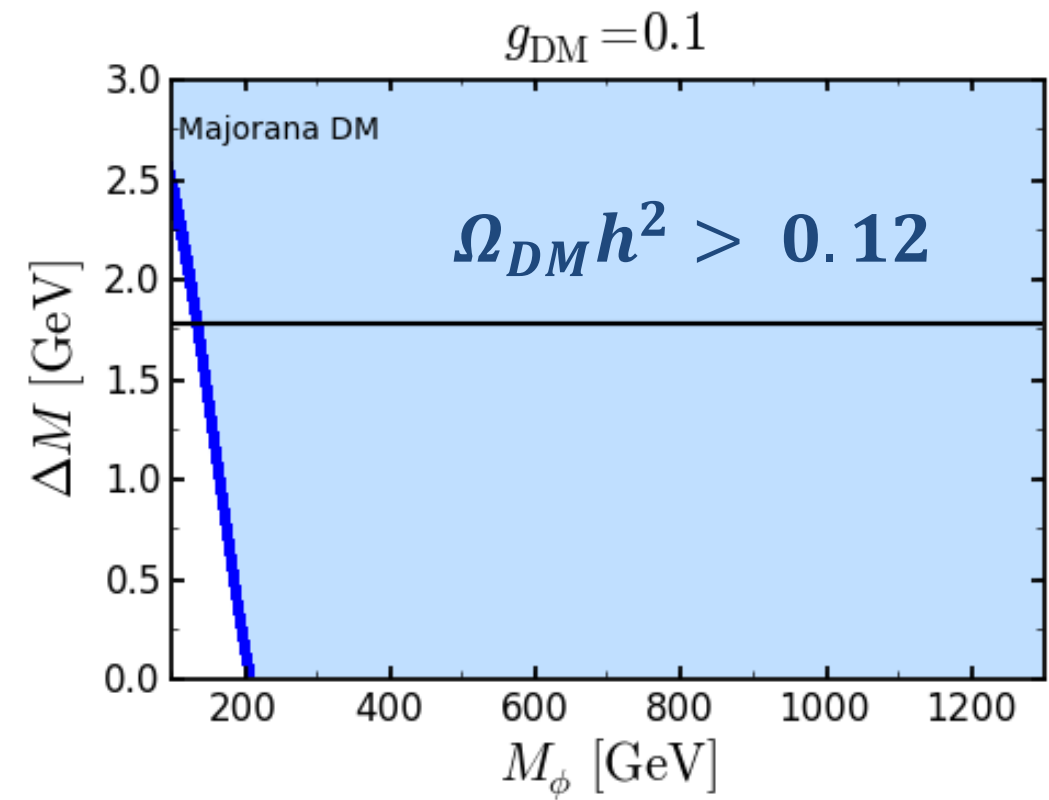
DM

χ

CAP

ϕ

$$\phi^* (\chi \tau_R) \subset \mathcal{L}$$



Model 1a

- Dark Matter = Majorana fermion
- Co-annihilation Partner = complex scalar charged under U(1)_Y only
 $(\eta^+, \eta^-) = (\phi^*, \phi) = (\phi^+, \phi^-)$

$$\begin{aligned}\mathcal{L} &= \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{DM}} + \mathcal{L}_{\text{CAP}} + \mathcal{L}_{\text{int}}, \\ \mathcal{L}_{\text{DM}} &= \frac{1}{2} \chi (i \not{\partial} - m_{\text{DM}}) \chi, \\ \mathcal{L}_{\text{CAP}} &= |D_\mu \phi|^2 - M_\phi^2 |\phi|^2, \\ \mathcal{L}_{\text{int}} &= g_{\text{DM}} \phi^* \chi \tau_R + \text{h.c.},\end{aligned}$$

3 parameters:
 $g_{\text{DM}}, m_{\text{DM}}, \Delta M < m_\tau$

- No SU(2) (and no SU(3)) interactions with the Dark Sector
- Can be the *Bino + right-handed Stau* co-annihilation strip in SUSY
- More importantly: this Simplified Model is interesting on its own right - it does not require a UV completion.

Model 1a

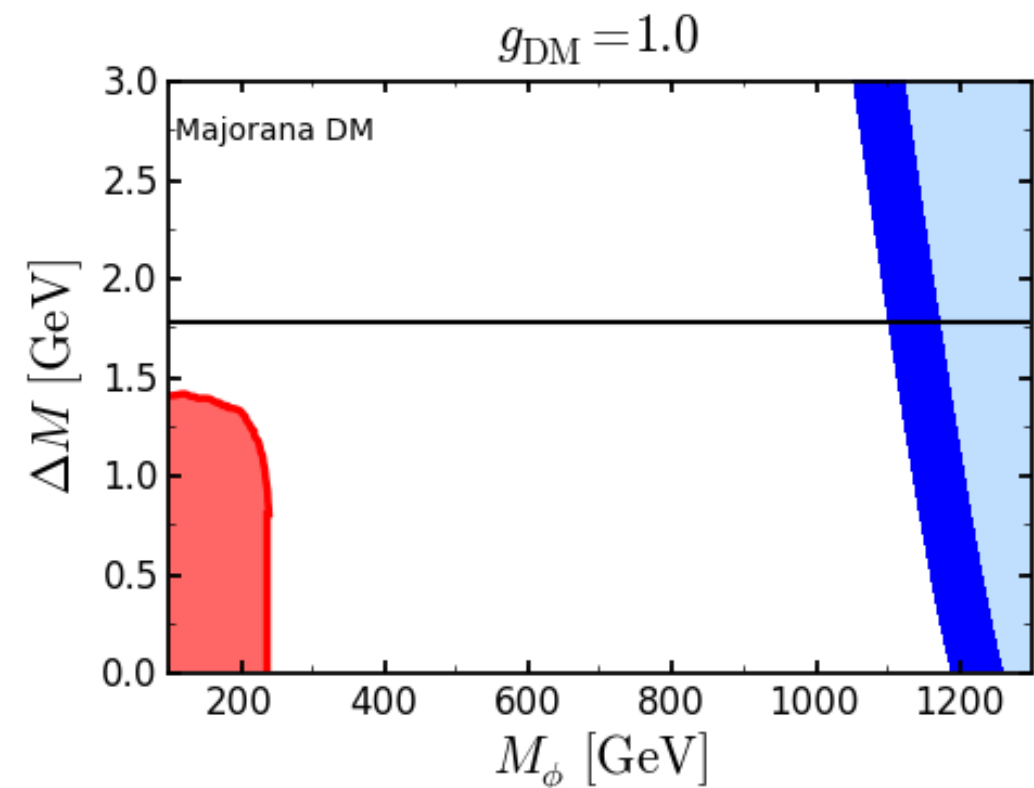
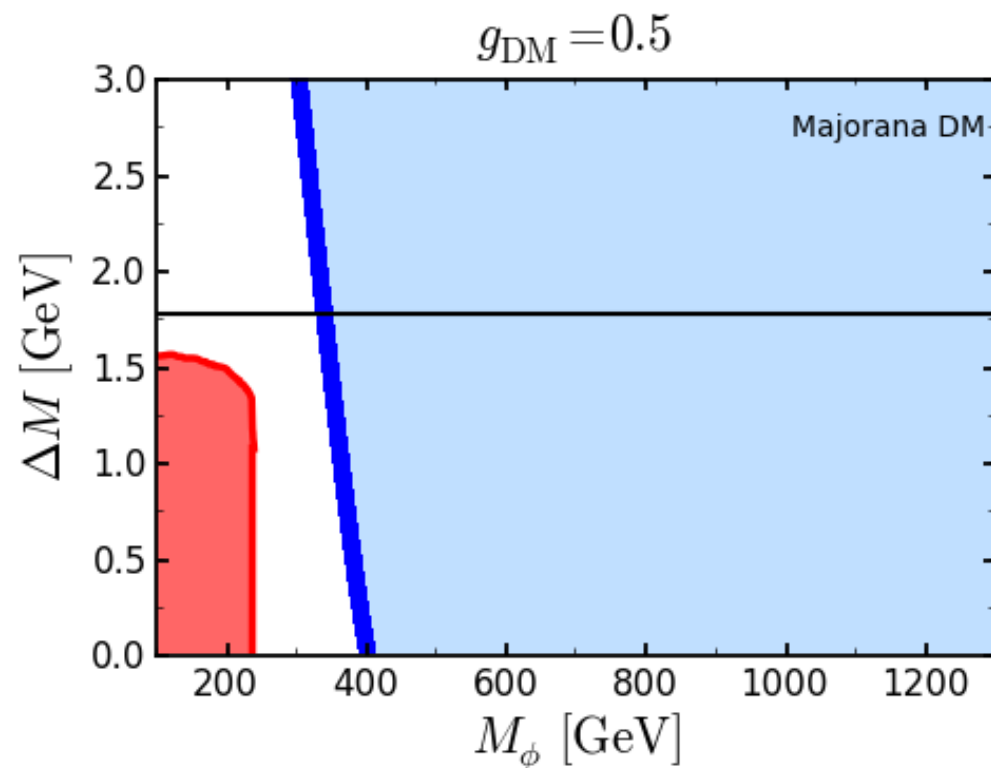
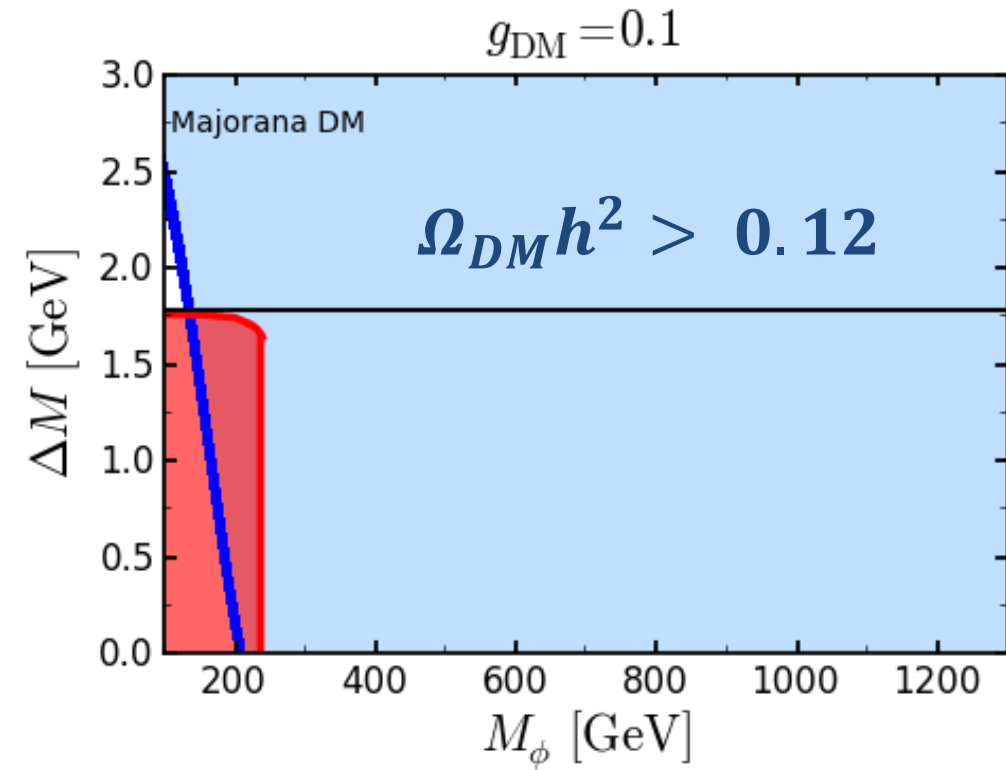
Majorana Scalar

DM **CAP** ($Y = 1$ $L_\tau = 1$)

χ ϕ

$$\phi^* (\chi \tau_R) \subset \mathcal{L}$$

ATLAS & CMS
8 TeV 18.8 fb⁻¹



Model 1a

Majorana Scalar

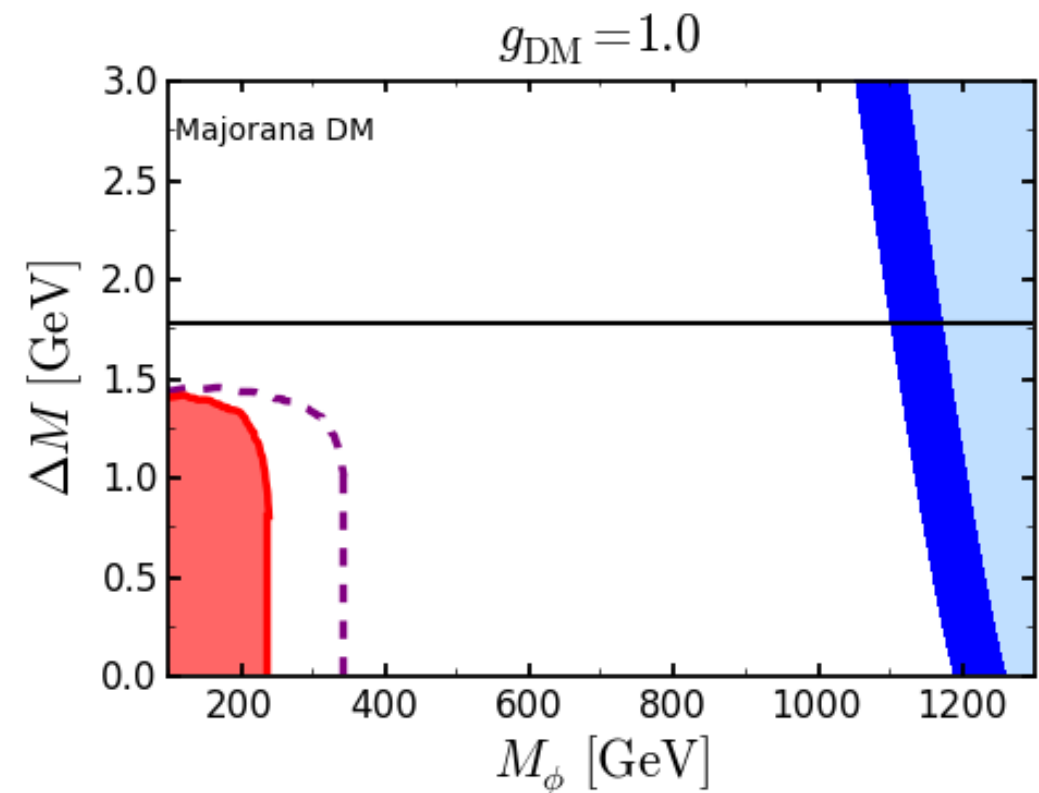
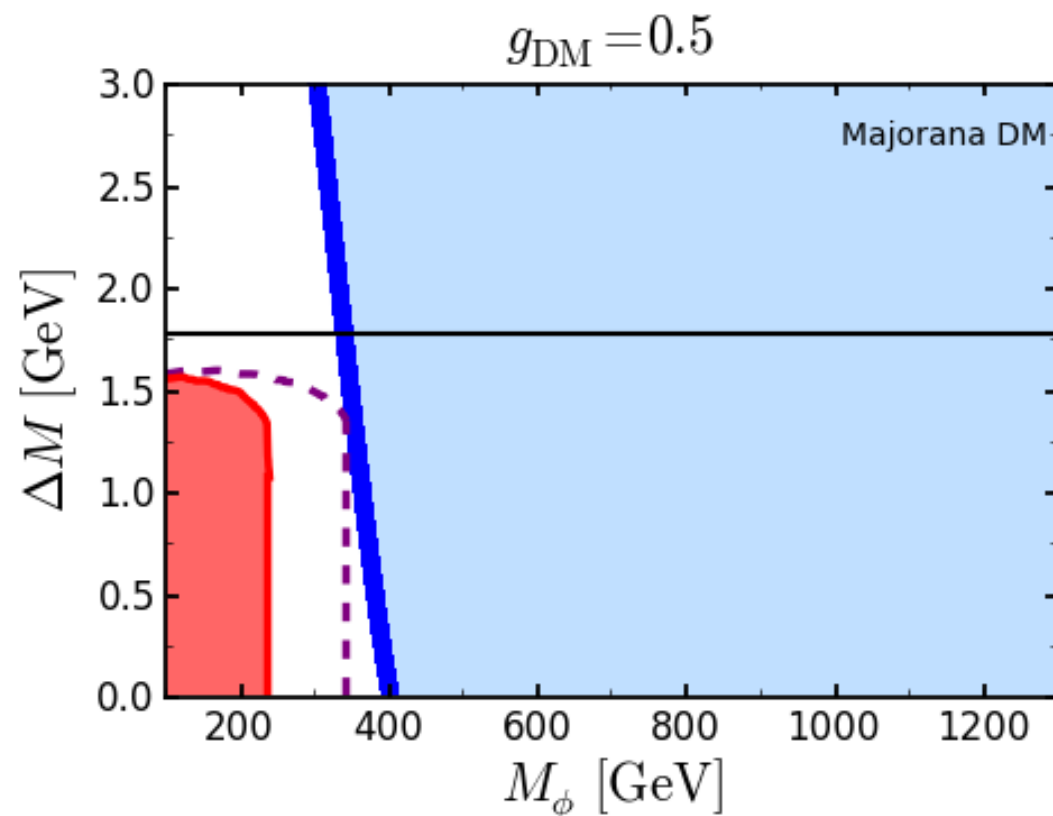
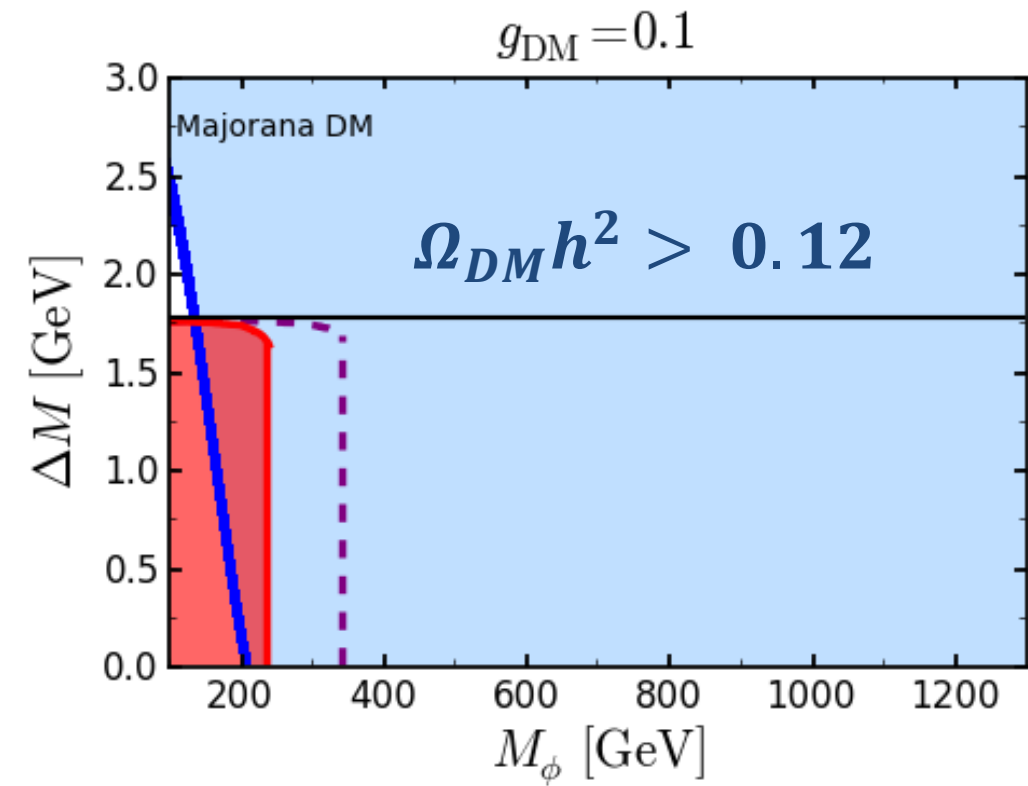
DM **CAP** ($Y = 1$ $L_\tau = 1$)

χ

ϕ

$$\phi^* (\chi \tau_R) \subset \mathcal{L}$$

13 TeV 30 fb⁻¹



Model 1a

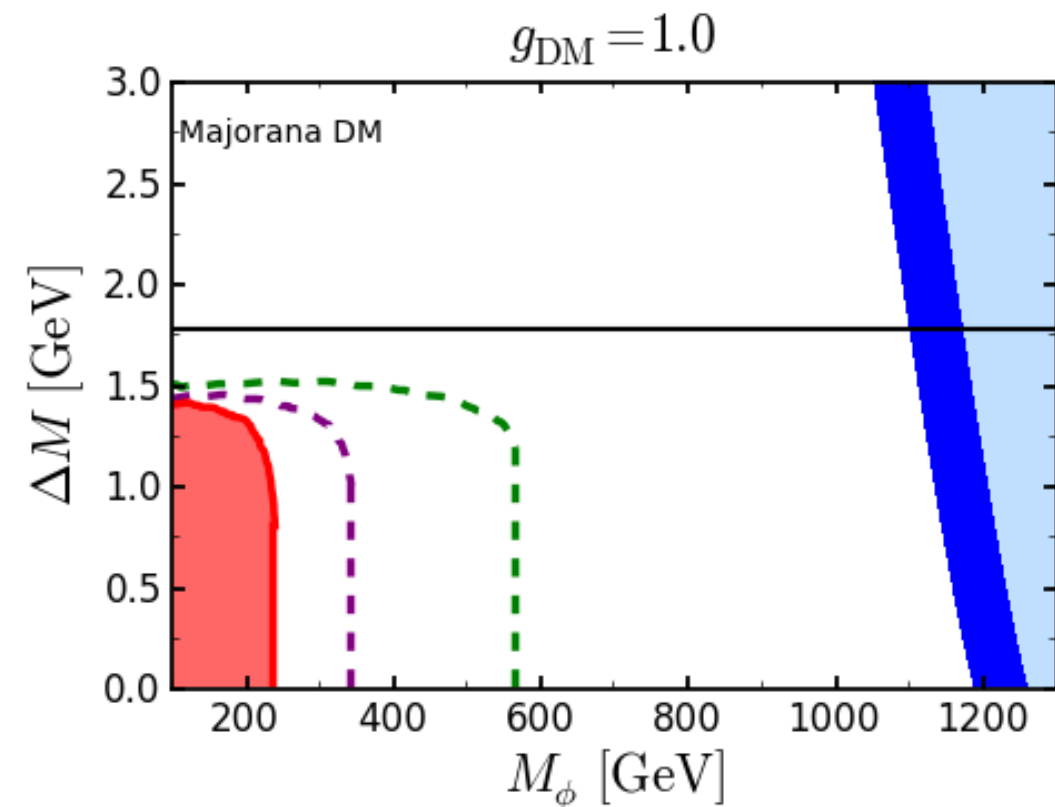
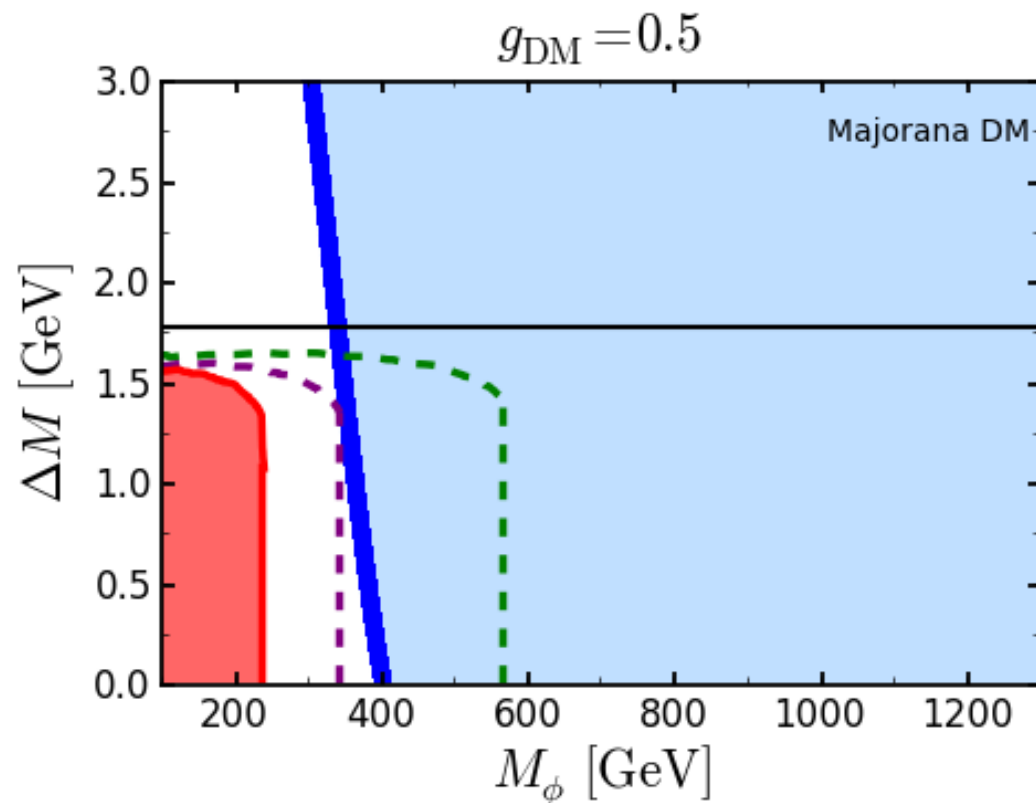
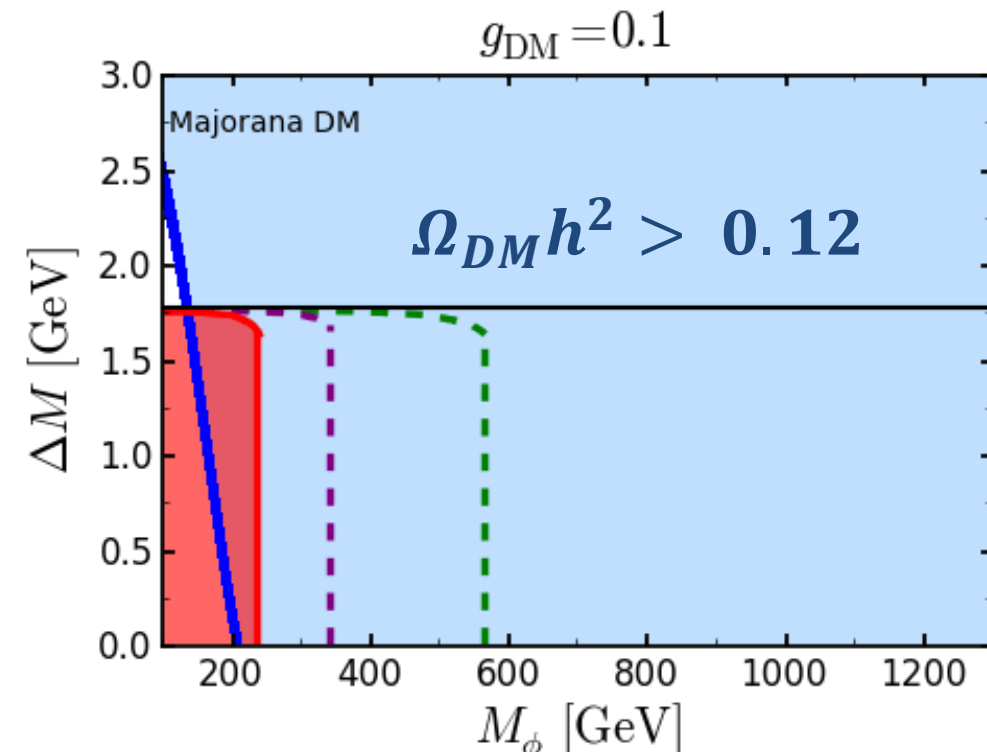
Majorana Scalar

DM **CAP** ($Y = 1$ $L_\tau = 1$)

χ ϕ

$$\phi^* (\chi \tau_R) \subset \mathcal{L}$$

13 TeV 300 fb⁻¹



Model 1a

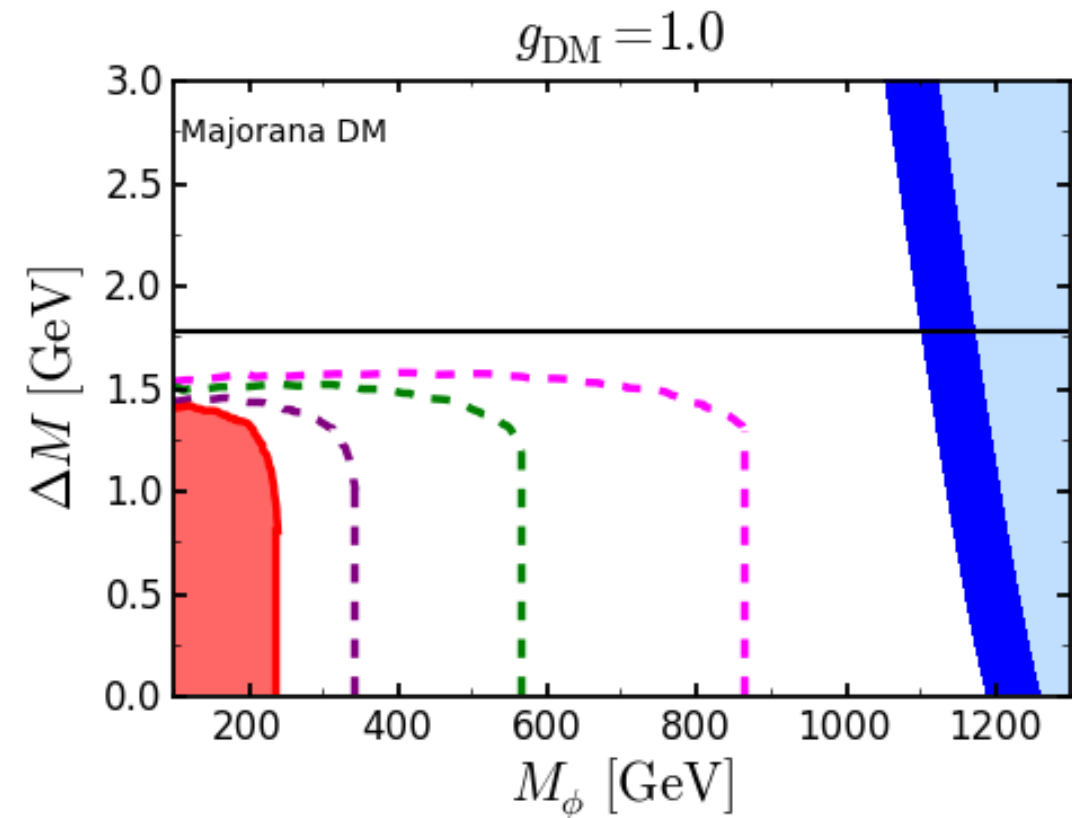
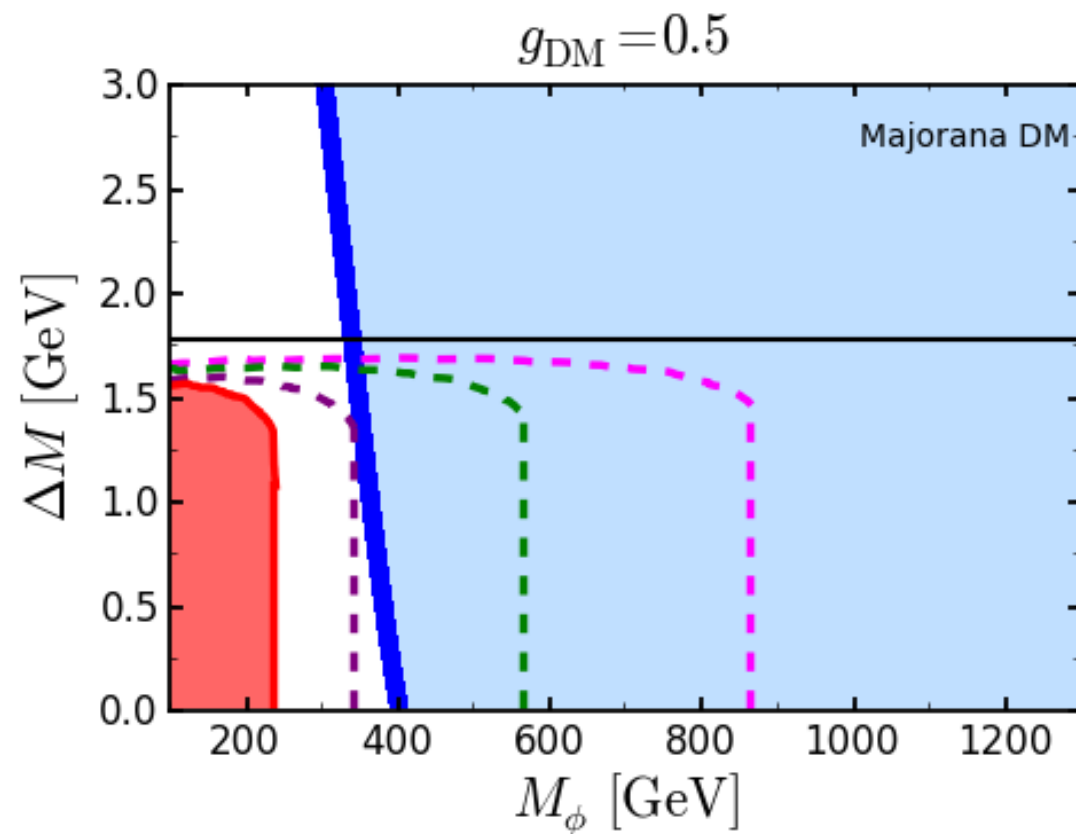
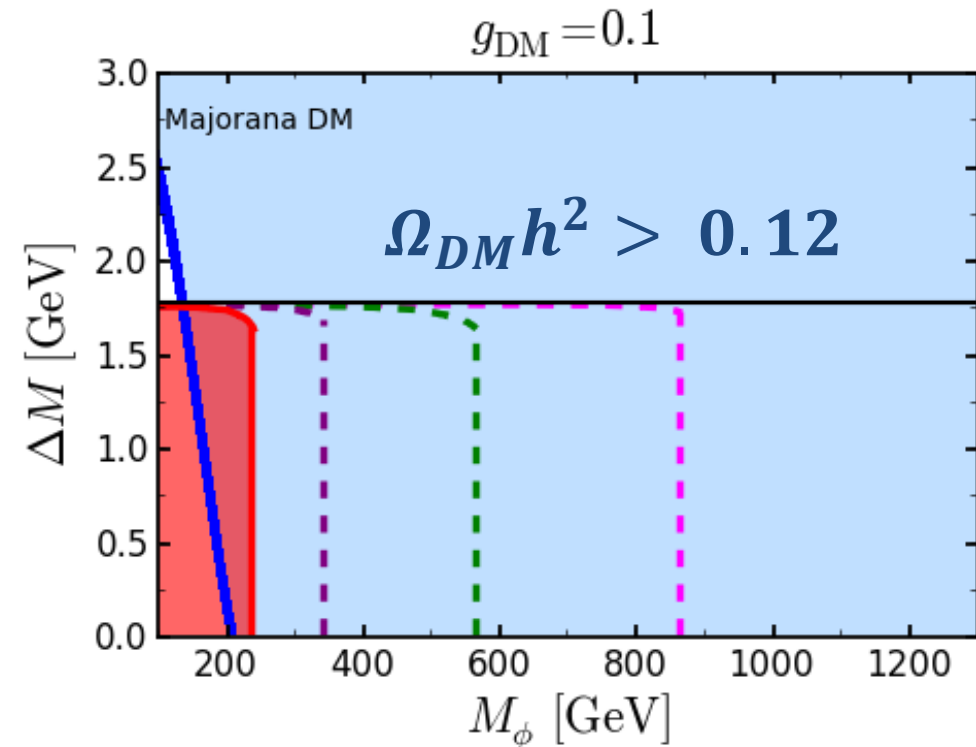
Majorana Scalar

DM **CAP** ($Y = 1$ $L_\tau = 1$)

χ ϕ

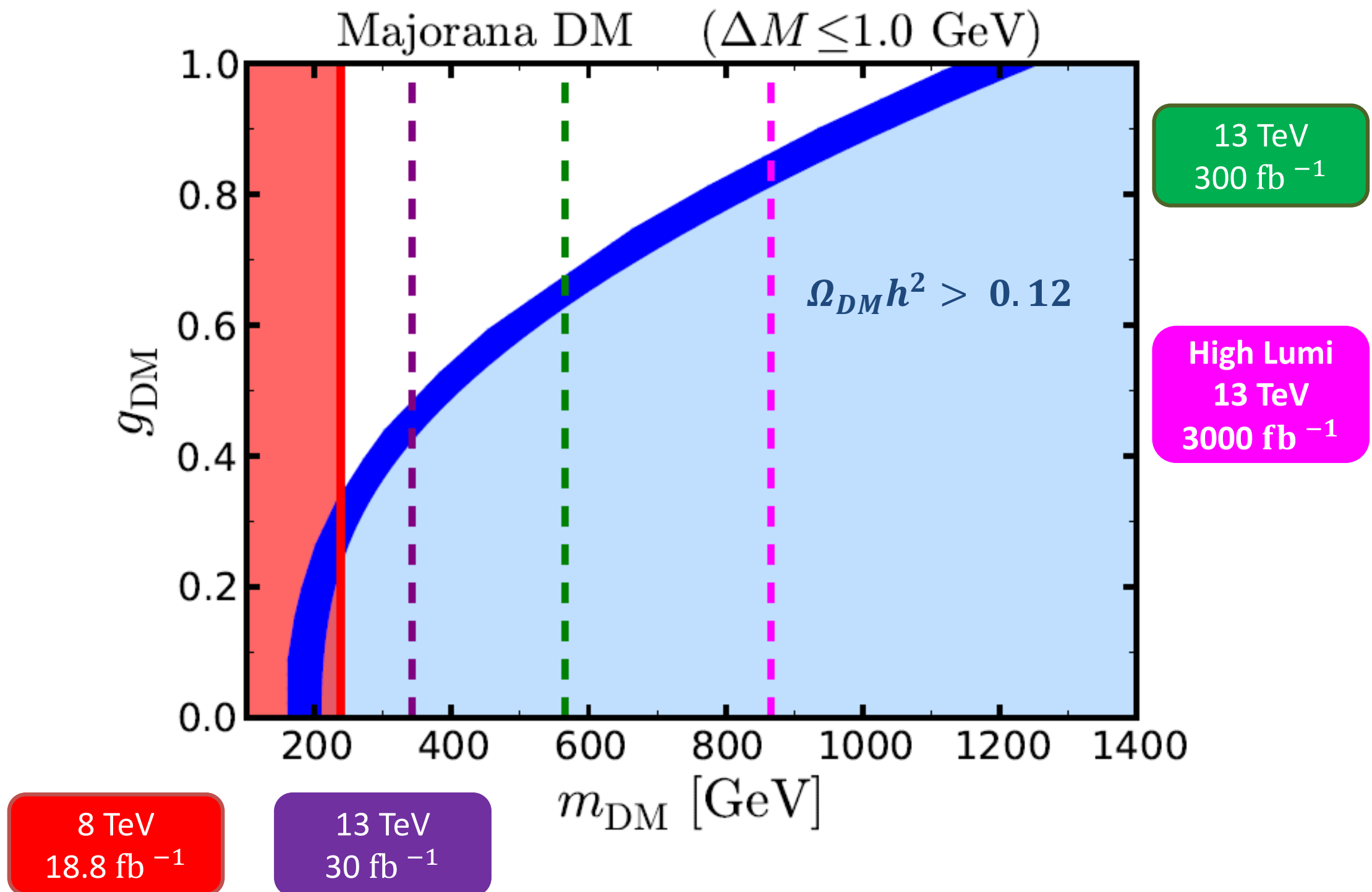
$$\phi^* (\chi \tau_R) \subset \mathcal{L}$$

High Lumi 13 TeV 3000 fb⁻¹



Model 1a

Now on the (g_{DM} , m_{DM}) plane:



Model 1b

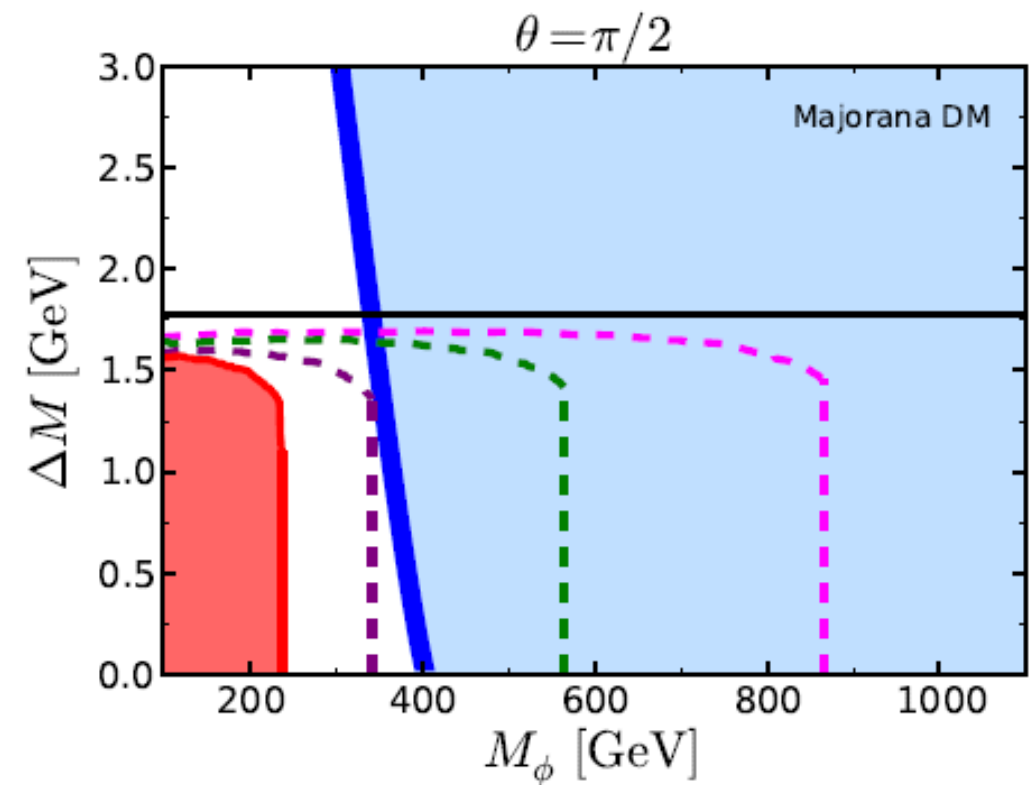
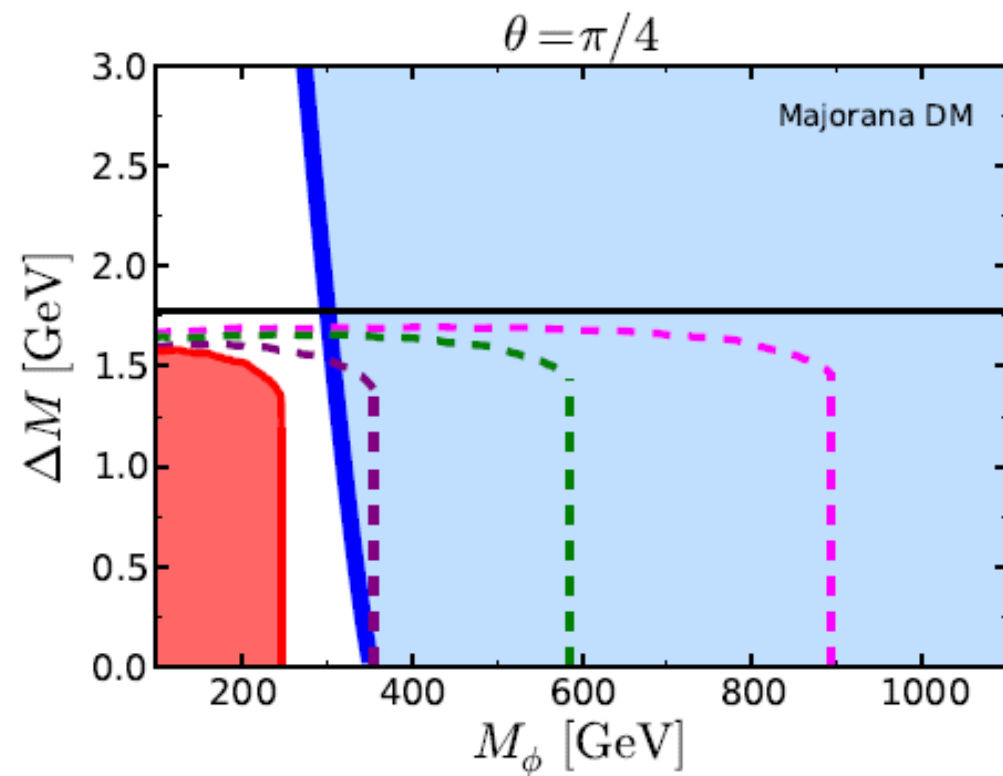
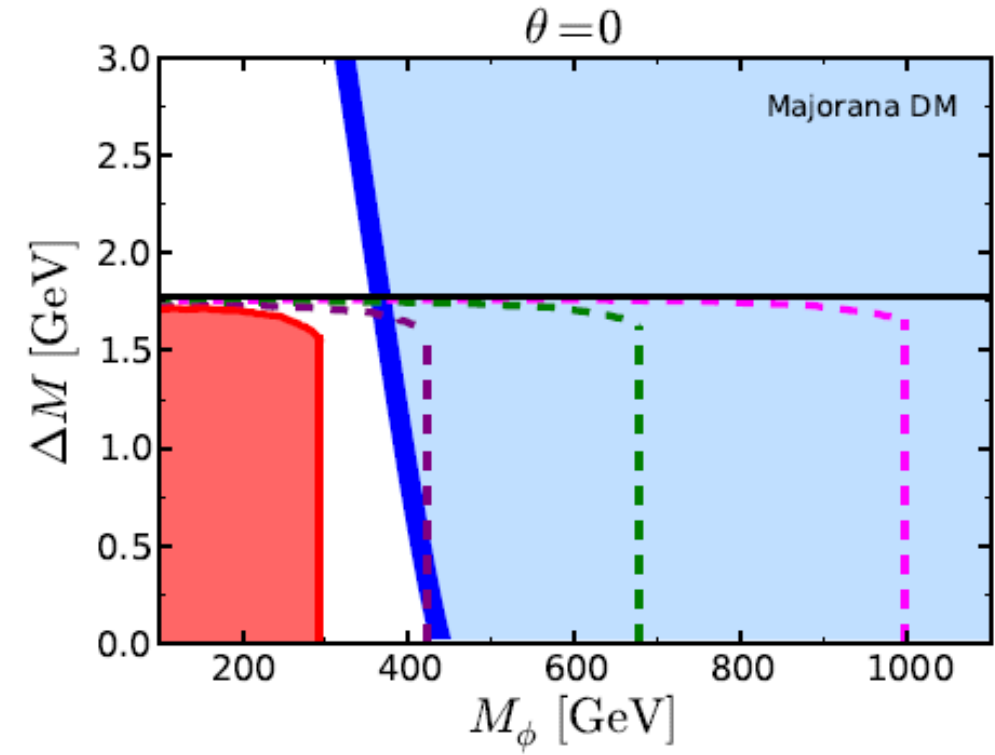
Majorana Scalar

DM **CAP (Q=-1)**
 χ ϕ

$$g_R \phi^* (\chi \tau_R) + g_L \phi^* (\chi \tau_L) \subset \mathcal{L}$$

$$\phi = \cos \theta \phi_L + \sin \theta \phi_R$$

$$g_L = \frac{1}{\sqrt{2}} g' \cos \theta, \quad g_R = -\sqrt{2} g' \sin \theta$$



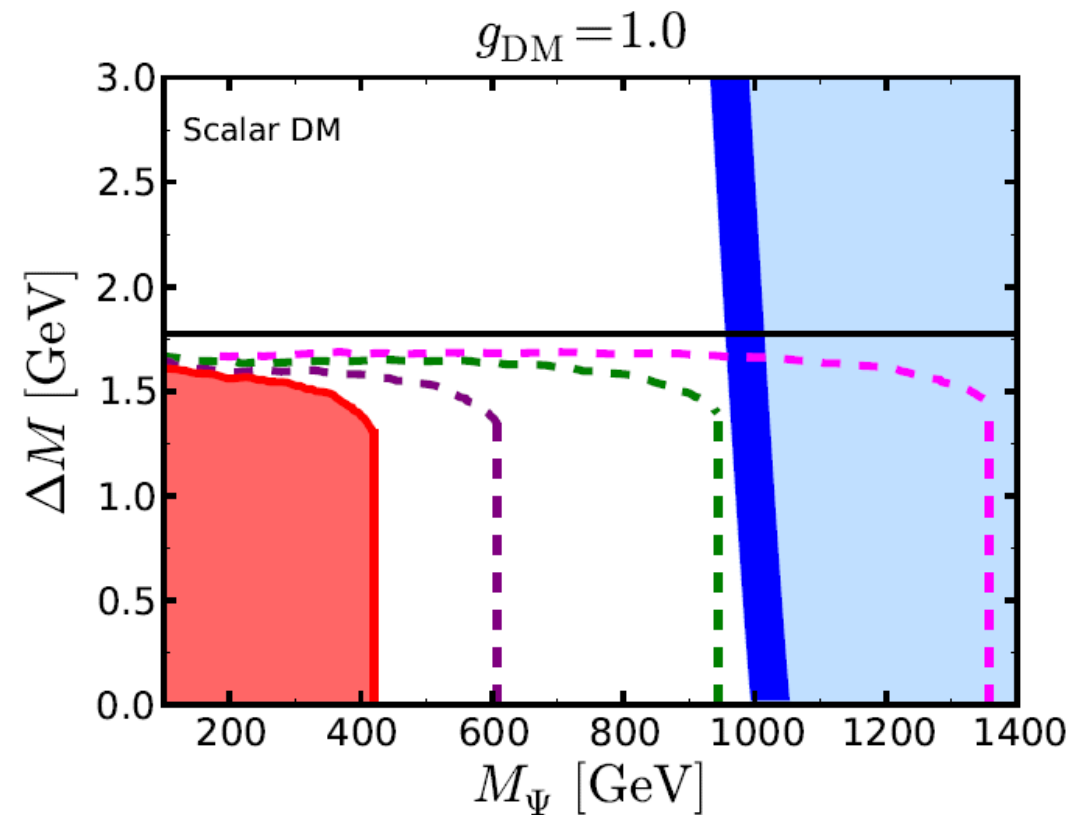
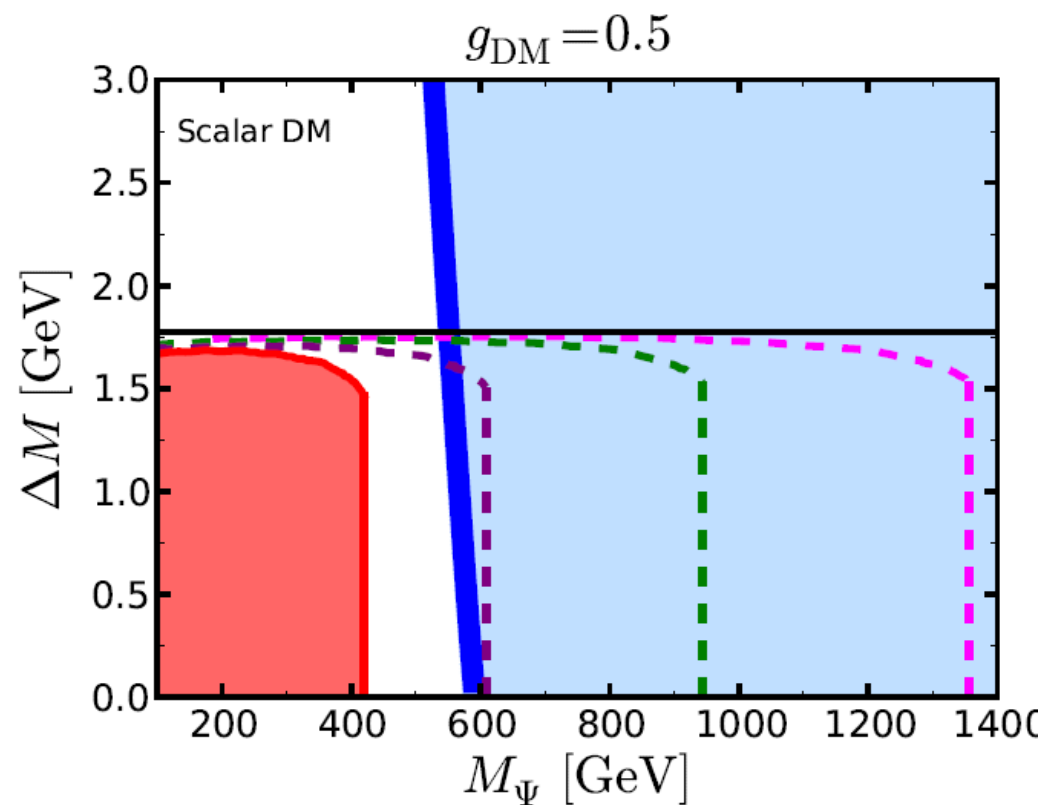
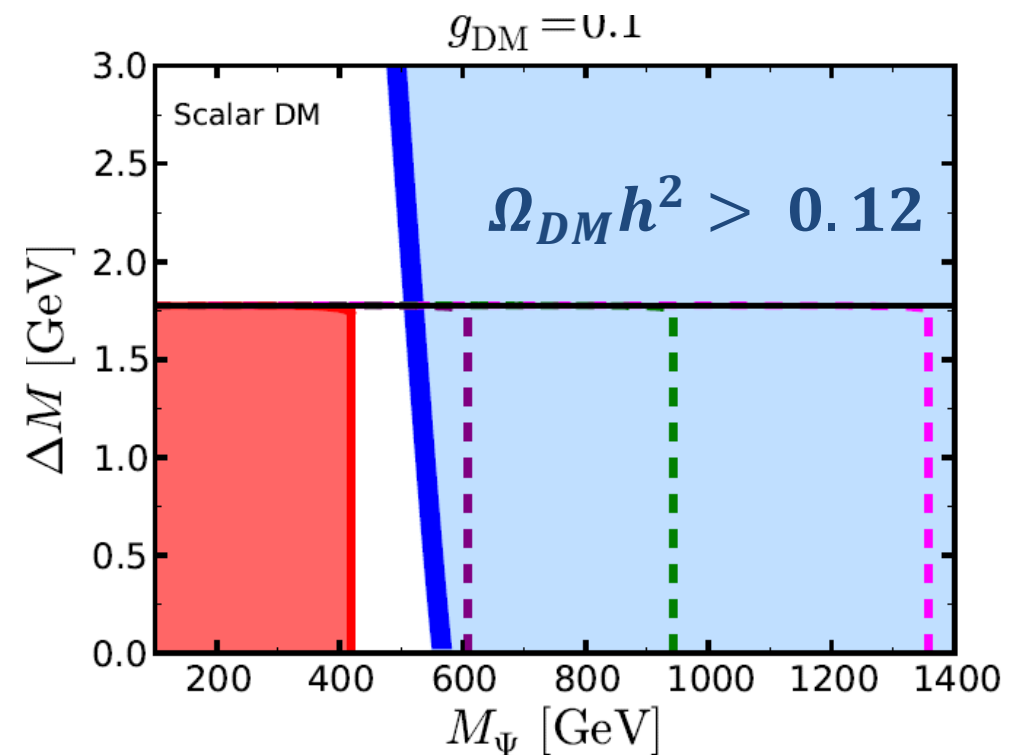
Model 2

R Scalar Dirac ferm

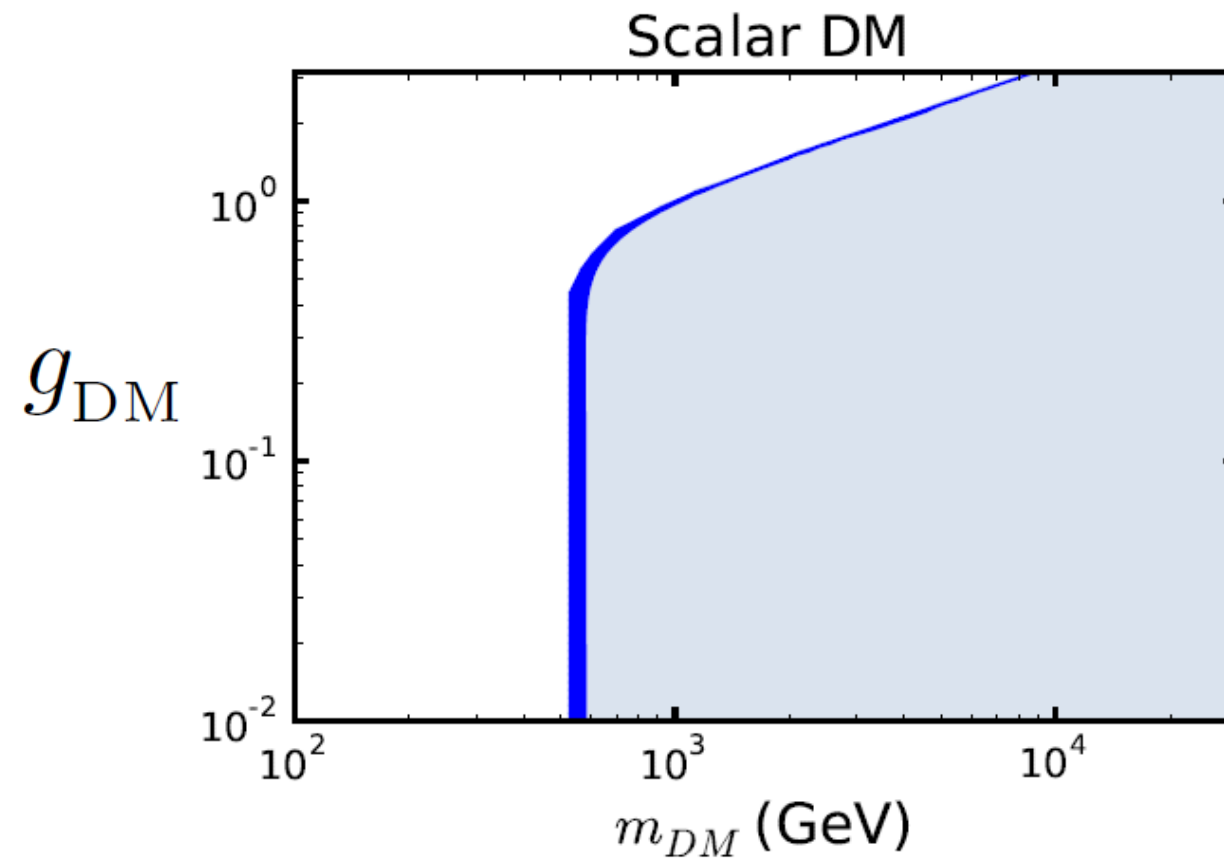
DM **CAP** ($Y = 1$ $L_\tau = 1$)
 S **Ψ**

$$\mathcal{S}(\bar{\Psi} \tau_R) \subset \mathcal{L}$$

**Gauge-invariant and renormalizable,
no problems of unitarity**



Model 2

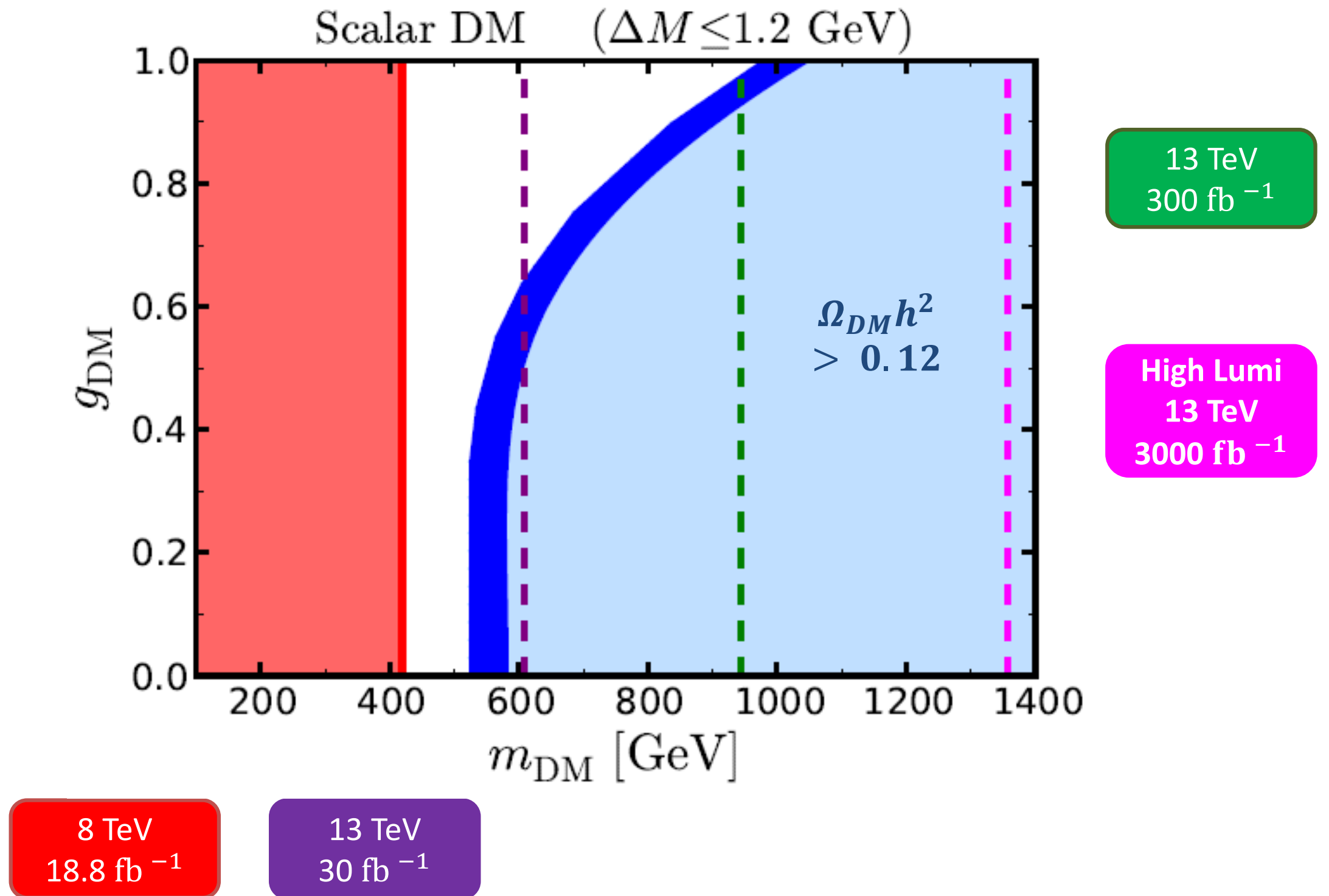


Intermediate region
 $0.6 < g_{DM} < 0.8$
 co-annihilations are
 relevant

$g_{DM} \gtrsim 0.6$ $\left\{ \begin{array}{l} \Psi\bar{\Psi} \text{ annihilation into SM final states} \\ \text{is the dominant contribution to } \Omega_{DM} h^2 \end{array} \right.$ Independent of g_{DM}

$g_{DM} \gtrsim 0.8$ $\left\{ \begin{array}{l} \Psi\Psi \rightarrow \tau\tau \\ \Psi\bar{\Psi} \rightarrow \tau\bar{\tau} \end{array} \right.$ Cross-section goes as g_{DM}^4

Model 2



Model 3

Can be KK photon

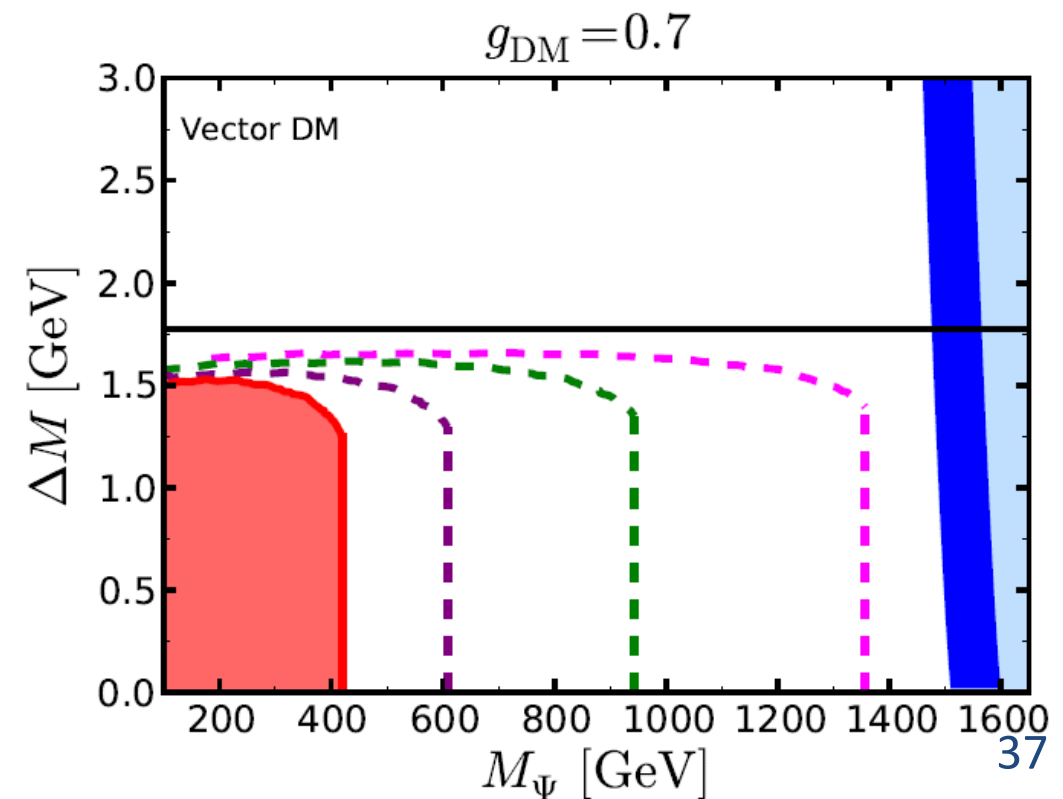
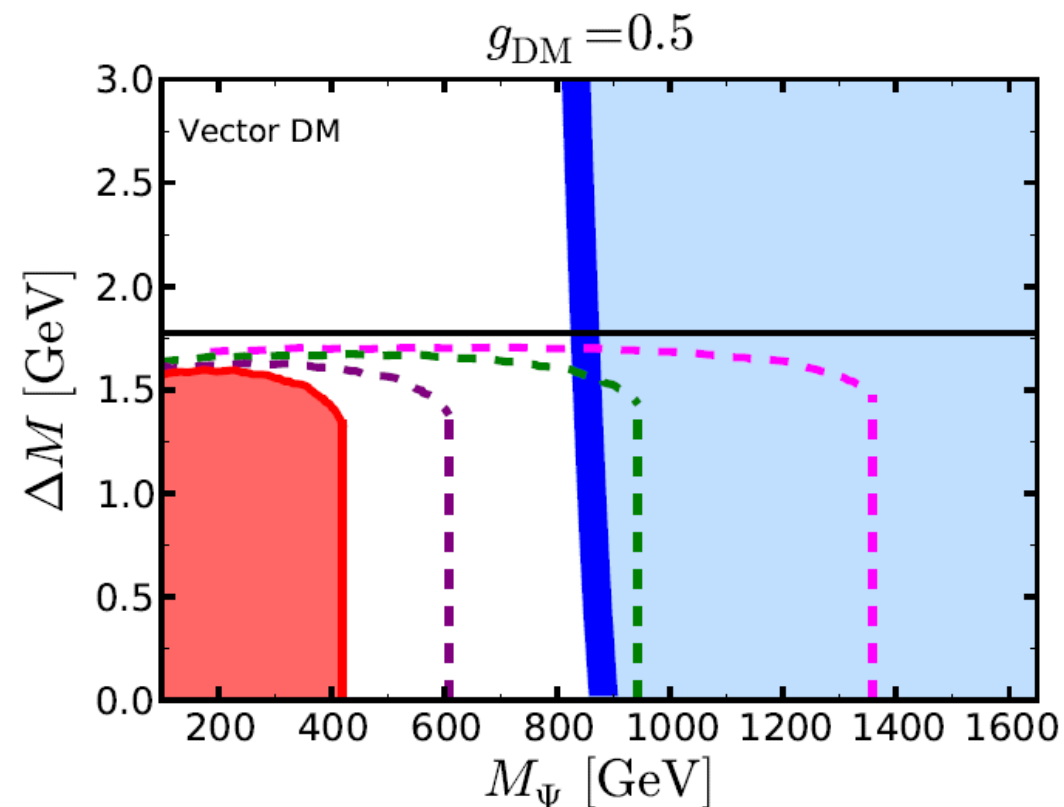
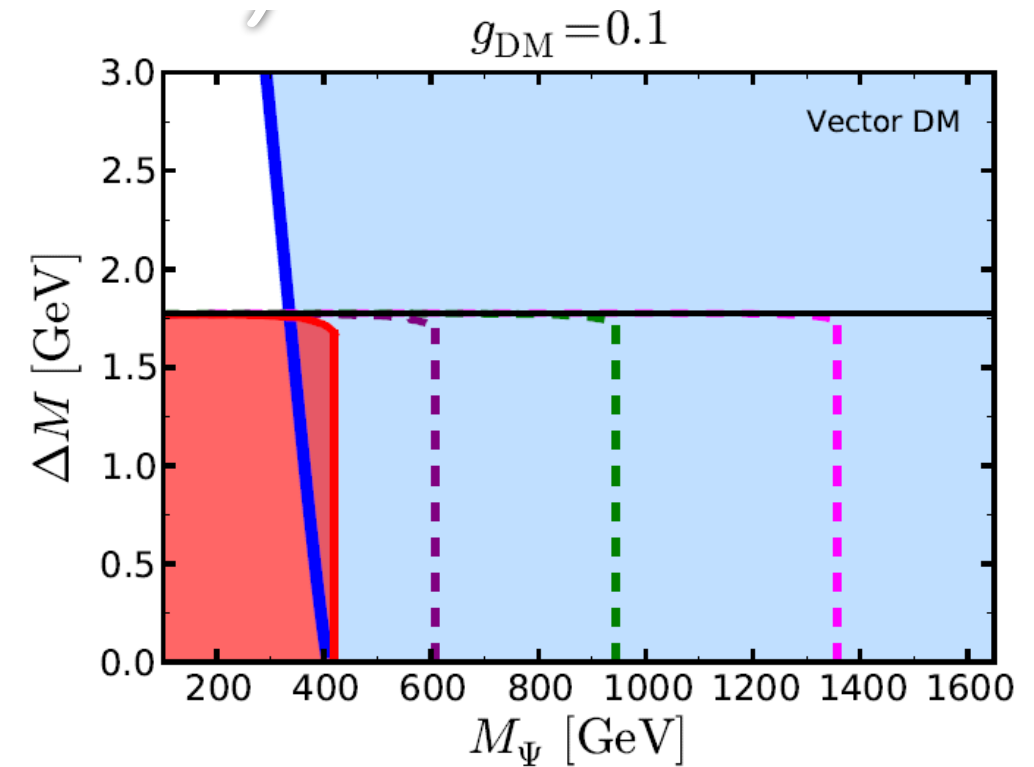
Vector Dirac ferm

DM **CAP** ($Y = 1$ $L_\tau = 1$)

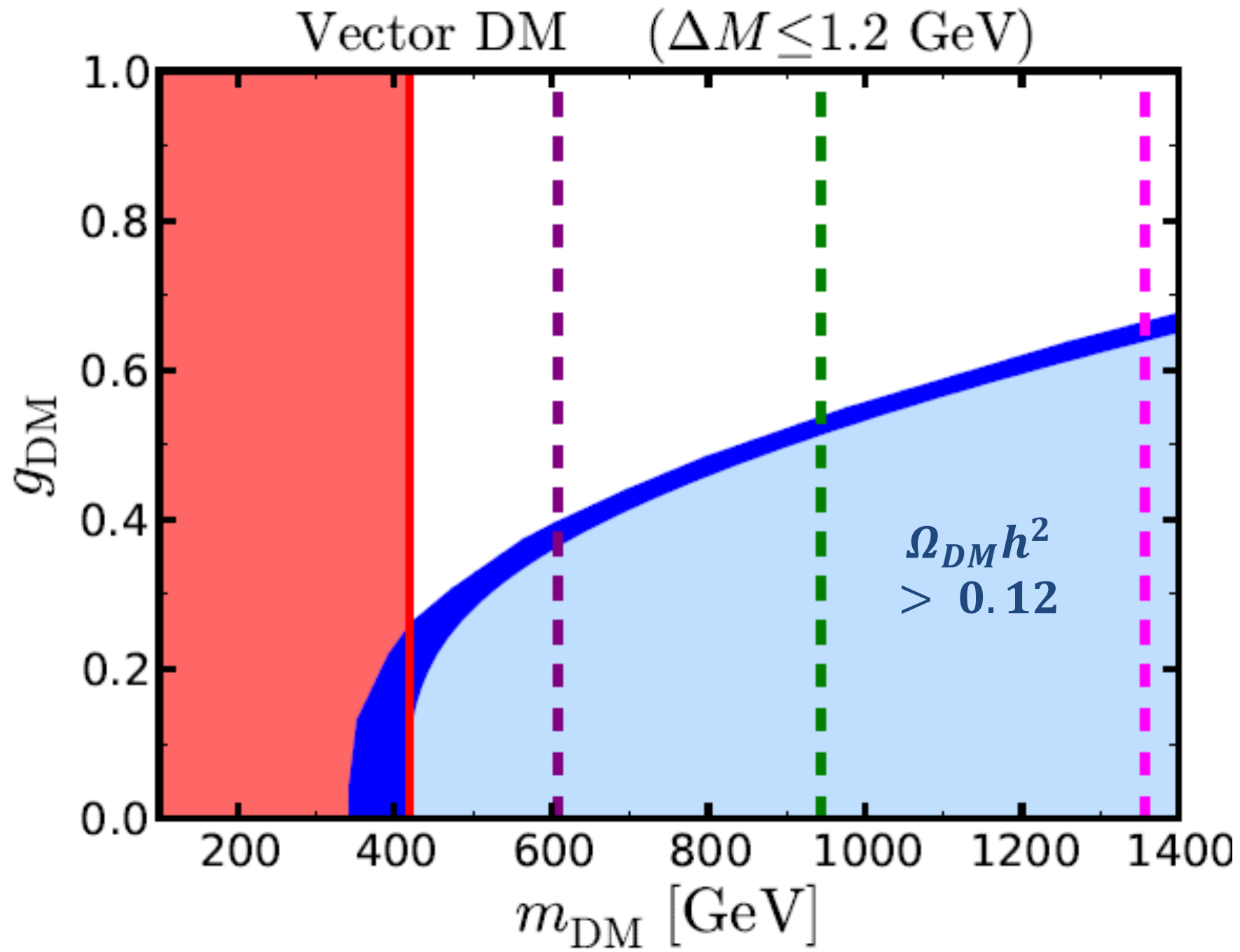
V_μ Ψ

$$V^\mu \bar{\Psi} \gamma_\mu P_R \tau \subset \mathcal{L}$$

NOT gauge-invariant, requires UV-completion, e.g. Extra-Dimensions



Model 3



13 TeV
300 fb⁻¹

High Lumi
13 TeV
3000 fb⁻¹

8 TeV
18.8 fb⁻¹

13 TeV
30 fb⁻¹

Conclusions

- Presented a set of simplified models which are distinct from and complimentary to the standard mediator-based simplified DM models
 - 1a - Majorana DM & Complex scalar CAP:** Gauge-invariant and renormalizable
 - 1b - Majorana DM L-R mixing & Complex scalar CAP:** UV theory can include SUSY
 - 2 - Scalar DM & Dirac Fermion CAP:** Gauge-invariant and renormalizable
 - 3 - Vector DM & Dirac Fermion CAP:** UV theory can be a Kaluza-Klein model in extra dimensions
- The signal we study for collider searches is the pair-production of CAPs that then ultimately decay into cosmologically stable dark matter
- Collider strategy: long-lived charged particles; distinct from jets + MET
- Collider searches here are well ahead of DD (no tree-level ints. with quarks) and ID (velocity suppression)
- The only limitation for successful LHC searches is the requirement of $\sim 1\text{GeV}$ compressed mass spectrum
- Co-annihilation partner plays an important role in lowering the relic density

- We presented a set of simplified models which are distinct from and complimentary to the standard mediator-based simplified DM models set:

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Scalar DM

Model-2			
Component	Field	Charge	Interaction (5.6)
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Vector DM

Model-3			
Component	Field	Charge	Interaction (5.7)
DM	Vector (V_μ)	$Y = 0$	$V_\mu(\bar{\Psi}\gamma^\mu P_R\tau) + \text{h.c.}$
CAP	Dirac fermion (Ψ)	$Y = -1$	

- These simplified modes can be used by ATLAS & CMS to interpret searches for long-lived charged particles to explore the new range of DM scenarios

Back-up slides

To draw simple collider reach projections (used above) for higher Energy and Luminosity:

- For a given theory model, e.g. Model 1a, compute the signal cross-section as a function of m_{DM} (for a given Energy)
- Plot the predicted number of events $= L \times \text{cross-section}$
- Plot the experimental exclusion curve $= \text{max number of allowed events (at a given confidence level)}$ as a function of m_{DM}
- The intersection of the theory and experimental excl. curves at m_{DM}^* gives the excluded region $m_{DM} < m_{DM}^*$
- Assume that the experimental exclusion curve [excluded number of events] remains the same at higher E and higher L
- Theory curve is recomputed and the new bound found $m_{DM} < m_{DM}^*(E, L)$

...this follows the 'Collider Reach' philosophy...

The same 2-to-2 cartoon involving DM and SM particles describes the DD, ID and Collider searches (only the initial state varies).

[One has to be extra careful with drawing conclusions from this in abridged theories - Simplified Models. The choice of the Dark sec & the SM particles involved in the dominant process in each case can be completely different.]

