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- $Au = \lambda u$   
 $Av = \mu v$

 $[A, B] = 0$

a)  $\lambda \neq \mu$  (non-deg.)

$$Bu \stackrel{?}{=} v u$$

$$A(Bu) = B Au = B(\lambda u) = \lambda (Bu)$$

$\Rightarrow Bu$  is eigenvector of  $A$  w/ e.v.  $\lambda$

$\lambda$  non-deg.  $\Rightarrow u$  is only e.v. (direction) with  $\lambda$

$$\Rightarrow Bu = \alpha u \quad (\alpha \in \mathbb{R})$$

But then  $u$  is e.v. of  $B$  w/ e.v.  $\alpha$ !

|| by for  $v$ .

b)  $\lambda = \mu$ .

$\leadsto$  don't need  $Bu \parallel u$ . spans space:  $(u, v)$

Find eigenvector,  $|\psi\rangle$ , of  $B$  (with eigenvalue  $b$ )

$\Rightarrow$  consider  $|\psi\rangle = \alpha |u\rangle + \beta |v\rangle$  ↓ diff.  $\alpha$

(what are  $\alpha, \beta$ ?)

$$\begin{aligned} A|\psi\rangle &= \alpha A|u\rangle + \beta A|v\rangle \\ &= \alpha \lambda |u\rangle + \beta \lambda |v\rangle \end{aligned}$$

recall:  $A|u\rangle = \lambda |u\rangle$

$A|v\rangle = \lambda |v\rangle$

↑  $\lambda = \mu$

$$= \alpha \lambda |u\rangle + \beta \lambda |v\rangle$$

$$= \lambda (\alpha |u\rangle + \beta |v\rangle)$$

$$= \lambda |\psi\rangle$$

$$B|\psi\rangle \stackrel{?}{=} b|\psi\rangle$$

If true:

$$B|\psi\rangle = B(\alpha |u\rangle + \beta |v\rangle) = \underbrace{\alpha B|u\rangle + \beta B|v\rangle}_{\text{matrix repr. of } B} = \underbrace{b(\alpha |u\rangle + \beta |v\rangle)}_{\text{vector repr. of } |\psi\rangle} = \alpha b |u\rangle + \beta b |v\rangle$$

side:

$$(\langle u|u\rangle = 1, \langle u|v\rangle = 0, \langle v|v\rangle = 1)$$

$$\bullet \times \langle u| : \quad \underbrace{\alpha \langle u|B|u\rangle}_{B_{uu}} + \underbrace{\beta \langle u|B|v\rangle}_{B_{uv}} = \alpha b$$

$$\bullet \times \langle v| : \quad \underbrace{\alpha \langle v|B|u\rangle}_{B_{vu}} + \underbrace{\beta \langle v|B|v\rangle}_{B_{vv}} = \beta b$$

$$\Rightarrow \underbrace{\begin{pmatrix} B_{uu} & B_{uv} \\ B_{vu} & B_{vv} \end{pmatrix}}_{\text{matrix repr. of } B \text{ in } (u,v) \text{ space}} \underbrace{\begin{pmatrix} \alpha \\ \beta \end{pmatrix}}_{\text{vector repr. of } |\psi\rangle = \alpha|u\rangle + \beta|v\rangle} = b \underbrace{\begin{pmatrix} \alpha \\ \beta \end{pmatrix}}_{\text{vector repr. of } |\psi\rangle = \alpha|u\rangle + \beta|v\rangle}$$

Solve e.va, e.vec of  $B \Rightarrow$  get  $|\psi\rangle$  (get  $\alpha, \beta$ ).

$$\langle \Phi | H_0 | \Phi \rangle = ?$$

$$H_0 = h(r_1, p_1) + h(r_2, p_2) \quad ; \quad h(r, p) = -\frac{\hbar^2}{2m} \nabla^2 + V(r)$$

$$\begin{aligned} \Phi &= \frac{1}{\sqrt{2}} \begin{vmatrix} \psi_1(r_1) \chi_1(1) & \psi_1(r_2) \chi_1(2) \\ \psi_2(r_1) \chi_2(1) & \psi_2(r_2) \chi_2(2) \end{vmatrix} \\ &= \frac{1}{\sqrt{2}} \left[ \psi_1(r_1) \psi_2(r_2) \chi_1(1) \chi_2(2) \right. \\ &\quad \left. - \psi_1(r_2) \psi_2(r_1) \chi_1(2) \chi_2(1) \right] \end{aligned}$$

$$\text{so } \langle \Phi | h(r_1, p_1) | \Phi \rangle$$

$$\begin{aligned} &= \frac{1}{2} \iint d\mathbf{r}_1 d\mathbf{r}_2 \left\{ \left[ \psi_1^*(r_1) \psi_2^*(r_2) \chi_1^\dagger(1) \chi_2^\dagger(2) \right. \right. \\ &\quad \left. \left. - \psi_1^*(r_2) \psi_2^*(r_1) \chi_1^\dagger(2) \chi_2^\dagger(1) \right] \right. \\ &\quad \times h(r_1, p_1) \\ &\quad \times \left[ \psi_1(r_1) \psi_2(r_2) \chi_1(1) \chi_2(2) \right. \\ &\quad \left. - \psi_1(r_2) \psi_2(r_1) \chi_1(2) \chi_2(1) \right] \left. \right\} \end{aligned}$$

check notes

$$= \frac{1}{2} \left\{ \int d^3 r_1 \psi_1^*(r_1) h(r_1, p_1) \psi_1(r_1) \underbrace{\int d^3 r_2 \psi_2^*(r_2) \psi_2(r_2)}_{=1 \text{ (norm.)}} \underbrace{\chi_1^\dagger(1) \chi_1(1)}_{=1} \underbrace{\chi_2^\dagger(2) \chi_2(2)}_{=1} \right.$$

$$- \int d^3 r_1 \psi_1^*(r_1) h(r_1, p_1) \psi_2(r_1) \underbrace{\int d^3 r_2 \psi_2^*(r_2) \psi_1(r_2)}_{=0 \text{ (orthog.)}} \chi_1^\dagger(1) \chi_1(2) \chi_2^\dagger(2) \chi_2(1)$$

$$- \int d^3 r_1 \psi_2^*(r_1) h(r_1, p_1) \psi_1(r_1) \underbrace{\int d^3 r_2 \psi_1^*(r_2) \psi_2(r_2)}_{=0 \text{ (orthog.)}} \chi_1^\dagger(2) \chi_1(1) \chi_2^\dagger(1) \chi_2(2)$$

$$+ \int d^3 r_1 \psi_2^*(r_1) h(r_1, p_1) \psi_2(r_1) \underbrace{\int d^3 r_2 \psi_1^*(r_2) \psi_1(r_2)}_{=1 \text{ (norm.)}} \underbrace{\chi_1^\dagger(2) \chi_1(2)}_{=1} \underbrace{\chi_2^\dagger(1) \chi_2(1)}_{=1} \left. \right\}$$

$$= \frac{1}{2} \left\{ \int d^3 r_1 \psi_1^*(r_1) h(r_1, p_1) \psi_1(r_1) + \int d^3 r_1 \psi_2^*(r_1) h(r_1, p_1) \psi_2(r_1) \right\}$$

$$= \frac{1}{2} \sum_{i=1}^2 \int d^3 r_i \psi_i^*(\underline{r}_i) h(\underline{r}_i, p_i) \psi_i(\underline{r}_i)$$

$$= \frac{1}{2} \sum_i \left\{ \int d^3 r_i \psi_i^*(\underline{r}_i) \left[ -\frac{\hbar^2}{2m} \nabla^2 \psi_i(\underline{r}_i) \right] + \int d^3 r_i |\psi_i(\underline{r}_i)|^2 V(\underline{r}_i) \right\}$$

but  $r_i$  is just a dummy variable:

$$\langle \Phi | h(\underline{r}_1, p_1) | \Phi \rangle$$

$$= \frac{1}{2} \sum_i \left\{ \int d^3 r \psi_i^*(\underline{r}) \left[ -\frac{\hbar^2}{2m} \nabla^2 \psi_i(\underline{r}) \right] + \int d^3 r \rho(\underline{r}) V(\underline{r}) \right\}$$

$$\text{so } \langle \Phi | H_0 | \Phi \rangle = \langle \Phi | (h(\underline{r}_1, p_1) + h(\underline{r}_2, p_2)) | \Phi \rangle$$

$$= \frac{1}{2} \sum_i \left\{ \int d^3 r \psi_i^*(\underline{r}) \left[ -\frac{\hbar^2}{2m} \nabla^2 \psi_i(\underline{r}) \right] + \int d^3 r \rho(\underline{r}) V(\underline{r}) \right\}$$

$$+ \frac{1}{2} \sum_i \left\{ \int d^3 r \psi_i^*(\underline{r}) \left[ -\frac{\hbar^2}{2m} \nabla^2 \psi_i(\underline{r}) \right] + \int d^3 r \rho(\underline{r}) V(\underline{r}) \right\}$$

$$= \sum_i \left\{ \int d^3 r \psi_i^*(\underline{r}) \left[ -\frac{\hbar^2}{2m} \nabla^2 \psi_i(\underline{r}) \right] + \int d^3 r \rho(\underline{r}) V(\underline{r}) \right\}$$

□

a) Yes, since  $\phi_A, \phi_B$  are, and there's a  $\frac{1}{\sqrt{2}}$  in  $\psi^{\text{nd}}$ . Explicitly:

$$\begin{aligned} \langle \psi_{AB}^{\text{dis}} | \psi_{AB}^{\text{dis}} \rangle &= \int d^2r |\psi_{AB}^{\text{dis}}(\underline{r}, \underline{r}')|^2 \\ &= \iint d\mathbf{r} d\mathbf{r}' |\phi_A(\mathbf{r}) \phi_B(\mathbf{r}')|^2 \\ &= \int d\mathbf{r} |\phi_A(\mathbf{r})|^2 \int d\mathbf{r}' |\phi_B(\mathbf{r}')|^2 \\ &= 1 \cdot 1 = 1 \quad \checkmark \text{yes.} \end{aligned}$$

$$\langle \psi_{AB}^{\text{nd}} | \psi_{AB}^{\text{nd}} \rangle = \int d^2r |\psi_{AB}^{\text{nd}}(\underline{r}, \underline{r}')|^2 \quad \langle \psi | \psi \rangle = \int d\mathbf{r} |\psi(\mathbf{r})|^2$$

$$= \frac{1}{2} \int d^2r |\phi_A(\mathbf{r}) \phi_B(\mathbf{r}') \pm \phi_B(\mathbf{r}) \phi_A(\mathbf{r}')|^2$$

$$= \frac{1}{2} \int d^2r \left( \phi_A(\mathbf{r}) \phi_B(\mathbf{r}') \pm \phi_B(\mathbf{r}) \phi_A(\mathbf{r}') \right)^* \left( \phi_A(\mathbf{r}) \phi_B(\mathbf{r}') \pm \phi_B(\mathbf{r}) \phi_A(\mathbf{r}') \right)$$

$$\begin{aligned} &= \frac{1}{2} \int d^2r \left\{ \underbrace{|\phi_A(\mathbf{r})|^2 |\phi_B(\mathbf{r}')|^2}_{\int d\mathbf{r} |\phi_A(\mathbf{r})|^2 \int d\mathbf{r}' |\phi_B(\mathbf{r}')|^2 = 1 \cdot 1 = 1} \pm \underbrace{\phi_A^*(\mathbf{r}) \phi_A(\mathbf{r}') \phi_B(\mathbf{r}) \phi_B^*(\mathbf{r}')}_{=0} \pm \underbrace{\phi_A^*(\mathbf{r}') \phi_A(\mathbf{r}) \phi_B^*(\mathbf{r}) \phi_B(\mathbf{r}')}_{\int d\mathbf{r} \phi_A(\mathbf{r}) \phi_B^*(\mathbf{r}) \int d\mathbf{r}' \phi_A^*(\mathbf{r}') \phi_B(\mathbf{r}')} \right. \\ &\quad \left. + \underbrace{|\phi_B(\mathbf{r})|^2 |\phi_A(\mathbf{r}')|^2}_{=1} \right\} \end{aligned}$$

= 0  
since wfs don't overlap.  
 $\Leftrightarrow$  orthogonal.

$$= \frac{1}{2} \{1 + 1\}$$

$$= 1$$

$$b) \quad \hat{O}_1 = O_1(\underline{r}) + O_1(\underline{r}')$$

$$i) \quad \langle \psi_{AB}^{ind} | \hat{O}_1 | \psi_{AB}^{ind} \rangle$$

$$= \int d^2 \underline{r} \quad \psi_{AB}^{ind*}(\underline{r}, \underline{r}') \hat{O}_1(\underline{r}, \underline{r}') \psi_{AB}^{ind}(\underline{r}, \underline{r}')$$

$$= \frac{1}{2} \int d^2 \underline{r} \left( \phi_A(\underline{r}) \phi_B(\underline{r}') \pm \phi_B(\underline{r}) \phi_A(\underline{r}') \right)^* \left[ O_1(\underline{r}) + O_1(\underline{r}') \right] \left( \phi_A(\underline{r}) \phi_B(\underline{r}') \pm \phi_B(\underline{r}) \phi_A(\underline{r}') \right)$$

again, cross terms vanish due to orthogonality of  $\phi_A$  &  $\phi_B$ ,

$$= \frac{1}{2} \left[ \int d^2 \underline{r} \quad \phi_A^*(\underline{r}) \phi_B^*(\underline{r}') \left( O_1(\underline{r}) + O_1(\underline{r}') \right) \phi_A(\underline{r}) \phi_B(\underline{r}') \right. \\ \left. + \int d^2 \underline{r} \quad \phi_B^*(\underline{r}) \phi_A^*(\underline{r}') \left( O_1(\underline{r}) + O_1(\underline{r}') \right) \phi_B(\underline{r}) \phi_A(\underline{r}') \right]$$

and the integral goes to unity by normalisation:

$$= \frac{1}{2} \left[ \int d\underline{r} \quad \phi_A^*(\underline{r}) O_1(\underline{r}) \phi_A(\underline{r}) + \int d\underline{r}' \quad \phi_B^*(\underline{r}') O_1(\underline{r}') \phi_B(\underline{r}') \right. \\ \left. + \int d\underline{r}' \quad \phi_A^*(\underline{r}') O_1(\underline{r}') \phi_A(\underline{r}') + \int d\underline{r} \quad \phi_B^*(\underline{r}) O_1(\underline{r}) \phi_B(\underline{r}) \right]$$

$$= \sum_{I \in \{A, B\}} \int d\underline{r} \quad \phi_I^*(\underline{r}) O_1(\underline{r}) \phi_I(\underline{r})$$

$$ii) \quad \langle \psi_{AB}^{dis} | \hat{O}_1(\underline{r}, \underline{r}') | \psi_{AB}^{dis} \rangle$$

$$= \int \int d^3 \underline{r} \quad \phi_A^*(\underline{r}) \phi_B^*(\underline{r}') \left( O_1(\underline{r}) + O_1(\underline{r}') \right) \phi_A(\underline{r}) \phi_B(\underline{r}')$$

similarly,

$$= \int d\underline{r} \phi_A^*(\underline{r}) \mathcal{O}_1(\underline{r}) \phi_A(\underline{r}) + \int d\underline{r}' \phi_B^*(\underline{r}') \mathcal{O}_1(\underline{r}') \phi_B(\underline{r}')$$

ie the same expression.

$$c) \hat{\mathcal{O}}_2 = \mathcal{O}_2(|\underline{r} - \underline{r}'|)$$

$$i) \langle \psi_{AB}^{dis} | \hat{\mathcal{O}}_2 | \psi_{AB}^{dis} \rangle$$

$$= \int d^2 \underline{r} \phi_A^*(\underline{r}) \phi_B^*(\underline{r}') \mathcal{O}_2(|\underline{r} - \underline{r}'|) \phi_A(\underline{r}) \phi_B(\underline{r}')$$

$$= \int d^2 \underline{r} |\phi_A(\underline{r})|^2 |\phi_B(\underline{r}')|^2 \mathcal{O}_2(|\underline{r} - \underline{r}'|)$$

$$ii) \langle \psi_{AB}^{ind} | \hat{\mathcal{O}}_2 | \psi_{AB}^{ind} \rangle$$

$$= \frac{1}{2} \int d^2 \underline{r} \left( \phi_A(\underline{r}) \phi_B(\underline{r}') \pm \phi_B(\underline{r}) \phi_A(\underline{r}') \right)^* \mathcal{O}_2(|\underline{r} - \underline{r}'|) \left( \phi_A(\underline{r}) \phi_B(\underline{r}') \pm \phi_B(\underline{r}) \phi_A(\underline{r}') \right)$$

similarly: cross terms vanish leaving

$$= \int d^2 \underline{r} |\phi_A(\underline{r})|^2 |\phi_B(\underline{r}')|^2 \mathcal{O}_2(|\underline{r} - \underline{r}'|)$$

$$\text{so } \langle \psi_{AB}^{ind} | \hat{\mathcal{O}}_i | \psi_{AB}^{ind} \rangle = \langle \psi_{AB}^{dis} | \hat{\mathcal{O}}_i | \psi_{AB}^{dis} \rangle \text{ for } i \in \{1, 2\}.$$

↖ indistinguishable

ie. we have identical particles, so their wfs must respect exchange symmetry, but since they are localized (separately), the symmetry

(fermionic/bosonic) doesn't affect observable quantities (expectation values of operators) & can be treated as distinguishable.



↑ ... as convenient.

↑  
which means simpler expressions.

# problem-four

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a) H atom: 1 electron :  $\Psi(x, y, z) = \Psi(r)$

$10^3$  points (coarse grid; cube of side length 10) [10 points per cartesian axis]

$\Rightarrow 10^3$  numbers to be stored (32 bit integers)

$\Rightarrow 4 \times 10^3$  bytes

$\Rightarrow \underline{\underline{4 \text{ kB}}}$

b) Li atom: 3 electrons :  $\Psi(\underline{r}_1, \underline{r}_2, \underline{r}_3) = \Psi(x_1, y_1, z_1, x_2, y_2, z_2, x_3, y_3, z_3)$

$\Rightarrow$  now we need  $10^9$  numbers

$\Rightarrow 4 \text{ GB}$  (fits on DVD  $\checkmark$ )

c) atom with  $N$  electron requires:  $(10^3)^N$  numbers  $\Rightarrow 4 \times 10^{3N}$  bytes

$m_{\text{DVD}} \sim 1 \text{ g} = 10^{-3} \text{ kg}$

i)  $M_{\text{earth}} \sim 6 \times 10^{24} \text{ kg} = 6 \times 10^{27} \text{ DVDs}$

DVD stores  $\sim 10 \text{ GB} = 10 \times 10^9 \text{ bytes} = 10^{10} \text{ bytes}$

$\Rightarrow 6 \times 10^{37} \text{ bytes}$

$$\frac{4 \times 10^{3N}}{6 \times 10^{37}} > 1$$

$$\Rightarrow \frac{2}{3} \times 10^{3N-37} \stackrel{\text{crit}}{=} 1$$

$$10^{3N-37} = 1.5$$

$$3N - 37 = \log_{10} 1.5$$

$$N = \frac{37 + \log 1.5}{3} = 12.4 \quad (3sf)$$

$$\Rightarrow N=13 \text{ exceeds } m_{earth} \quad (Al)$$

$$ii) \quad m_{\odot} \sim 2 \times 10^{30} \text{ kg} = 2 \times 10^{33} \text{ DVDs} = 2 \times 10^{43} \text{ bytes}$$

$$\text{so } N > \frac{43 + \log 1.5}{3} = 14.4$$

$$\Rightarrow N=15 \quad (P)$$

$$iii) \quad m_{MW} \sim 5 \times 10^{11} m_{\odot} = 10^{42} \text{ kg} = 10^{45} \text{ DVDs} = 10^{55} \text{ bytes}$$

$$N > \frac{55 + \log 1.5}{3} = 18.4$$

$$\Rightarrow N=19 \quad (K)$$