

Drell-Yan in SHERPA

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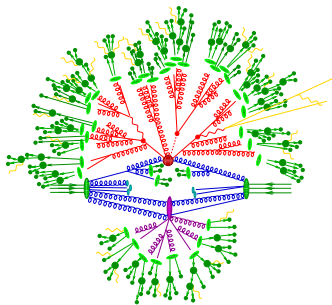
THE
ROYAL
SOCIETY

SHERPA

- latest release SHERPA-2.2.10
- includes
 - two matrix element generators
 - two parton showers
 - soft-photon resummation
 - multi-parton interactions
 - hadronisation
 - hadron decays
- typical setup for $V + \text{jets}$

$$pp \rightarrow \ell^+ \ell^- + 0, 1, 2j @ \text{NLO QCD} + \text{EW}_{\text{approx}} + 3, 4j @ \text{LO}$$

including resummed final-state QED corrections in the YFS
soft-photon resummation scheme



EW corrections

Kallweit, Lindert, Maierhöfer, Pozzorini, MS arXiv:1511.08692

Approximate EW corrections / EW_{virt} approximation

- incorporate approximate electroweak corrections in SHERPA's NLO QCD multijet merging (MEPS@NLO)
- tailored to large- p_T regions where EW corrections dominated by virtual W/Z exchange and RG running
- modify MC@NLO \bar{B} -function to include NLO EW virtual corrections

optionally include subleading Born

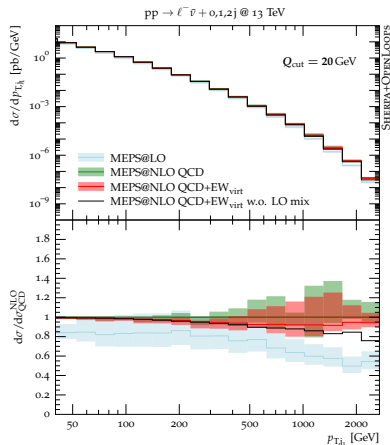
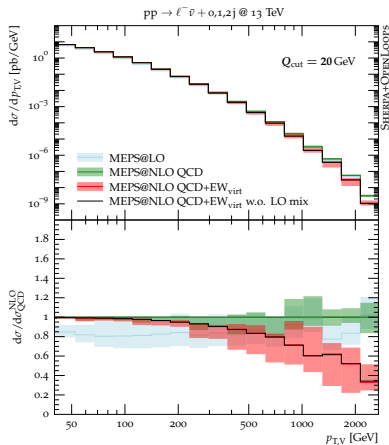
$$\bar{B}_{n,\text{QCD}+\text{EW}_{\text{virt}}}(\Phi_n) = \bar{B}_{n,\text{QCD}}(\Phi_n) + V_{n,\text{EW}}(\Phi_n) + I_{n,\text{EW}}(\Phi_n) + B_{n,\text{sub}}(\Phi_n)$$

exact virtual contribution

approximate integrated real contribution

- real QED radiation can be recovered through YFS resummation
- simple stand-in for proper QCD+EW matching and merging

EW corrections



⇒ particle level events including dominant EW corrections

EW corrections

- approx. EW corrections in EW_{virt} scheme available as alternative event weights (same as scale and PDF variations)
 - separately for EW, EW+subLO, EW+subLO+subLO2, ...
- calculated in three schemes
 - additive

$$\overline{B}_n^{\text{QCD+EW}} = \overline{B}_n^{\text{QCD}} + B_n \delta_{\text{EW}}^{\text{approx}} + B_n^{\text{sub}}$$

- multiplicative

$$\overline{B}_n^{\text{QCD} \times \text{EW}} = \overline{B}_n^{\text{QCD}} (1 + \delta_{\text{EW}}^{\text{approx}}) + B_n^{\text{sub}}$$

- exponentiated (caution: only Sudakov logs exponentiate, not complete EW_{virt})

$$\overline{B}_n^{\text{QCD} \times \text{EW exp}} = \overline{B}_n^{\text{QCD}} \exp(\delta_{\text{EW}}^{\text{approx}}) + B_n^{\text{sub}}$$



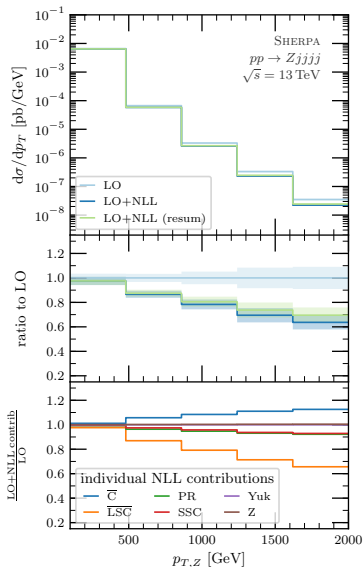
not the same as fixed-order multiplicative scheme



EW corrections

Upcoming in SHERPA-3.0.0

- fixed-order NLO EW
[MS arXiv:1712.07975](#)
- EW Sudakov approximation
for parton showered event
generation
[Bothmann, Napoletano
arXiv:2006.14635](#)



EW clustering

In the interpretation of a $Z + \text{jets}$ configuration allow for the possibility that this signature was produced by dijet configuration radiating a Z boson.

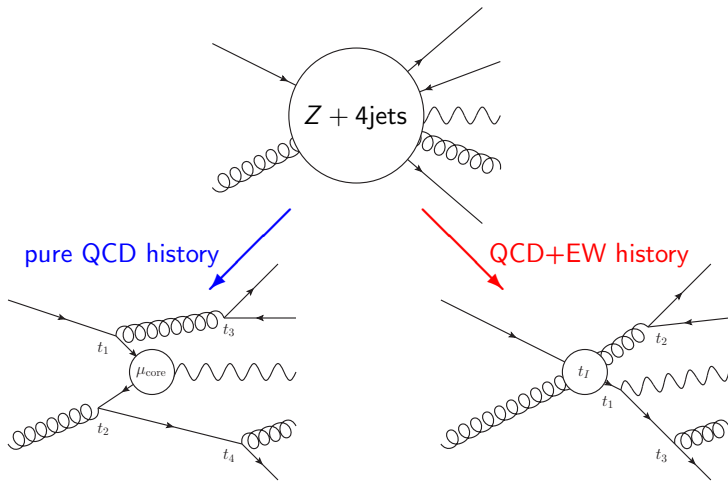
- identify parton shower history to determine CKKW scales and calculate intermediate Sudakov factors

$$\alpha_s^{n+k}(\mu_R^2) = \alpha_s^k(\mu_{\text{core}}^2) \prod_{i=1}^n \alpha_s(t_i)$$

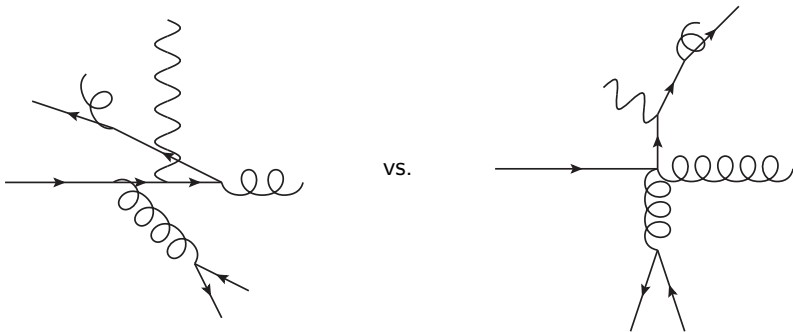
- identifying an EW splitting will “simply skip a step”, α taken as constant (as in ME), no EW Sudakov applied
- use spin averaged EW splitting functions in $E \gg m$ limit

This affects scale choices, no higher-order EW corrs. introduced.

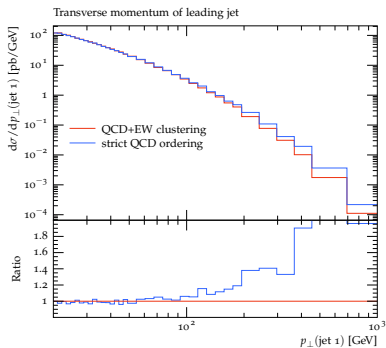
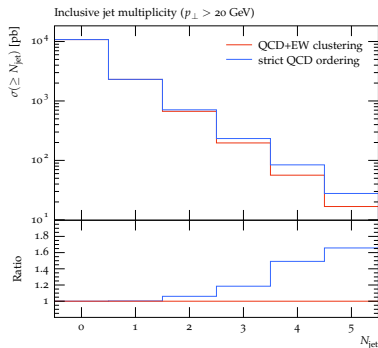
Multijet merging – identifying a history



QCD multijet merging – identifying a history



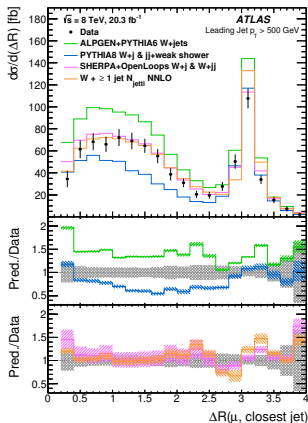
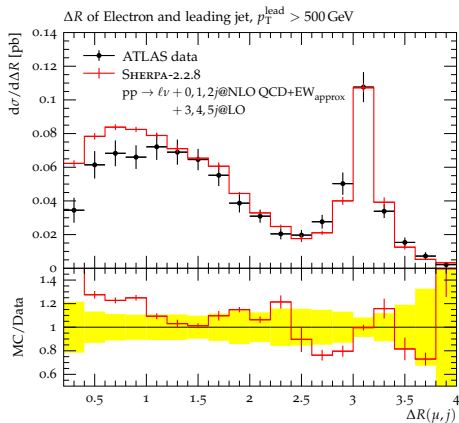
Importance of electroweak clustering



⇒ large impact at high p_{\perp} and multiplicity

EW clustering

ATLAS arXiv:1609.07045



Electroweak input schemes

For computing reasons, LHC collaborations need need $Z + \text{jets}$ samples that are useful (properly describe) **all** possible observables

- EW precision measurements ($A_i, \sin \theta, A_{\text{FB}}, \dots$)
- $p_{\text{T}}(\ell\ell)$
- underlying event in DY
- jet properties
- TeV-scale searches/measurements (DM backgrounds)
- V_{jj} properties (VBF-Z backgrounds)
- τ spin correlations in $Z \rightarrow \tau\tau$ (strongly dependent on $\sin \theta$)
- \vdots

Necessitates an input scheme that contains m_Z and $\sin \theta$ and is consistent at NLO EW.

Electroweak input schemes

Historically used inconsistent EW input scheme to have both m_Z and $\sin \theta$ at their measured values.

Fine at LO EW, impossible if higher-order EW effects are to be included.

G_μ and $\alpha(m_Z)$ schemes have m_Z as input, but value of $\sin \theta$ is off.

New input schemes with m_Z and $\sin \theta$ implemented in SHERPA-2.2.10.

21 input $(G_F, m_Z, \Gamma_Z, \sin \theta)$, derived $(\alpha, m_W, \Gamma_W, \text{vev})$

22 input $(\alpha(m_Z), m_Z, \Gamma_Z, \sin \theta)$, derived $(m_W, \Gamma_W, \text{vev})$

Schemes 21 and 22 also implemented by OpenLoops including appropriate renormalisation

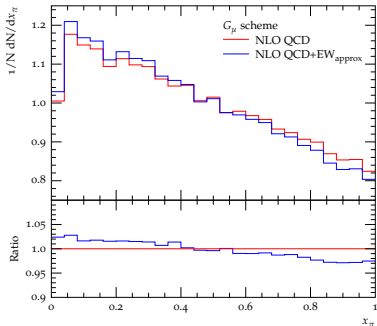
[Chiesa, Piccinini, Vicini arXiv:1906.11569](#)

⇒ they can be used to include approximate EW corrections in a sample

Impact of EW parameters – τ polarisation

In a consistent setup, having $\sin\theta$ as a dependent parameter may have unintended consequences, eg. τ spin correlations in $Z \rightarrow \tau\tau$ events.

Problem not apparent when using Too11 for ME, Too12 for PS, Too13 for τ decays as EW parameters typically not synchronised.



Using scheme with $\sin\theta$ as input fixes this issue.

QED corrections

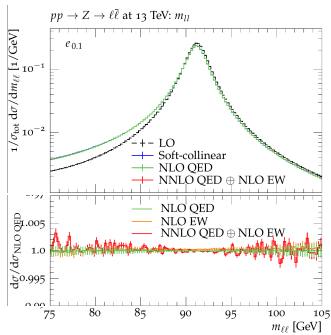
Sherpa employs a soft-photon resummation in the YFS scheme for higher-order QED corrections

up to SHERPA-2.2.10

- YFS \otimes NLO QED

Upcoming in SHERPA-3.0.0

- YFS \otimes NNLO QED + NLO EW



Krauss, Lindert, Linten, MS arXiv:1809.10650

Conclusions

- SHERPA-2.2.10 state-of-the-art all-round sample including
 - QCD multijet merging
 - approx. EW corrections tailored to TeV regime (no accuracy for incl. observables)
 - QCD+EW interpretation of multijet final states
 - parton shower QCD resummation
 - YFS soft-photon resummation w/ NLO QED corrections
 - suitable EW input schemes for DY precision observables
- currently most higher-order EW corrections included targeted to TeV scale physics, with the exception of soft-photon resummation
- many other features: elastic photon-induced procs. in EPA, ...
- additional features in upcoming SHERPA-3.0.0

<https://sherpa-team.gitlab.io>

Backup