Motivation	Status	Parton Showers	Future Experiments	Soft QCD	Developments
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# Precision QCD simulations

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#### April $6^{\rm th}$ 2023 – FNAL – Theory Seminar



Motivation	Status	Parton Showers	Future Experiments	Soft QCD	Developments
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- why precision tools?
- current precision
- better parton showers
- future experiments
- soft QCD as limiting factor
- what is achievable?

Motivation	Status	Parton Showers	Future Experiments	Soft QCD	Developments
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### why precision

#### (carrying coal to Newcastle)

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- to date no discovery of new physics (BSM)
- hope for "simple" discoveries is waning (don't expect
- push into precision tests of the Standard Model
- statistical uncertainties approach zero (because of fantastic work of accelerator, DAQ, etc.)
- systematic exp. uncertainties decrease
- theoretical uncertainties are or become dominant (obstacle to full explitation of LHC)



Motivation

(a pity, but that's Nature)

(find it or constrain "subtle"!)

No Match

100% Comparison Match

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(don't expect anything glaringly obvious)

(because of ingenious experimental work)

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Motivation

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Future Experiment

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#### how to build an event generator

- paradigm: "divide et impera"
- divide simulation in distinct phases, with (logarithmically) separated scales
- start with signal event

(fixed order perturbation theory)

• dress partons with parton shower

(resummed perturbatkon theory)

add underlying event

(phenomenological models)

hadronize partons

(phenomenological mod

decay hadrons

(effective theories, simple symmetries & data)



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#### current precision

(where we are)



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#### status: executive summary

 $\checkmark~$  NNLO  $\otimes$  parton shower for colour singlet production

(MINNLO: 1309.4634, 1407.2940, ..., 2208.12660; UNNLOPS: 1405.4607, 1407.3773)

 $\checkmark~$  NNLO  $\otimes$  parton shower for heavy quarks

(MINNLO: 2112.04168 (tt), 2302.01645 (bb))

✓ MEPS@NLO: NLO multijet merging

(SHERPA: 1207.5030; MADGRAPH: 1209.6215; PYTHIA: 1211.7278; HERWIG: 1705.06700 plus follow-ups & refinements)

 $\checkmark\,$  all of the above including EW@NLO

(explicit: 1511.08692, 1705.00598, ..., 2204.07652; Sudakov approximation: hep-ph/0010201, 2111.13453)

 $\checkmark$  (N)NLO  $\otimes$  N<sup>1,2,3</sup>LL $\otimes$  parton shower

(GENEVA: 1211.7049, 1508.01475, 2102.08390, ...)

(2107.01224, 2208.02276 (not covered here))

- multijet merging with TMDs
- improving parton showers

((next-to leading) logarithmic accuracy (see below); amplitude evolution: 1802.08531, ... (not covered here))



• relative size argument:  $\alpha_s^2 \approx \alpha_W$ : must include NLO EW corrections for  $\mathcal{O}(1 - 10\%)$  accuracy  $\implies$  automated in OPENLOOPS, RECOLA, aMC@NLO \_ MADGRAPH

(1705.00598, 1704.05783, 1405.0301)

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#### SM precision simulation in a nutshell: Drell-Yan

- current "accuracy standard(s)":
  - fixed-order: N<sup>3</sup>LO for inclusive, NNLO for Vj
  - matching: NNLOPS for inclusive V
  - merging: MEPs@NLO for  $V + \leq 2$  jets at NLO  $V + \geq 3$  jets at LO
- dominating QCD effects:  $\mathcal{O}(10-30\%)$ 
  - low- $p_{\perp}$  region dominated by parton shower
  - high- $p_{\perp}$  region dominated by (multi-) jet topologies
  - higher accuracy in rate (and some shapes) through NNLO matching
- must add EW corrections for %-level precision
  - EW correction at large scales  $\mathcal{O}(10\%)$
  - QED FSR + EW for V line shapes at  $\mathcal{O}(1\%)$

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# NNLOPS for Z production: MINNLO & UNNLOPS

#### (1407.2904, 1405.3607)

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- different logic of achieving NNLO precision
- available for H, V production (both) and VV production (MINNLO)

Status

# merging example: $p_{\perp,\gamma\gamma}$ in MEPS@LO vs. NNLO

#### (arXiv:1211.1913 [hep-ex])



#### 

### better parton showers?

(the story never gets old)



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#### parton showers, compact notation

Sudakov form factor (no-decay probability)

$$\Delta_{ij,k}^{(\mathcal{K})}(t,t_0) = \exp\left[-\int_{t_0}^t \frac{\mathrm{d}t}{t} \frac{\alpha_s}{2\pi} \int \mathrm{d}z \frac{\mathrm{d}\phi}{2\pi} - \underbrace{\mathcal{K}_{ij,k}(t,z,\phi)}_{\text{splitting kernel for}}\right]$$

• evolution parameter t defined by kinematics

generalised angle (HERWIG ++) or transverse momentum (PYTHIA, SHERPA)

• will replace  $\frac{\mathrm{d}t}{t}\mathrm{d}z\frac{\mathrm{d}\phi}{2\pi}\longrightarrow\mathrm{d}\Phi$ 

(subtle differences important for theoretical accuracy (LL vs. NLL)!)

- scale choice for strong coupling:  $\alpha_{s}(k_{\perp}^{2})$
- regularisation through cut-off t<sub>0</sub>

resums classes of higher logarithms

scale for onset of non-perturbative effects (hadronization)!

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#### factorisation of amplitudes: colour coherence

collinear:

$$\sum_{\lambda,\lambda'=\pm} {n < 1, \ldots, n \rangle_n} \xrightarrow{i \parallel j} \sum_{\lambda,\lambda'=\pm} {n-1 \langle 1, \ldots, \lambda(ij), \ldots, \lambda$$

with spin-dependent splitting function  $P_{(ij)i}^{\lambda\lambda'}(z)$ 

soft:

$${}_{n}\langle 1,\ldots,n|1,\ldots,n\rangle_{n} \xrightarrow{p_{j}\to 0} \\ -8\pi\alpha_{s}\sum_{i,k\neq j} {}_{n-1}\langle 1,\ldots,j,\ldots,n|\mathsf{T}_{i}\mathsf{T}_{k}\mathsf{w}_{ik,j}|1,\ldots,j,\ldots,n\rangle_{n-1}$$

with colour-insertion operators  $\mathbf{T}_{i,k}$  & soft eikonal

$$w_{ik,j} = \frac{p_i p_k}{(p_i p_j)(p_j p_k)} = \frac{W_{ik,j}}{E_j^2} = \frac{1}{E_j^2} \frac{1 - \cos \theta_{ik}}{(1 - \cos \theta_{ij})(1 - \cos \theta_{jk})}$$

(obviously, frame-dependent when expressed by energies & angles)

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#### soft eikonals, decomposed

• textbook decomposition (pink bible):  $W_{ik,j} = \tilde{W}^i_{ik,j} + \tilde{W}^k_{ki,j}$ with "radiator functions"  $\tilde{W}^i_{ik,j}$ : (identify "splitters" to combine with collinear terms)

$$\tilde{W}^i_{ik,j} = \frac{1}{2} \left( \frac{1 - \cos \theta_{ik}}{(1 - \cos \theta_{ij})(1 - \cos \theta_{jk})} + \frac{1}{1 - \cos \theta_{ij}} - \frac{1}{1 - \cos \theta_{jk}} \right)$$

• express  $\theta_{jk}$  for use in *i*-splitter term:

$$\cos\theta_{jk} = \cos\theta_{ij}\cos\theta_{ik} + \sin\theta_{ij}\sin\theta_{ik}\cos\phi_{jk}^{i}\ldots$$

• ... and average over azimuth  $\phi_{ik}^{i}$ :

$$\frac{1}{2\pi} \int_0^{2\pi} \mathrm{d}\phi^i_{jk} \tilde{W}^i_{ik,j} = \frac{\tilde{l}^i_{ik,j}}{1 - \cos\theta^i_j} , \qquad \text{where} \qquad \tilde{l}^i_{ik,j} = \left\{ \begin{array}{cc} 1 & \quad \text{if} \quad \theta^i_j < \theta^i_k \\ 0 & \quad \text{else} \end{array} \right.$$

(this is the well-known source of angular ordering)



• azimuthally integrated radiator function (normalised to  $2\pi$ ):



• need to include azimuth modulation, if observables sensitive to it

• but: naive inclusion bound to fail (MC efficiency  $\rightarrow$  0)

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### a new approach: PANSCALES

(resolving problems)



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Motivation	Status	Parton Showers	Future Experiments	Soft QCD	Developments
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comparison with fixed-order results

(PANSCALES: 1805.09327, 2205.02237)



- comparison with fixed-order results
- logarithmic accuracy in FS showering

(PANSCALES: 1805.09327, 2205.02237)

(PANSCALES: 2002.11114, 2011.10054)

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- comparison with fixed-order results
- logarithmic accuracy in FS showering
- logarithmic accuracy in IS showering

(PANSCALES: 1805.09327, 2205.02237)

(PANSCALES: 2002.11114, 2011.10054)

(PANSCALES: 2207.09467)



- comparison with fixed-order results
- logarithmic accuracy in FS showering
- logarithmic accuracy in IS showering
- spin correlations

(NPB 310 (1988) 571, hep-ph/0110108, PANSCALES: 2103.16526, 2111.01161)



(PANSCALES: 1805.09327, 2205.02237)

(PANSCALES: 2207.09467)

(PANSCALES: 2002.11114, 2011.10054)

- comparison with fixed-order results
- logarithmic accuracy in FS showering
- logarithmic accuracy in IS showering

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(PANSCALES: 2002.11114, 2011.10054)

(PANSCALES: 2207.09467)

- spin correlations (NPB 310 (1988) 571, hep-ph/0110108, PANSCALES: 2103.16526, 2111.01161)
- $\checkmark$  PANSCALES encodes massless parton showers at NLL accuracy

 $(1^{st} \text{ matching in } e^+e^-: 2301.09645)$ 

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### a new approach: ALARIC

(progress in SHERPA)



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### soft eikonals, decomposed again

#### • define positive definite radiators:

(borrowing from Catani & Seymour, Nucl. Phys. B485 (1997) 291)

$$ar{W}^i_{ik,j} = rac{1-\cos heta_{ik}}{(1-\cos heta_{ij})(2-\cos heta_{ij}-\cos heta_{jk})}$$

• same result after azimuth averaging, but  $\tilde{I}^{i}_{ik,j} \longrightarrow \bar{I}^{i}_{ik,j}$  with

$$ar{I}^{i}_{ik,j} = rac{1}{\sqrt{(ar{A}^{i}_{ij,k})^2 - (ar{B}^{i}_{ij,k})^2}}$$

where

$$\bar{A}_{ij,k}^{i} = \frac{2 - \cos \theta_{j}^{i}(1 + \cos \theta_{k}^{i})}{1 - \cos \theta_{k}^{i}} , \bar{B}_{ij,k}^{i} = \frac{\sqrt{(1 - \cos^{2} \theta_{j}^{i})(1 - \cos^{2} \theta_{k}^{i})}}{1 - \cos \theta_{k}^{i}}$$



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(2208.06057)

#### F. Krauss Precision QCD simulations



#### matching with collinear terms

• collinear limit of eikonal factors:

$$w_{ik,j} \xrightarrow{i||j} w_{ik,j}^{(\text{coll})}(z) = \frac{1}{2p_i p_j} \frac{2z}{1-z} , \quad \text{where} \quad z \xrightarrow{i||j} \frac{E_i}{E_i + E_j}$$

• compare with leading (1 - z)-terms of splitting functions

 $(1/z \text{ term in } g \rightarrow gg \text{ captured with other "dipole"})$ 

$$P_{qq}(z) = C_F \left(\frac{2z}{1-z} + (1-z)\right) ,$$
  

$$P_{gg}(z) = C_A \left(\frac{2z}{1-z} + z(1-z)\right) ,$$
  

$$P_{gq}(z) = T_R \left(1 - 2z(1-z)\right) .$$

 $\longrightarrow$  defines "collinear remnant"



#### kinematics: birds-eye view

- kinematics as main obstacle to NLL accuracy in dipole showers: recoil of subsequent soft emissions may change "NLL history"
- construct new mapping  $\{ \widetilde{p}_l \} \longrightarrow \{ p_l \}$

(inspired by Catani & Seymour's treatment of identified hadrons)

• logic: disentangle colour spectator  $\tilde{p}_k$  and recoil partner  $\tilde{K}$ 

(i.e. define a global recoil scheme, use spectator for eikonal/azimuth)



basis to prove NLL accuracy analytically

#### set-up of numerical tests

- compare results in  $\alpha_S \rightarrow 0$  limit with NLL result
- set-up for checks
  - fixed  $\alpha_s$
  - leading colour  $C_A = 2C_F = 3$
  - all partons massless
- example: azimuth angle between two leading Lund-plane declusterings

(should be  $\Delta \Psi_{12} = 0$ )



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Motivation	Status	Parton Showers	Future Experiments	Soft QCD	Developments
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#### numerics: event shapes

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- NLL

Alaric

Dire

 $a_r = 0.005$ 

 $a_s = 0.0025$ 

 $\lambda = \alpha_s \ln B_T$ 

 $a_{e} = 0.01$ 

-0.5 -0.4

 $\alpha_{-} = 0.00$ 







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Motivation	Status	Parton Showers	Future Experiments	Soft QCD	Developments
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### how about data?

#### (always nice to see practical impact, innit?)

#### set-up of data comparison

- compare hadron-level results with LEP data
- perturbative set-up
  - no higher orders (no matching or merging)
  - running two-loop  $\alpha_s$  with  $\alpha_s(M_z) = 0.118$
  - use CMW scheme for soft eikonal parts
  - all partons massless, masses emulated through simplistic thresholds
  - leading colour  $C_A = N_c = 3$ ,  $C_F = \frac{N_c^2 1}{2N_c}$
- non-perturbative set-up
  - need to use PYTHIA hadronization

(ALARIC not yet ready for heavy hadron decays)

• default parameters of PYTHIA 6.4, but

PARJ(21) = 0.3, PARJ(41) = 0.4, PARJ(42) = 0.36(ALARIC)/0.45(DIRE)

Future Experiment

Soft QCD 000000 Developments 00000000

#### differential jet rates & event shapes



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Motivation	Status	Parton Showers	Future Experiments	Soft QCD	Developments
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# MC4EIC?

#### (preparing for the future)

### precision QCD studies at HERA

- "old" HERA data and analysis as boot-camp for EIC
- HERA = unique test-bed for (non-)perturbative QCD:
  - $\bullet~{\sf large-}\,Q^2$  DIS has no MPI  $\longrightarrow$  initial state showering "clean"
  - large- $Q^2$  DIS has no MPI  $\longrightarrow$  beam fragmentation "clean"
  - add HERA data to hadronization tunes?
- also: large photo-production cross section:
  - test hadronic structure of photon (relevant for EIC)
  - nota bene: last fits of photon-PDF are 20 years old
  - new fits urgently needed for EIC

(that is, if we want to treat collinear factorisation as limiting case for TMD's etc..)

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H1

 $ln(\lambda_{i}^{1})$ 

#### precision QCD studies at HERA

• incidentally, recent paper by H1

2303.13620

uses modern MC's (HERWIG 7, PYTHIA 8, SHERPA)







#### precision QCD studies at HERA

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2303.13620

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#### resolved photon processes

- "resolved" photons (i.e. with QCD structure/PDF) at LEP & HERA
- fits date from early 2000's: can supplement with NLO MC machinery
- first steps (SHERPA) below



#### ZEUS, hep-ex/0112029

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#### resolved photon processes

- "resolved" photons (i.e. with QCD structure/PDF) at LEP & HERA
- fits date from early 2000's: can supplement with NLO MC machinery

first steps towards NLO (SHERPA) below



ZEUS, hep-ex/1205.6153

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Motivation	Status	Parton Showers	Future Experiments	Soft QCD	Developments
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### soft physics strikes back

(the ultimate frontier?)

# g ightarrow Q ar Q — a systematic nightmare

 parton showers geared towards collinear & soft emissions of gluons

(double log structure)

- g 
  ightarrow q ar q only collinear
- old measurements at LEP of g 
  ightarrow bb and  $g 
  ightarrow car{c}$  rate
- fix this at LHC for modern showers

(important for ttbb)

• questions: kernel, scale in  $\alpha_s$ 

(example:  $k_{\perp}$  vs.  $m_{bb}$ )



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# Motivation Status Parton Showers Future Experiments Soft QCD 000 000000 0000 0000 000000

#### latest ATLAS measurements

(arXiv:1812.09283, 1705.03374 [hep-ex])

- use *b*-tagged jets with R = 0.2 (left)
- use muons in  $B \to J/\Psi(\mu\mu) + X$  and  $B \to \mu + X$  as proxies (right)



#### massive quarks are tricky - encore

- heavy quarks also problematic in initial state: no PDF support for  $Q^2 \leq m_Q^2 \longrightarrow$  quarks stop showering
- possible solutions:
  - naive: ignore and leave for beam remnants (SHERPA)
  - better: enforce splitting in region around  $m_Q^2$  (PYTHIA)

 $\longrightarrow$  effectively produces collinear  ${\it Q}$  and gluon in IS

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• will need to check effect on precision obsevables:  $p_{\perp}^{(W)}/p_{\perp}^{(Z)}$ 

Future Experimen

Soft QCD 0000●0

### soft physics: strange strangeness

- universality of hadronization assumed
- parameters tuned to LEP data in particular: strangeness suppression
- for strangeness: flat ratios but data do not reproduce this
- looks like SU(3) restoration not observed for protons
- needs to be investigated (see next)





#### hadronization issues

(illustrative plots from arXiv:1610.09818 [hep-ph])

Image: A math a math

- "ridge" in *pp* collisions:
- HI-like behaviour: unexpected, unexplained
- play with hadronization?



Motivation	Status	Parton Showers	Future Experiments	Soft QCD	Developments
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### Developments & Challenges

(achievable goals in the next 5)

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### "predictions" for analytic tools

- expect progress in higher-order calculations
  - (semi-)automated NNLO $_{\rm QCD}$  for multi-particle FS & mixed exact NLO $_{\rm QCD}\otimes$  NLO $_{EW}$  for multi-particle FS

( I expect this to be driven by constructing efficient basis of master integrals and improving and automating IR subtraction methods)

 $\bullet\,$  approximate and/or exact  $N^3LO_{\rm QCD}$  beyond  $2\rightarrow 1$  topologies

(similar problems as before: master integrals, ...)

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- higher precision in PDF's
- don't know enough about resummation to have a clear picture, will concentrate on parton shower simulations instead

# simulation

- reminders/lessons from past decade:
  - NLO matching with parton showers has been trivial

(parton shower kernels equal NLO subtraction kernels)

 MEPS@LO and MEPS@NLO usually agree at 10% level, reduced scale uncertainties at NLO

(this is not yet fully explored when searching for efficient event generation)

- extrapolation to higher-orders in simulation:
  - NNLO will not be quite as simple will need  $\mathcal{O}(\alpha_s^2)$  kernels

(tricky, think about it as fully automated NNLO subtraction kernels)

- opens parton shower for systematic uncertainty analysis
  - $\longrightarrow$  expect main effect in further reduction of scale uncertainties
- non-trivial impact of choices (kinematics) on log accuracy

Motivation	Status	Parton Showers	Future Experiments	Soft QCD	Developments
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#### where does it stop?

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### theory limitations

"The great advances in science usually result from new tools rather than from new doctrines" (Freeman Dyson)

• technical challenges: speed & stability of fixed-order calculations

• numerical issues 1: phase-space integration efficiency

(current approach: multi-channel with process-specific mappings)

(hard to see how ML can make a massive difference - until now only "blanks")

• numerical issues 2: special functions in multi-loop master integrals

(convergence of series expansion, maybe need to go to quadruple precision)

• numerical issues 3: stability of numerical (N-1)-loop results

(example: NLO inputs to NNLO calculation, supremacy of compact analytic expressions)

 $\rightarrow$  this will necessitate highly technical, barely publishable work

(traditionally this is often a dead-end for careers in theory)

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### theory limitations

#### o physics challenges:

#### • physics challenge 1: sources of uncertainty

(simple scale variations by a factor of two may not be sufficient: central scale in multi-scale processes? in addition: new sources of uncertainty. example: kinematics scheme in parton showers and impact on log accuracy)

#### → community effort: agree on robust procedures

#### • physics challenge 2: complexity & control

(as the calculations become increasingly complex we need more and more independent – algorithmic, implementation, ... – checks. examples: two calculations for fixed order, two implementations of parton showers, etc.)

#### 

#### physics challenge 3: input parameters as source of uncertainty

(my favourite example: PDF's with different values for  $\alpha_S(M_Z)$  and, yes, there's more: PDF's etc.)

→ community effort: harvest "old data"

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### theory limitations

- "philosophical" challenges:
  - philosophical challenge 1: breakdown of factorisation

(studies suggest that at some high order factorisation breaks down for basic 2  $\rightarrow$  2 QCD processes it is not clear how this translates to more complicated processes)

 $\longrightarrow$  don't know if we should care

 philosophical challenge 2: impact of higher-twist (this is closely related to 1. simple back-of-the envelope: at DY m<sub>P</sub>/M<sub>7</sub> ≈ 1% ≈ α<sup>2</sup><sub>c</sub> - is this a problem?)

 $\longrightarrow$  this may need a careful analysis – is it linear or quadratic?

 philosophical challenge 3: impact of soft physics uncertainties (lots of them a a hadron collider: MPI/underlying event, hadronization, ...quite often models aiming to describe these effects are too similar, and based on identical data)

 $\longrightarrow$  only more data will help here

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