



Exercise 4

1. Consider the quantized Klein-Gordon field $\hat{\phi}(x)$.

(a) Show that the Pauli-Jordan function, defined via $i\Delta(x) = [\hat{\phi}(x), \hat{\phi}(0)]$, satisfies

$$\frac{\partial}{\partial t} \Delta(x)|_{x_0=0} = -\delta^3(\vec{x}).$$

(b) Define positive and negative frequent Δ -functions according to

$$i\Delta^\pm(x-y) = \pm \int \frac{d^4k}{(2\pi)^3} \theta(\pm k_0) \delta(k^2 - m^2) e^{-ik(x-y)}.$$

How can the Pauli-Jordan function and the Feynman propagator Δ_F be constructed from Δ^+ and Δ^- ?¹

(c) Show that Δ_F is a Greens function to the Klein-Gordon equation. Is this also true for Δ^+ and Δ^- ? How about the Pauli-Jordan function?

2. (a) Classically, the scattering angle for the scattering of a test particle off a spherically symmetric potential is a function of the projectile energy and the impact parameter only:

$$\vartheta(E_0, b) = \pi - 2b \int_{r_{\min}}^{\infty} \frac{dr}{r^2 [1 - V(r)/E_0 - b^2/r^2]^{1/2}}. \quad (1)$$

Starting from Eq. (1) calculate $\vartheta(E_0, b)$ for

- the scattering of two hard balls with radii R_1 and R_2 ,
- the case of Coulomb scattering.

Use the relation

$$\frac{d\sigma}{d\Omega} = \frac{b db}{d \cos \vartheta}$$

to derive the classical cross section for both processes.

(b) In terms of nonrelativistic quantum mechanics the differential cross section reads

$$\frac{d\sigma}{d\Omega} = |f(\vartheta)|^2 \quad \text{where} \quad f(\vartheta) = \frac{1}{k} \sum_l \frac{2l+1}{2i} (e^{2i\delta_l} - 1) P_l(\cos \vartheta). \quad (2)$$

Use Eq. (2) to prove the optical theorem²

$$\sigma = \frac{4\pi}{k} \text{Im}\{f(0)\}.$$

(turn over)

¹Hint: Use the Fourier representation of the Heaviside function

$$\theta(x) = \lim_{\varepsilon \rightarrow 0} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi i} \frac{e^{i\omega x}}{\omega - i\varepsilon}$$

²Hint: $\int_{-1}^{+1} dz P_l(z) P_{l'}(z) = \frac{2}{2l+1} \delta_{ll'}$ $P_l(\pm 1) = (\pm 1)^l$

3. Let $\hat{\phi}$ be an arbitrary quantum field. Show explicitly that the time ordered product of three field operators is given by

$$\begin{aligned}
 T[\hat{\phi}(x_A) \hat{\phi}(x_B) \hat{\phi}(x_C)] &= : \hat{\phi}(x_A) \hat{\phi}(x_B) \hat{\phi}(x_C) : \\
 &\quad + \langle 0 | T[\hat{\phi}(x_A), \hat{\phi}(x_B)] | 0 \rangle \hat{\phi}(x_C) \\
 &\quad + \langle 0 | T[\hat{\phi}(x_B), \hat{\phi}(x_C)] | 0 \rangle \hat{\phi}(x_A) \\
 &\quad + \epsilon_{AB} \langle 0 | T[\hat{\phi}(x_A), \hat{\phi}(x_C)] | 0 \rangle \hat{\phi}(x_B),
 \end{aligned}$$

where the sign of the permutation ϵ_{AB} is $+1$ for bosons and -1 for fermions.