

# $0\nu\beta\beta$ within the Left- Right Symmetric Model

Andres Olivares del Campo

2-11-2015 IPPP Student Seminar

Supervised by Prof. Allanach

# Outline

- Neutrino masses and seesaw mechanism

# Outline

- Neutrino masses and seesaw mechanism
- Experimental signatures of  $0\nu\beta\beta$

# Outline

- Neutrino masses and seesaw mechanism
- Experimental signatures of  $0\nu\beta\beta$
- Standard mechanism for  $0\nu\beta\beta$

# Outline

- Neutrino masses and seesaw mechanism
- Experimental signatures of  $0\nu\beta\beta$
- Standard mechanism for  $0\nu\beta\beta$
- Minimal Left-Right Symmetric Model (LRSM)

# Outline

- Neutrino masses and seesaw mechanism
- Experimental signatures of  $0\nu\beta\beta$
- Standard mechanism for  $0\nu\beta\beta$
- Minimal Left-Right Symmetric Model (LRSM)
- Extra contributions to  $0\nu\beta\beta$  and limits on new physics parameters

# Neutrino Masses

---

# Neutrino Masses

- Why do they have mass?

# Neutrino Masses

- Why do they have mass?

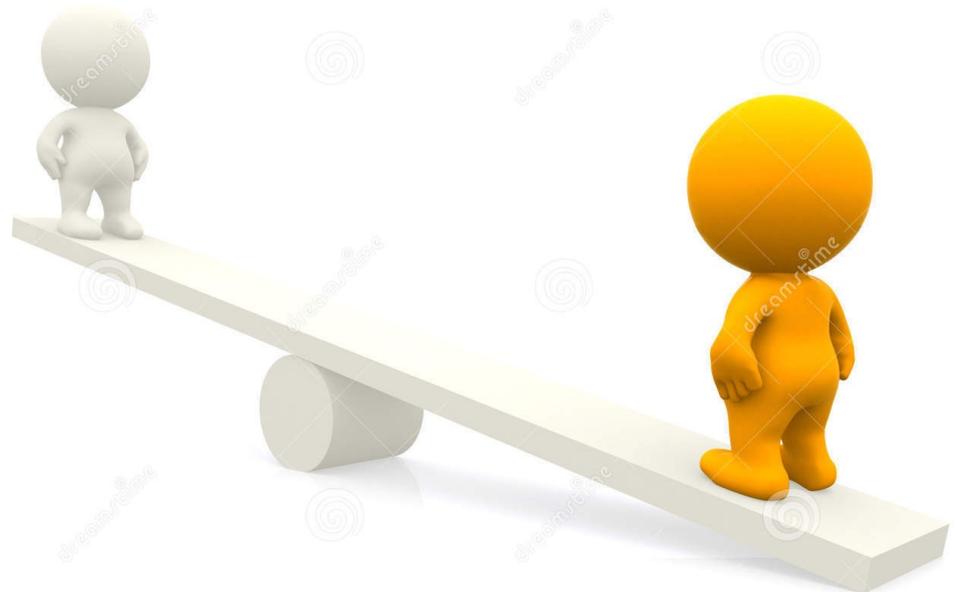
- Why small?

# Neutrino Masses

- Why do they have mass?

- Why small?

## Seesaw Mechanism



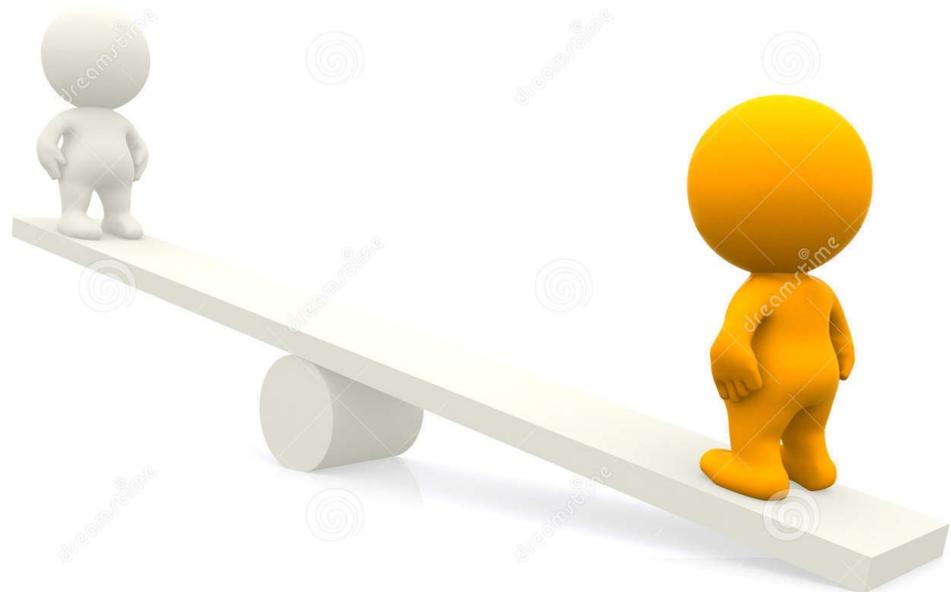
# Neutrino Masses

- Why do they have mass?

- Why small?

## Seesaw Mechanism

- Type I: Adding heavy right handed neutrinos  $N$
- Type II: Adding a scalar triplet  $\Delta$
- Type III: Adding a fermionic triplet



# Neutrino Masses

---

- Majorana or Dirac?

# Neutrino Masses

- Majorana or Dirac?
- Neutrinos are neutral so they could be their own antiparticles

# Neutrino Masses

- Majorana or Dirac?
- Neutrinos are neutral so they could be their own antiparticles

# Experimental Set Up

---

## Experimental Set Up

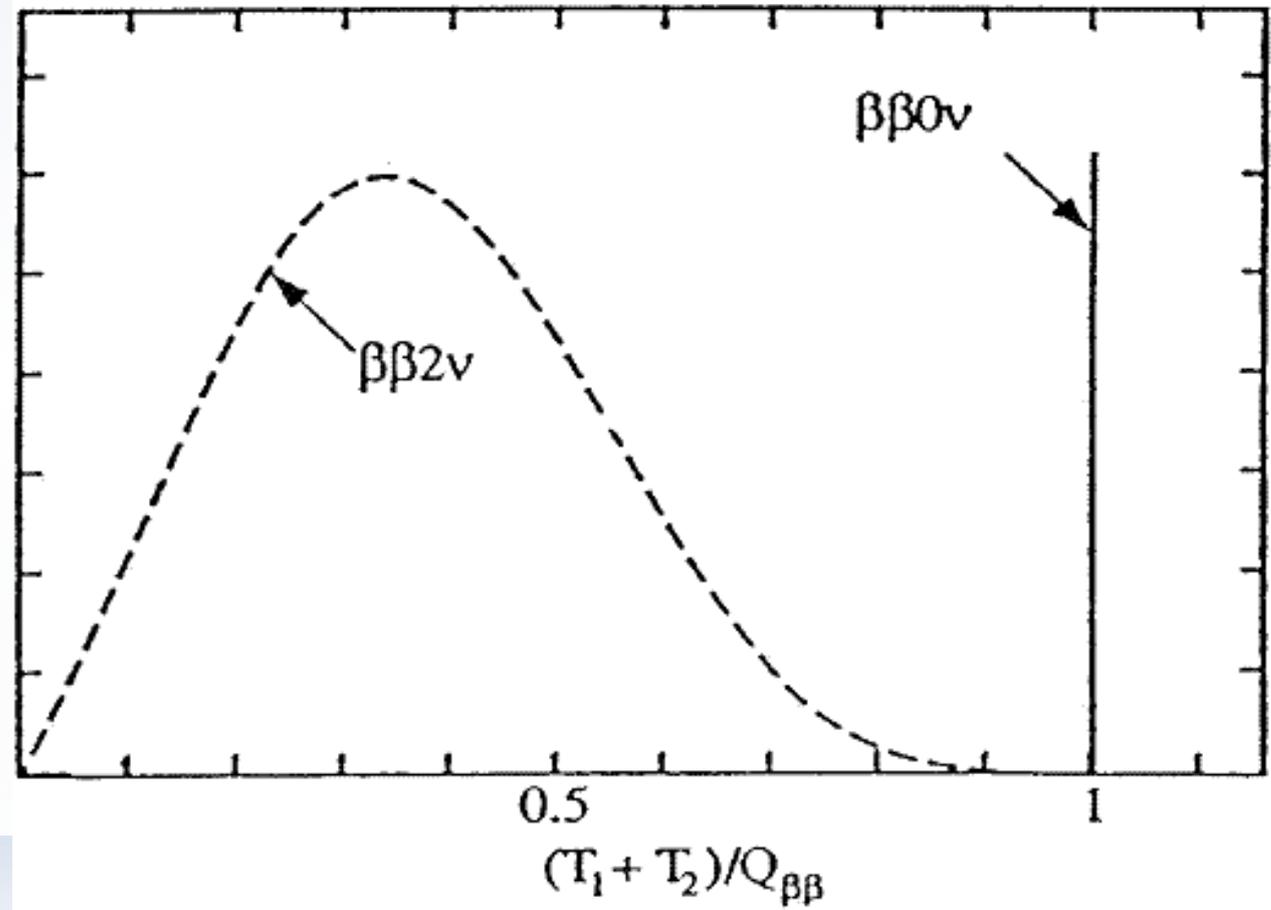
- $2\nu\beta\beta$ :  $N(A,Z) \rightarrow N(A, Z+2) + 2e^- + 2\bar{\nu}_e$
- $0\nu\beta\beta$ :  $N(A,Z) \rightarrow N(A, Z+2) + 2e^-$

## Experimental Set Up

- $2\nu\beta\beta$ :  $N(A,Z) \rightarrow N(A, Z+2) + 2e^- + 2\bar{\nu}_e$ 
  - 4-body decay
- $0\nu\beta\beta$ :  $N(A,Z) \rightarrow N(A, Z+2) + 2e^-$ 
  - 2-body decay

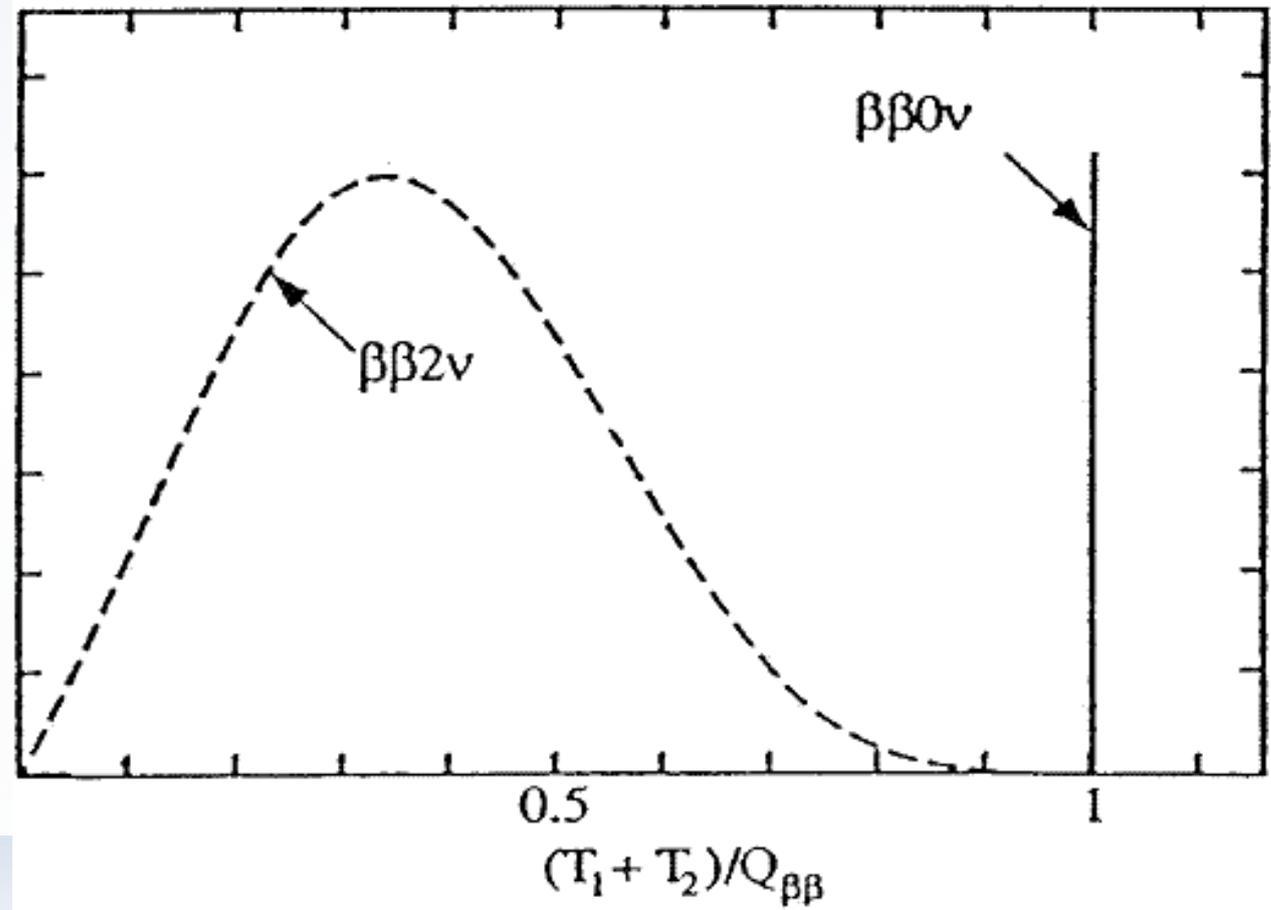
## Experimental Set Up

- $2\nu\beta\beta$ :  $N(A,Z) \rightarrow N(A, Z+2) + 2e^- + 2\bar{\nu}_e$ 
  - 4-body decay
- $0\nu\beta\beta$ :  $N(A,Z) \rightarrow N(A, Z+2) + 2e^-$ 
  - 2-body decay



# Experimental Set Up

- $2\nu\beta\beta$ :  $N(A,Z) \rightarrow N(A, Z+2) + 2e^- + 2\bar{\nu}_e$ 
  - 4-body decay
- $0\nu\beta\beta$ :  $N(A,Z) \rightarrow N(A, Z+2) + 2e^-$ 
  - 2-body decay
  - $\nu=\bar{\nu} \rightarrow$  Majorana



# Experimental Set Up

- $2\nu\beta\beta: N(A,Z) \rightarrow N(A, Z+2) + 2e^- + 2\bar{\nu}_e$

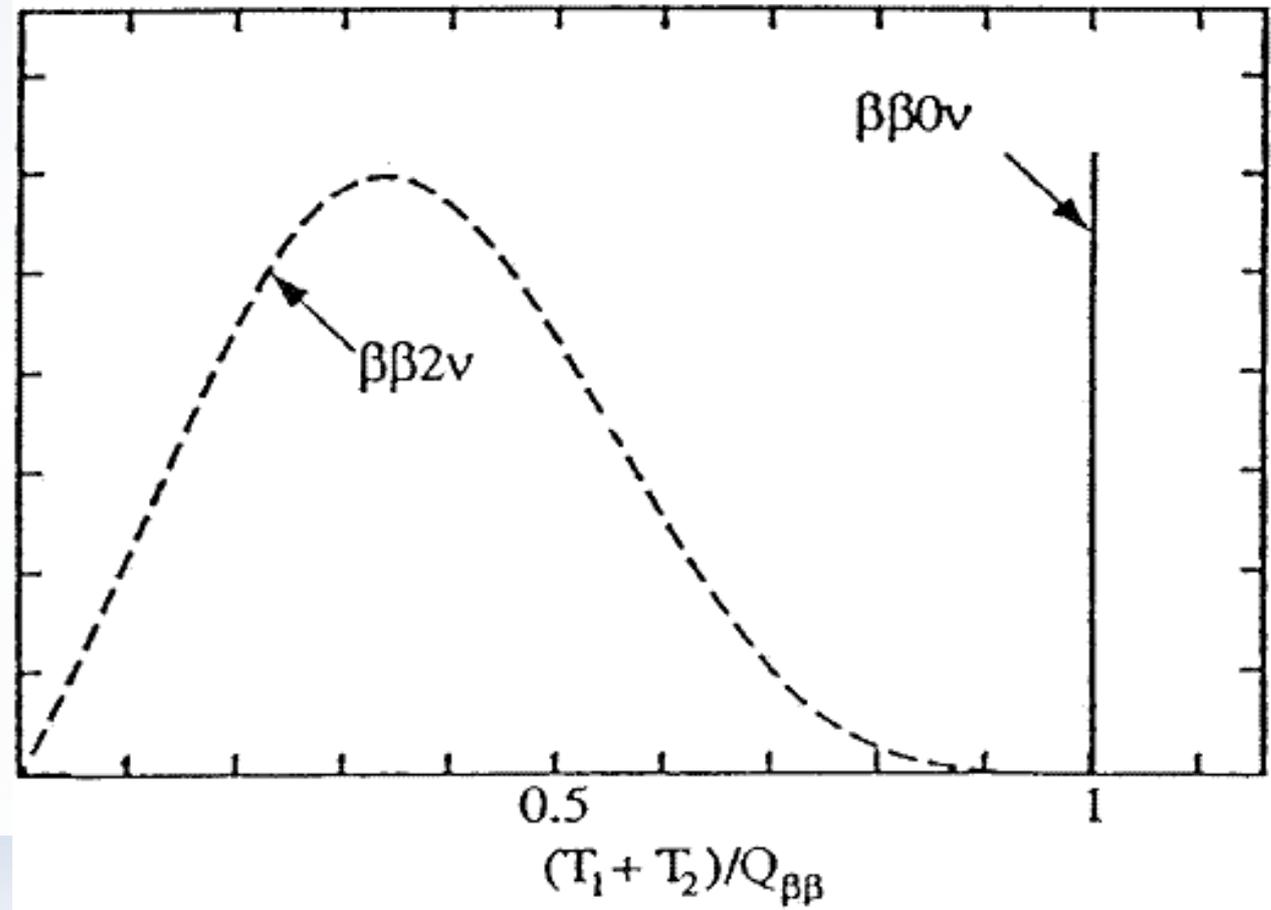
- 4-body decay

- $0\nu\beta\beta: N(A,Z) \rightarrow N(A, Z+2) + 2e^-$

- 2-body decay

- $\nu=\bar{\nu} \rightarrow$  Majorana

- Decay of  $^{76}\text{Ge}$   
and  $^{136}\text{Xe}$



# Half lives

---

## Half lives

- Many experiments looking at this process: Gerda, Heidelberg-Moscow, IGEX, Exo, KamLand-Zen
  - $T_{1/2}$  for  $^{76}\text{Ge} = 3.0 \times 10^{25}$  Yrs
  - $T_{1/2}$  for  $^{136}\text{Xe} = 1.9 \times 10^{25}$  Yrs

# Half lives

- Many experiments looking at this process: Gerda, Heidelberg-Moscow, IGEX, Exo, KamLand-Zen
  - $T_{1/2}$  for  $^{76}\text{Ge} = 3.0 \times 10^{25}$  Yrs
  - $T_{1/2}$  for  $^{136}\text{Xe} = 1.9 \times 10^{25}$  Yrs

$$T_{1/2}^{-1} = G_x(Q, Z) |\mathcal{M}_x(A, Z) \eta_x|^2$$

# Half lives

- Many experiments looking at this process: Gerda, Heidelberg-Moscow, IGEX, Exo, KamLand-Zen
- $T_{1/2}$  for  $^{76}\text{Ge} = 3.0 \times 10^{25}$  Yrs
- $T_{1/2}$  for  $^{136}\text{Xe} = 1.9 \times 10^{25}$  Yrs

$$T_{1/2}^{-1} = G_x(Q, Z) |\mathcal{M}_x(A, Z) \eta_x|^2$$

↑  
↓  
Phase Space Factor

↑  
↓  
Nuclear Matrix Element

# Half lives

- Many experiments looking at this process: Gerda, Heidelberg-Moscow, IGEX, Exo, KamLand-Zen
- $T_{1/2}$  for  $^{76}\text{Ge} = 3.0 \times 10^{25}$  Yrs
- $T_{1/2}$  for  $^{136}\text{Xe} = 1.9 \times 10^{25}$  Yrs

$$T_{1/2}^{-1} = G_x(Q, Z) |\mathcal{M}_x(A, Z) \eta_x|^2$$

↑  
↓  
Phase Space Factor

↑  
↓  
Nuclear Matrix Element

# Half lives

- Many experiments looking at this process: Gerda, Heidelberg-Moscow, IGEX, Exo, KamLand-Zen

- $T_{1/2}$  for  $^{76}\text{Ge} = 3.0 \times 10^{25}$  Yrs

- $T_{1/2}$  for  $^{136}\text{Xe} = 1.9 \times 10^{25}$  Yrs

$$T_{1/2}^{-1} = G_x(Q, Z) |\mathcal{M}_x(A, Z) \eta_x|^2$$

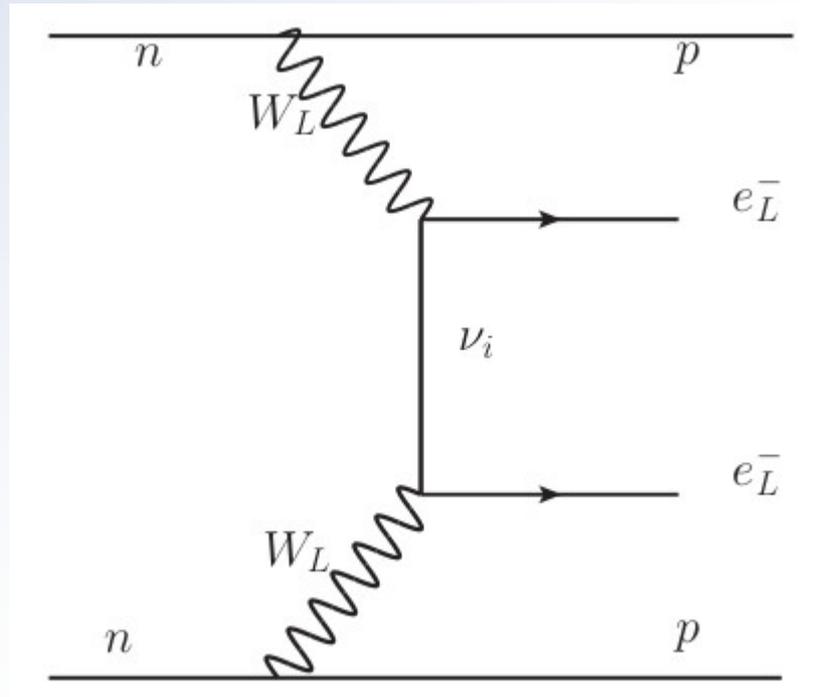
↑  
↓  
Phase Space Factor

↑  
↓  
Nuclear Matrix Element

- $\eta_x \propto$  Amplitude of different processes

# Standard $0\nu\beta\beta$

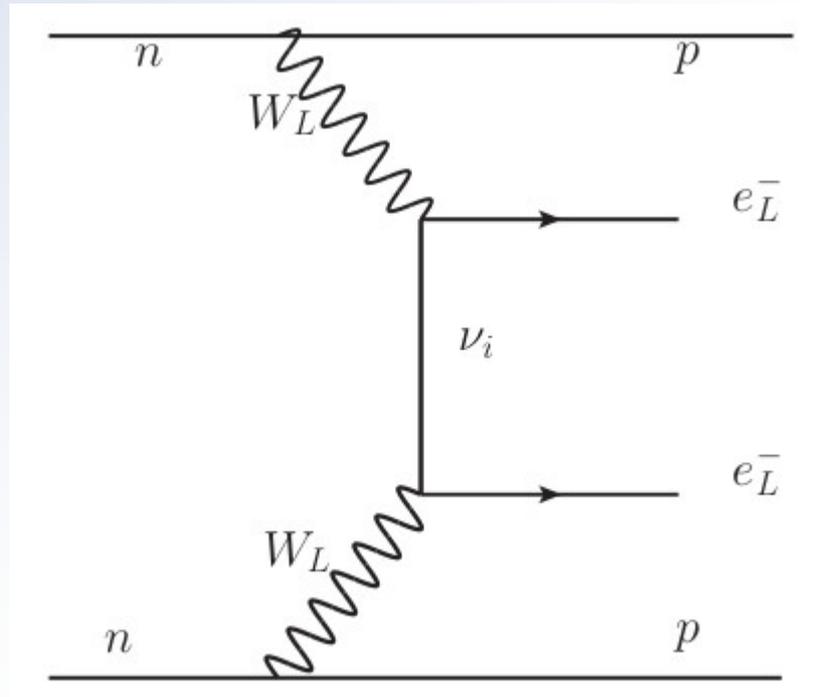
- Light neutrino exchange



arXiv:1204.2527

# Standard $0\nu\beta\beta$

- Light neutrino exchange

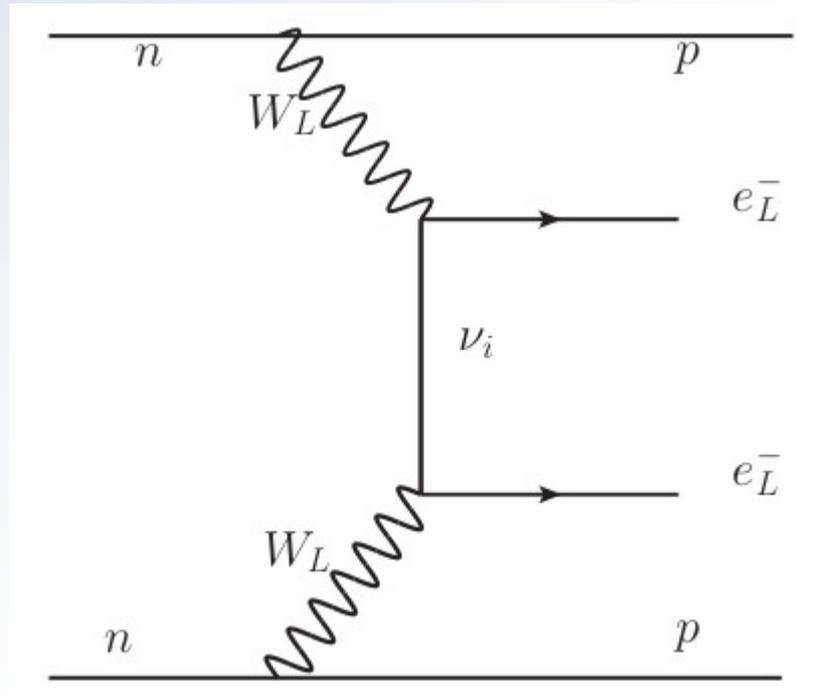


arXiv:1204.2527

$$\mathcal{A}_\nu = \sum_i^3 G_F^2 U_{ei}^2 \frac{m_i}{q^2 - m_i^2} \approx \sum_i^3 G_F^2 U_{ei}^2 \frac{m_i}{q^2} \propto \frac{m_\nu^{ee}}{q^2}$$

# Standard $0\nu\beta\beta$

- Light neutrino exchange

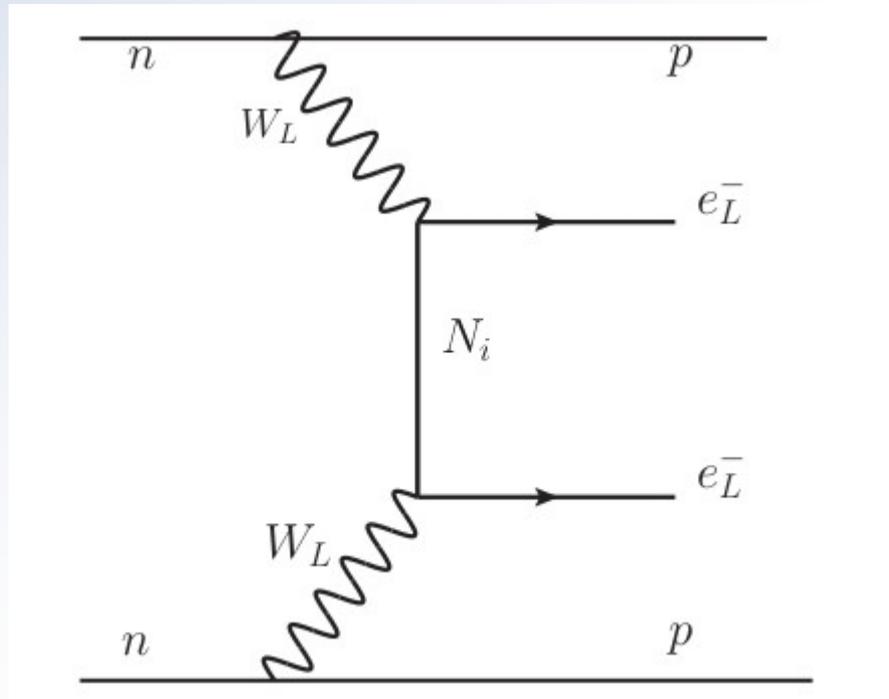


arXiv:1204.2527

$$A_\nu = \sum_i^3 G_F^2 U_{ei}^2 \frac{m_i}{q^2 - m_i^2} \approx \sum_i^3 G_F^2 U_{ei}^2 \frac{m_i}{q^2} \propto \frac{m_\nu^{ee}}{q^2}$$

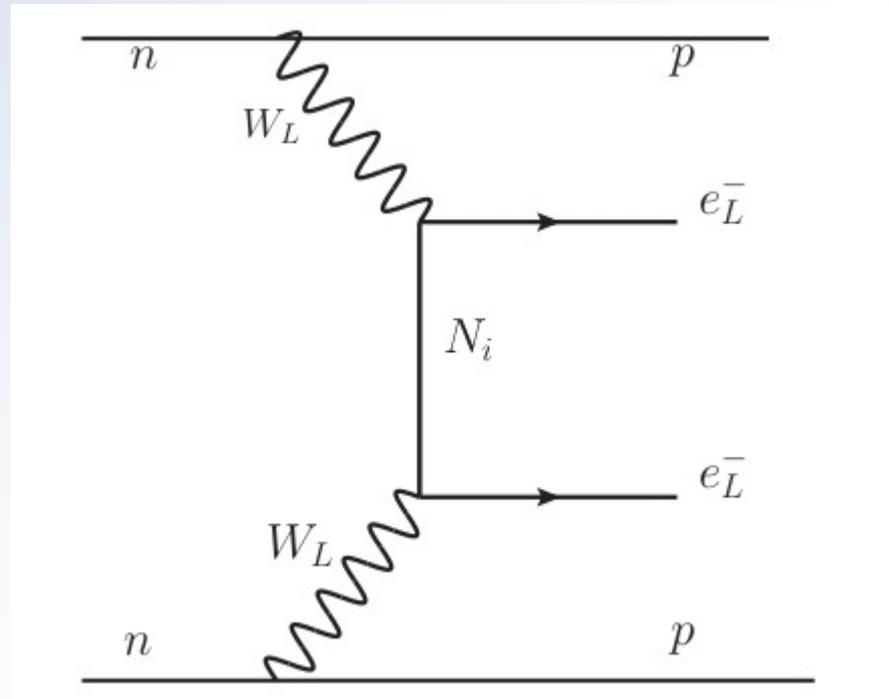
For  $q \approx 100 \text{ MeV} \gg m_i$

# Heavy neutrino exchange



arXiv:1204.2527

# Heavy neutrino exchange



arXiv:1204.2527

$$\mathcal{A}_N = \sum_i^3 G_F^2 T_{ei}^2 \frac{M_i}{q^2 - M_i^2} \approx \sum_i^3 G_F^2 T_{ei}^2 \frac{1}{M_i} \propto \frac{1}{m_N^{ee}}$$

Assuming  $M_i \gg q$

# Motivations for LRSM

---

# Motivations for LRSM

- Experimental: Predicts new particles that can be found at colliders

# Motivations for LRSM

- Experimental: Predicts new particles that can be found at colliders
- Theoretical: Can be embedded in GUT models like  $SO(10)$

# Motivations for LRSM

- Experimental: Predicts new particles that can be found at colliders
- Theoretical: Can be embedded in GUT models like  $SO(10)$
- Will introduce new channels that contribute to  $0\nu\beta\beta$

## Left-Right Symmetric Model

$$SU(2)_L \times U(1)_Y \rightarrow SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

## Left-Right Symmetric Model

$$SU(2)_L \times U(1)_Y \rightarrow SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

- Discrete Symmetry so that  $g_L = g_R$

# Left-Right Symmetric Model

$$SU(2)_L \times U(1)_Y \rightarrow SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

- Discrete Symmetry so that  $g_L = g_R$

Field	Form	$SU(2)_L$	$SU(2)_R$	$U(1)_{B-L}$
$Q_L$	$\begin{pmatrix} u \\ d \end{pmatrix}_L$	2	1	$\frac{1}{3}$
$Q_R$	$\begin{pmatrix} u \\ d \end{pmatrix}_R$	1	2	$\frac{1}{3}$
$\psi_L$	$\begin{pmatrix} \nu \\ l \end{pmatrix}_L$	2	1	-1
$\psi_R$	$\begin{pmatrix} \nu \\ l \end{pmatrix}_R$	1	2	-1
$\phi$	$\begin{pmatrix} \phi_1^0 & \phi_1^+ \\ \phi_2^- & \phi_2^0 \end{pmatrix}$	2	2	0
$\Delta_L$	$\begin{pmatrix} \frac{\Delta_L^+}{\sqrt{2}} & \Delta_L^{++} \\ \Delta_L^0 & \frac{-\Delta_L^+}{\sqrt{2}} \end{pmatrix}$	3	1	2
$\Delta_R$	$\begin{pmatrix} \frac{\Delta_R^+}{\sqrt{2}} & \Delta_R^{++} \\ \Delta_R^0 & \frac{-\Delta_R^+}{\sqrt{2}} \end{pmatrix}$	1	3	2

# Left-Right Symmetric Model

$$SU(2)_L \times U(1)_Y \rightarrow SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

- Discrete Symmetry so that  $g_L = g_R$

Field	Form	SU(2) <sub>L</sub>	SU(2) <sub>R</sub>	U(1) <sub>B-L</sub>
$Q_L$	$\begin{pmatrix} u \\ d \end{pmatrix}_L$	2	1	$\frac{1}{3}$
$Q_R$	$\begin{pmatrix} u \\ d \end{pmatrix}_R$	1	2	$\frac{1}{3}$
$\psi_L$	$\begin{pmatrix} \nu \\ l \end{pmatrix}_L$	2	1	-1
$\psi_R$	$\begin{pmatrix} \nu \\ l \end{pmatrix}_R$	1	2	-1
$\phi$	$\begin{pmatrix} \phi_1^0 & \phi_1^+ \\ \phi_2^- & \phi_2^0 \end{pmatrix}$	2	2	0
$\Delta_L$	$\begin{pmatrix} \frac{\Delta_L^+}{\sqrt{2}} & \Delta_L^{++} \\ \Delta_L^0 & \frac{-\Delta_L^+}{\sqrt{2}} \end{pmatrix}$	3	1	2
$\Delta_R$	$\begin{pmatrix} \frac{\Delta_R^+}{\sqrt{2}} & \Delta_R^{++} \\ \Delta_R^0 & \frac{-\Delta_R^+}{\sqrt{2}} \end{pmatrix}$	1	3	2

# Left-Right Symmetric Model

$$SU(2)_L \times U(1)_Y \rightarrow SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

- Discrete Symmetry so that  $g_L = g_R$

Type I



Field	Form	SU(2) <sub>L</sub>	SU(2) <sub>R</sub>	U(1) <sub>B-L</sub>
$Q_L$	$\begin{pmatrix} u \\ d \end{pmatrix}_L$	2	1	$\frac{1}{3}$
$Q_R$	$\begin{pmatrix} u \\ d \end{pmatrix}_R$	1	2	$\frac{1}{3}$
$\psi_L$	$\begin{pmatrix} \nu \\ l \end{pmatrix}_L$	2	1	-1
$\psi_R$	$\begin{pmatrix} \nu \\ l \end{pmatrix}_R$	1	2	-1
$\phi$	$\begin{pmatrix} \phi_1^0 & \phi_1^+ \\ \phi_2^- & \phi_2^0 \end{pmatrix}$	2	2	0
$\Delta_L$	$\begin{pmatrix} \frac{\Delta_L^+}{\sqrt{2}} & \Delta_L^{++} \\ \Delta_L^0 & \frac{-\Delta_L^+}{\sqrt{2}} \end{pmatrix}$	3	1	2
$\Delta_R$	$\begin{pmatrix} \frac{\Delta_R^+}{\sqrt{2}} & \Delta_R^{++} \\ \Delta_R^0 & \frac{-\Delta_R^+}{\sqrt{2}} \end{pmatrix}$	1	3	2

# Left-Right Symmetric Model

$$SU(2)_L \times U(1)_Y \rightarrow SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

- Discrete Symmetry so that  $g_L = g_R$

Type I



Field	Form	SU(2) <sub>L</sub>	SU(2) <sub>R</sub>	U(1) <sub>B-L</sub>
$Q_L$	$\begin{pmatrix} u \\ d \end{pmatrix}_L$	2	1	$\frac{1}{3}$
$Q_R$	$\begin{pmatrix} u \\ d \end{pmatrix}_R$	1	2	$\frac{1}{3}$
$\psi_L$	$\begin{pmatrix} \nu \\ l \end{pmatrix}_L$	2	1	-1
$\psi_R$	$\begin{pmatrix} \nu \\ l \end{pmatrix}_R$	1	2	-1
$\phi$	$\begin{pmatrix} \phi_1^0 & \phi_1^+ \\ \phi_2^- & \phi_2^0 \end{pmatrix}$	2	2	0
$\Delta_L$	$\begin{pmatrix} \frac{\Delta_L^+}{\sqrt{2}} & \Delta_L^{++} \\ \Delta_L^0 & \frac{-\Delta_L^+}{\sqrt{2}} \end{pmatrix}$	3	1	2
$\Delta_R$	$\begin{pmatrix} \frac{\Delta_R^+}{\sqrt{2}} & \Delta_R^{++} \\ \Delta_R^0 & \frac{-\Delta_R^+}{\sqrt{2}} \end{pmatrix}$	1	3	2

# Left-Right Symmetric Model

$$SU(2)_L \times U(1)_Y \rightarrow SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

- Discrete Symmetry so that  $g_L = g_R$

Field	Form	$SU(2)_L$	$SU(2)_R$	$U(1)_{B-L}$
$Q_L$	$\begin{pmatrix} u \\ d \end{pmatrix}_L$	2	1	$\frac{1}{3}$
$Q_R$	$\begin{pmatrix} u \\ d \end{pmatrix}_R$	1	2	$\frac{1}{3}$
$\psi_L$	$\begin{pmatrix} \nu \\ l \end{pmatrix}_L$	2	1	-1
Type I ← $\psi_R$	$\begin{pmatrix} \nu \\ l \end{pmatrix}_R$	1	2	-1
$\phi$	$\begin{pmatrix} \phi_1^0 & \phi_1^+ \\ \phi_2^- & \phi_2^0 \end{pmatrix}$	2	2	0
Type II ← $\Delta_L$	$\begin{pmatrix} \frac{\Delta_L^+}{\sqrt{2}} & \Delta_L^{++} \\ \Delta_L^0 & \frac{-\Delta_L^+}{\sqrt{2}} \end{pmatrix}$	3	1	2
$\Delta_R$	$\begin{pmatrix} \frac{\Delta_R^+}{\sqrt{2}} & \Delta_R^{++} \\ \Delta_R^0 & \frac{-\Delta_R^+}{\sqrt{2}} \end{pmatrix}$	1	3	2

## Extended Lagrangian

$$\begin{aligned} \mathcal{L}_Y = & h_1 \bar{\psi}_L \phi \psi_R + h_2 \bar{\psi}_L \tilde{\phi} \psi_R + h_3 \bar{Q}_L \phi Q_R + h_4 \bar{Q}_L \tilde{\phi} Q_R + \\ & + i h_5 \psi_L^T C \sigma_2 \Delta_L \psi_L + i h_6 \psi_R^T C \sigma_2 \Delta_R \psi_R + h.c. \end{aligned}$$

## Extended Lagrangian

$$\mathcal{L}_Y = h_1 \bar{\psi}_L \phi \psi_R + h_2 \bar{\psi}_L \tilde{\phi} \psi_R + h_3 \bar{Q}_L \phi Q_R + h_4 \bar{Q}_L \tilde{\phi} Q_R + \\ + i h_5 \psi_L^T C \sigma_2 \Delta_L \psi_L + i h_6 \psi_R^T C \sigma_2 \Delta_R \psi_R + h.c.$$

- Symmetry breaking occurs in two steps and  $\langle \Delta_R \rangle \gg \langle \Delta_L \rangle$

$$\phi = \begin{pmatrix} \kappa_1 & 0 \\ 0 & \kappa_2 \end{pmatrix}, \quad \Delta_{L,R} = \begin{pmatrix} 0 & 0 \\ \kappa_{L,R} & 0 \end{pmatrix}$$

## Extended Lagrangian

$$\mathcal{L}_Y = h_1 \bar{\psi}_L \phi \psi_R + h_2 \bar{\psi}_L \tilde{\phi} \psi_R + h_3 \bar{Q}_L \phi Q_R + h_4 \bar{Q}_L \tilde{\phi} Q_R + \\ + i h_5 \psi_L^T C \sigma_2 \Delta_L \psi_L + i h_6 \psi_R^T C \sigma_2 \Delta_R \psi_R + h.c.$$

- Symmetry breaking occurs in two steps and  $\langle \Delta_R \rangle \gg \langle \Delta_L \rangle$

$$\phi = \begin{pmatrix} k_1 & 0 \\ 0 & k_2 \end{pmatrix}, \quad \Delta_{L,R} = \begin{pmatrix} 0 & 0 \\ \kappa_{L,R} & 0 \end{pmatrix}$$

$$\mathcal{L}_{\phi, \Delta_{L,R}}^\nu = (h_1 k_1 + h_2 k_2) \nu_L^T C \nu_R^c \\ + h_5 \kappa_L \nu_L^T C \nu_L + h_6 \kappa_R (\nu_R^c)^T C \nu_R^c + h.c.$$

## Extended Lagrangian

$$\mathcal{L}_Y = h_1 \bar{\psi}_L \phi \psi_R + h_2 \bar{\psi}_L \tilde{\phi} \psi_R + h_3 \bar{Q}_L \phi Q_R + h_4 \bar{Q}_L \tilde{\phi} Q_R + \\ + i h_5 \psi_L^T C \sigma_2 \Delta_L \psi_L + i h_6 \psi_R^T C \sigma_2 \Delta_R \psi_R + h.c.$$

- Symmetry breaking occurs in two steps and  $\langle \Delta_R \rangle \gg \langle \Delta_L \rangle$

$$\phi = \begin{pmatrix} k_1 & 0 \\ 0 & k_2 \end{pmatrix}, \quad \Delta_{L,R} = \begin{pmatrix} 0 & 0 \\ \kappa_{L,R} & 0 \end{pmatrix}$$

$$\mathcal{L}_{\phi, \Delta_{L,R}}^\nu = (h_1 k_1 + h_2 k_2) \nu_L^T C \nu_R^c \\ + h_5 \kappa_L \nu_L^T C \nu_L + h_6 \kappa_R (\nu_R^c)^T C \nu_R^c + h.c$$

$$m_\nu = h_5 \kappa_L - \frac{k_1^2}{4\kappa_R} h_1^T h_6^{-1} h_1$$

$$m_N = h_6 \kappa_R$$

# Extended Lagrangian

$$\mathcal{L}_Y = h_1 \bar{\psi}_L \phi \psi_R + h_2 \bar{\psi}_L \tilde{\phi} \psi_R + h_3 \bar{Q}_L \phi Q_R + h_4 \bar{Q}_L \tilde{\phi} Q_R + \\ + i h_5 \psi_L^T C \sigma_2 \Delta_L \psi_L + i h_6 \psi_R^T C \sigma_2 \Delta_R \psi_R + h.c.$$

- Symmetry breaking occurs in two steps and  $\langle \Delta_R \rangle \gg \langle \Delta_L \rangle$

$$\phi = \begin{pmatrix} k_1 & 0 \\ 0 & k_2 \end{pmatrix}, \quad \Delta_{L,R} = \begin{pmatrix} 0 & 0 \\ \kappa_{L,R} & 0 \end{pmatrix}$$

$$\mathcal{L}_{\phi, \Delta_{L,R}}^\nu = (h_1 k_1 + h_2 k_2) \nu_L^T C \nu_R^c \\ + h_5 \kappa_L \nu_L^T C \nu_L + h_6 \kappa_R (\nu_R^c)^T C \nu_R^c + h.c$$

$$\kappa_R \gg k_1 \gg \kappa_L$$

$$h_1 k_1 \gg h_2 k_2$$

$$m_\nu = h_5 \kappa_L - \frac{k_1^2}{4\kappa_R} h_1^T h_6^{-1} h_1$$

$$m_N = h_6 \kappa_R$$

# Extended Lagrangian

$$\mathcal{L}_Y = h_1 \bar{\psi}_L \phi \psi_R + h_2 \bar{\psi}_L \tilde{\phi} \psi_R + h_3 \bar{Q}_L \phi Q_R + h_4 \bar{Q}_L \tilde{\phi} Q_R + \\ + i h_5 \psi_L^T C \sigma_2 \Delta_L \psi_L + i h_6 \psi_R^T C \sigma_2 \Delta_R \psi_R + h.c.$$

- Symmetry breaking occurs in two steps and  $\langle \Delta_R \rangle \gg \langle \Delta_L \rangle$

$$\phi = \begin{pmatrix} k_1 & 0 \\ 0 & k_2 \end{pmatrix}, \quad \Delta_{L,R} = \begin{pmatrix} 0 & 0 \\ \kappa_{L,R} & 0 \end{pmatrix}$$

$$\mathcal{L}_{\phi, \Delta_{L,R}}^\nu = (h_1 k_1 + h_2 k_2) \nu_L^T C \nu_R^c \\ + h_5 \kappa_L \nu_L^T C \nu_L + h_6 \kappa_R (\nu_R^c)^T C \nu_R^c + h.c$$

$$\kappa_R \gg k_1 \gg \kappa_L$$

$$h_1 k_1 \gg h_2 k_2$$

$$m_\nu = h_5 \kappa_L - \frac{k_1^2}{4\kappa_R} h_1^T h_6^{-1} h_1$$

$$m_N = h_6 \kappa_R$$

## New contributions to $0\nu\beta\beta$

- Light and heavy neutrino exchange with two  $W_R$  bosons

## New contributions to $0\nu\beta\beta$

- Light and heavy neutrino exchange with two  $W_R$  bosons
- Light and heavy neutrino exchange with a mix of  $W_L$  and  $W_R$  bosons

# New contributions to $0\nu\beta\beta$

- Light and heavy neutrino exchange with two  $W_R$  bosons
- Light and heavy neutrino exchange with a mix of  $W_L$  and  $W_R$  bosons
- $\Delta_L^-$  Higgs triplet decay

# New contributions to $0\nu\beta\beta$

- Light and heavy neutrino exchange with two  $W_R$  bosons
- Light and heavy neutrino exchange with a mix of  $W_L$  and  $W_R$  bosons
- $\Delta_L^-$  Higgs triplet decay
- $\Delta_R^-$  Higgs triplet decay

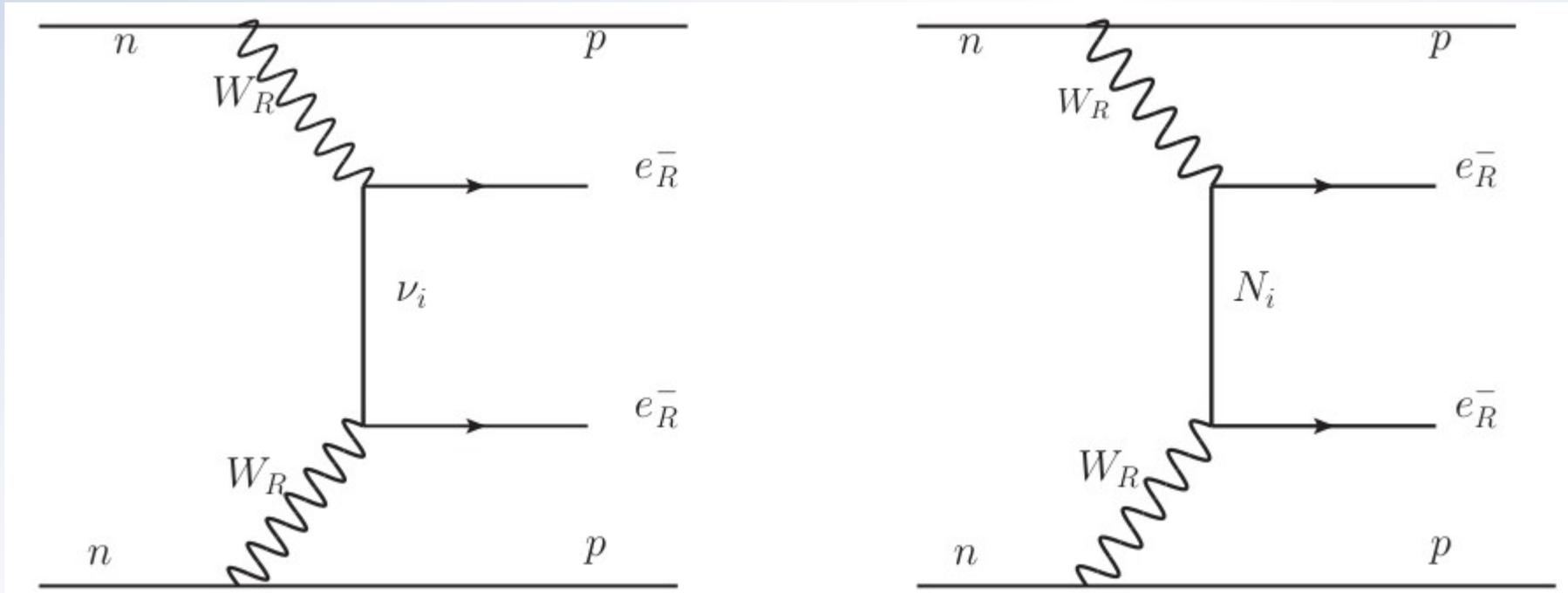
# Amplitudes I

- $W_R$  exchange

# Amplitudes I

- $W_R$  exchange

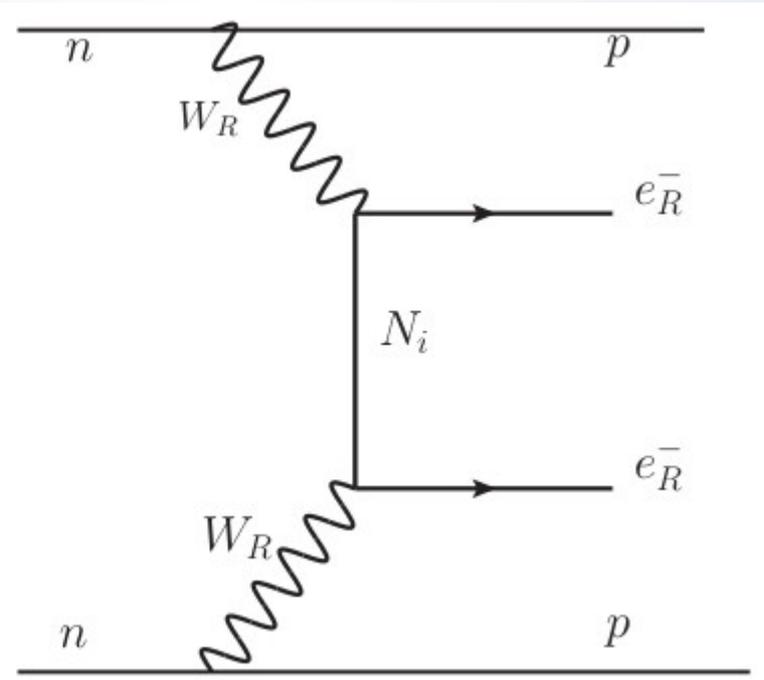
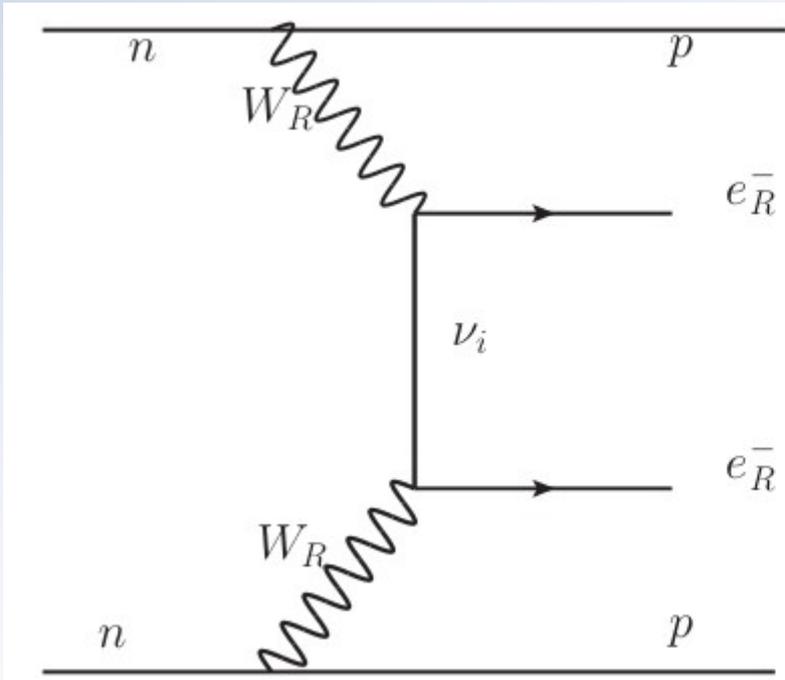
arXiv:1204.2527



# Amplitudes I

- $W_R$  exchange

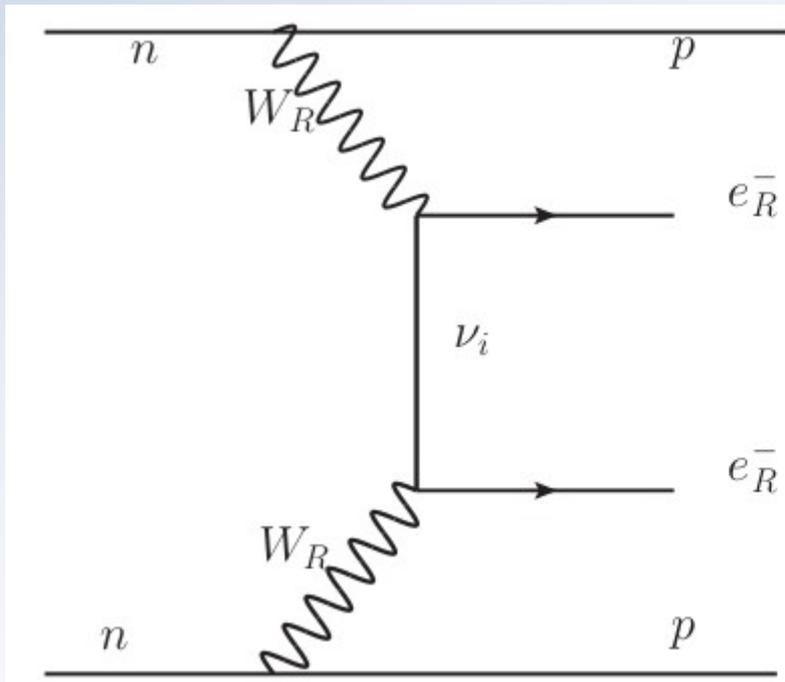
arXiv:1204.2527



$$A_{\nu}^{RR} \propto G_F^2 \frac{M_{W_L}^4}{M_{W_R}^4} \sum_i^3 \frac{(S^*)_{ei}^2 m_i}{q^2}$$

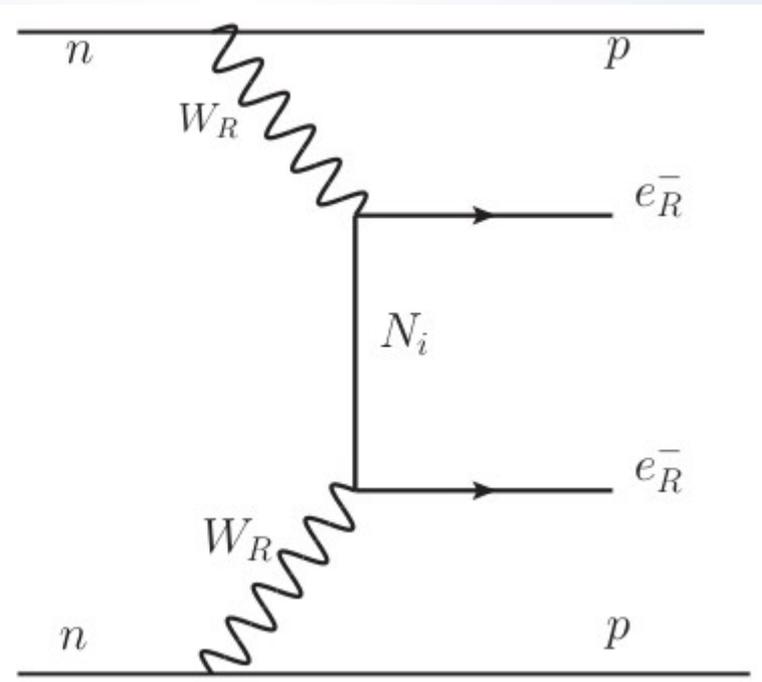
# Amplitudes I

- $W_R$  exchange



$$A_{\nu}^{RR} \propto G_F^2 \frac{M_{W_L}^4}{M_{W_R}^4} \sum_i^3 \frac{(S^*)_{ei}^2 m_i}{q^2}$$

arXiv:1204.2527

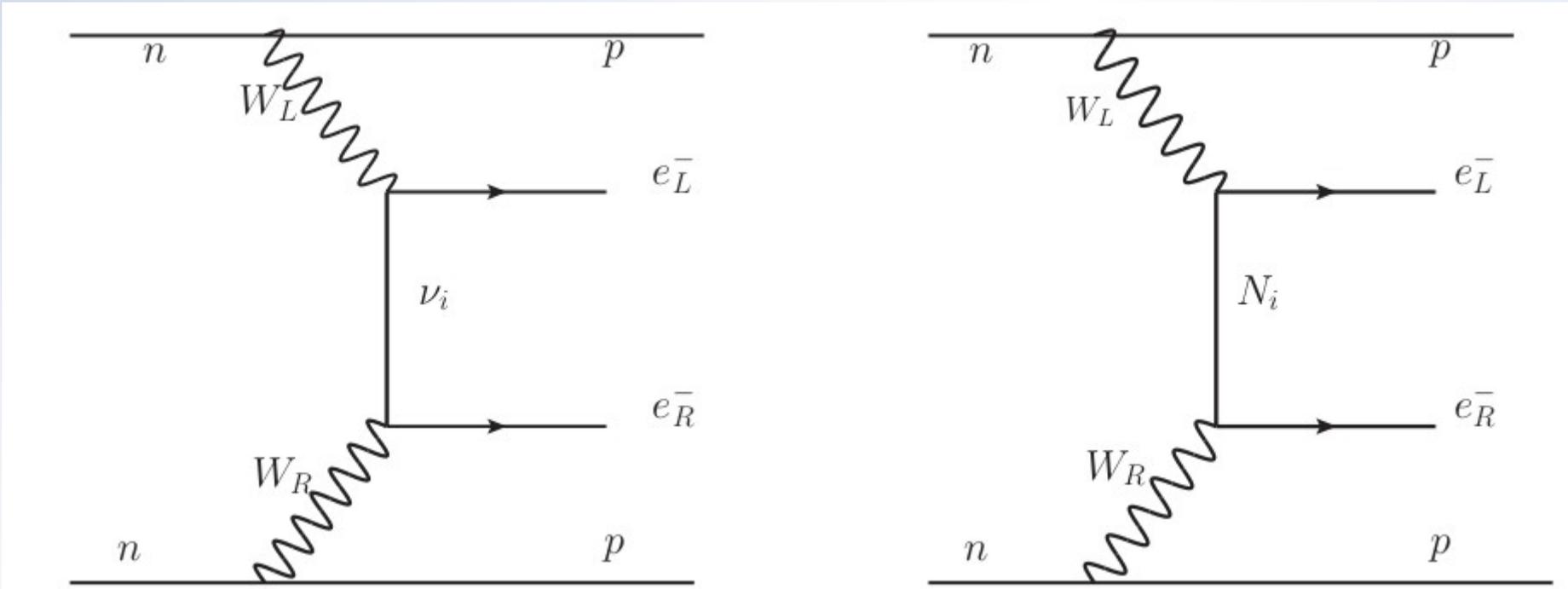


$$A_N^{RR} \propto G_F^2 \frac{M_{W_L}^4}{M_{W_R}^4} \sum_i^3 \frac{(U_R^*)_{ei}^2}{M_i}$$

# Amplitudes II

- $W_L$  and  $W_R$  exchange

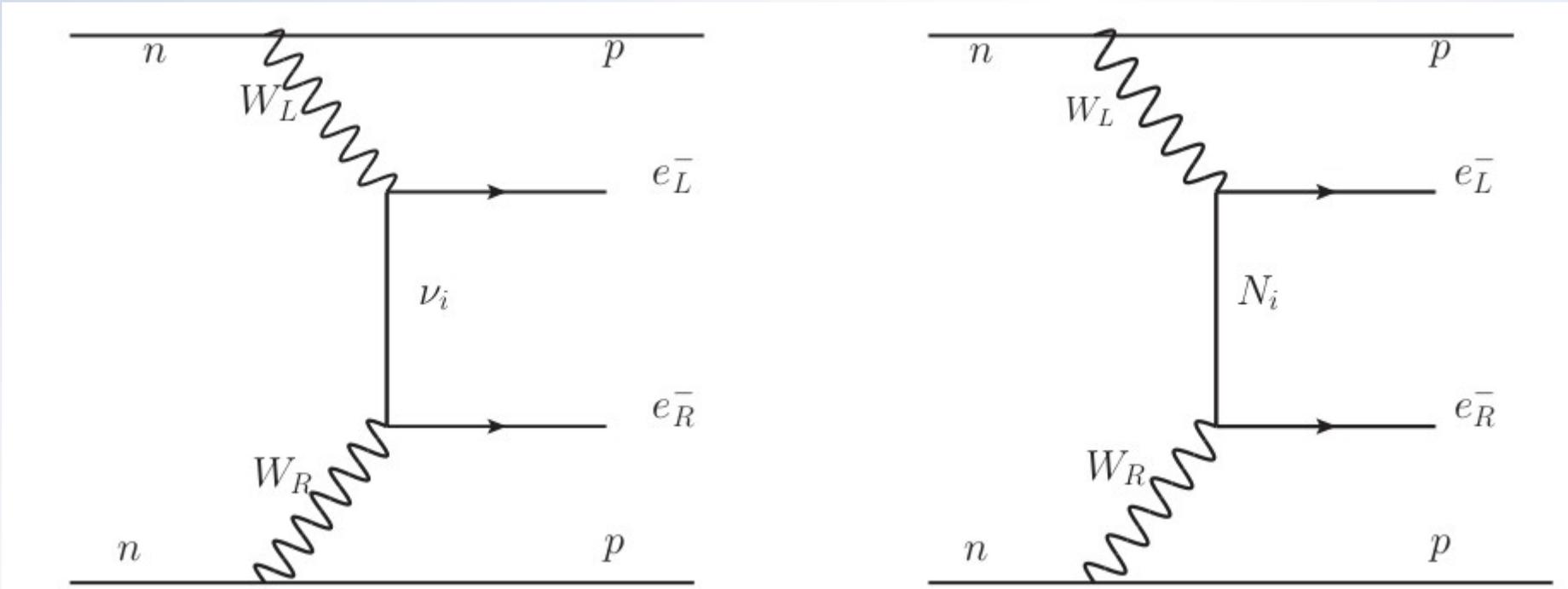
arXiv:1204.2527



# Amplitudes II

- $W_L$  and  $W_R$  exchange

arXiv:1204.2527

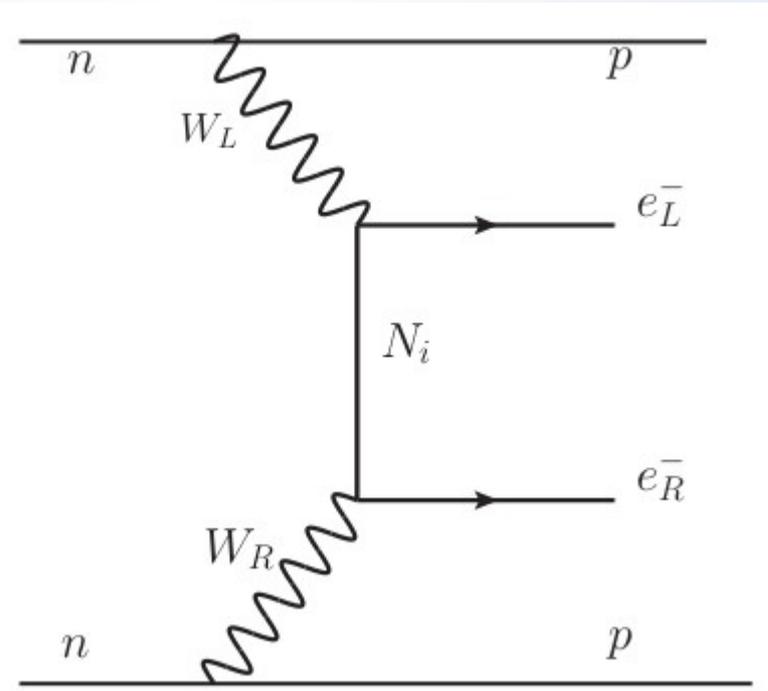
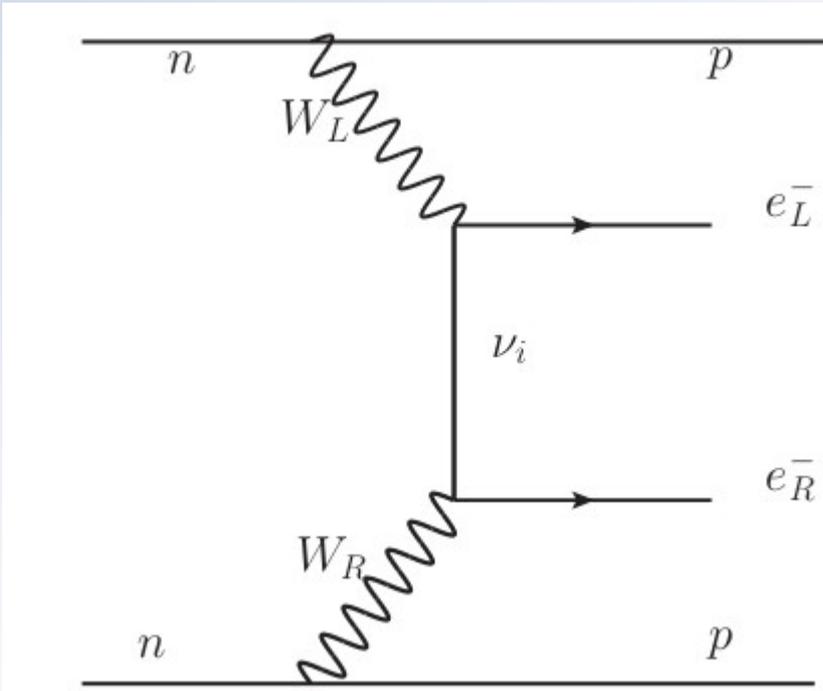


$$A_{\nu}^{LR} \propto G_F^2 \frac{M_{W_L}^2}{M_{W_R}^2} \sum_i^3 \frac{(U_L)_{ei} (S^*)_{ei}}{q}$$

# Amplitudes II

- $W_L$  and  $W_R$  exchange

arXiv:1204.2527



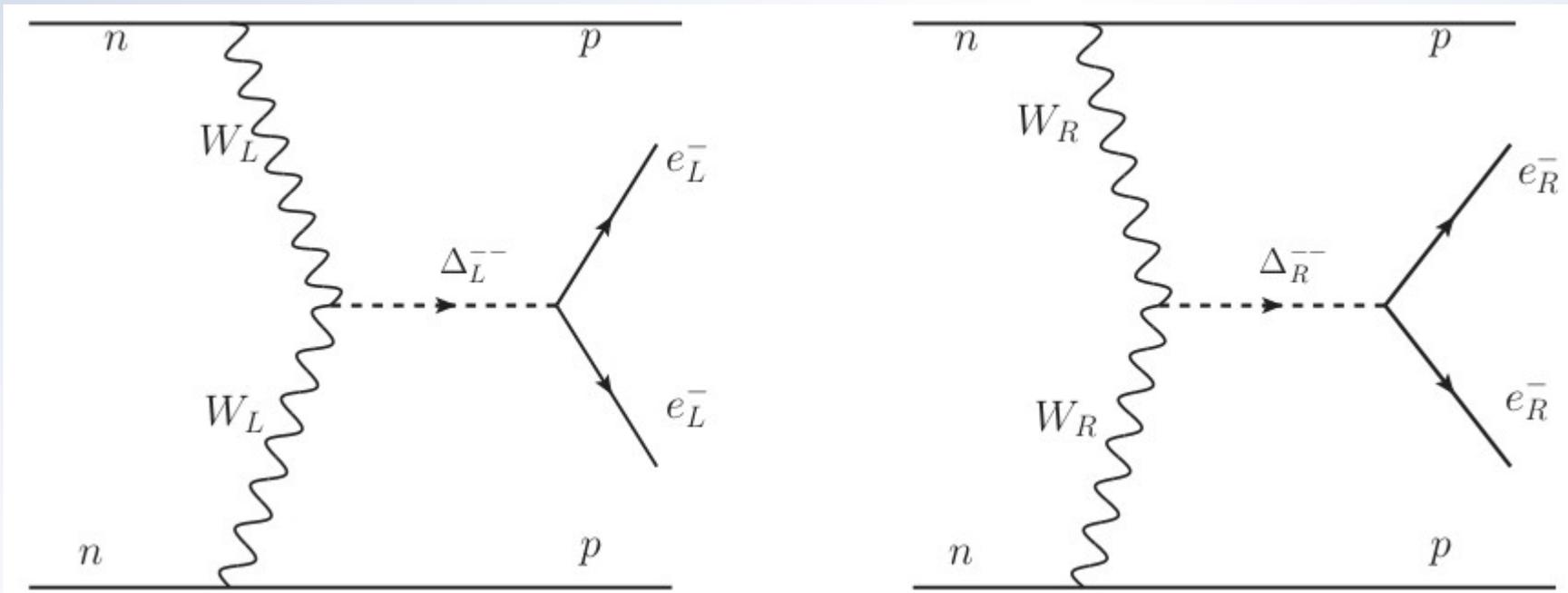
$$A_{\nu}^{LR} \propto G_F^2 \frac{M_{W_L}^2}{M_{W_R}^2} \sum_i^3 \frac{(U_L)_{ei} (S^*)_{ei}}{q}$$

$$A_N^{LR} \propto G_F^2 \frac{M_{W_L}^2}{M_{W_R}^2} \sum_i^3 \frac{(T)_{ei} (U_R^*)_{ei} q}{M_i^2}$$

# Amplitudes III

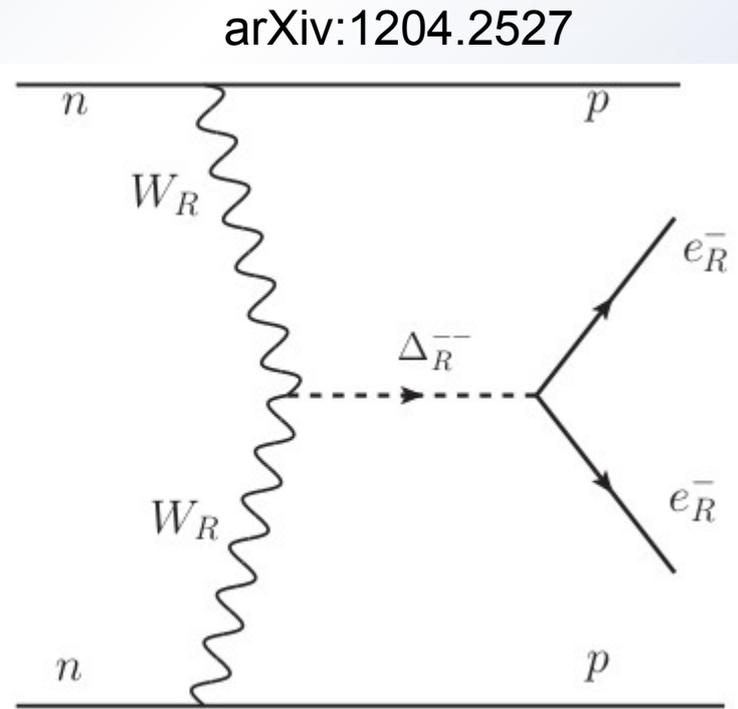
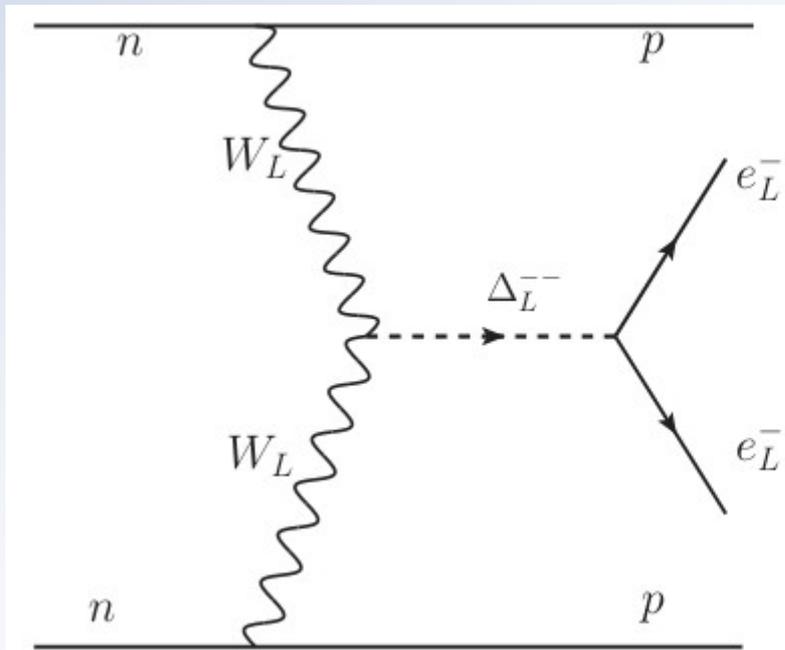
- $\Delta_L^-$  and  $\Delta_R^-$  exchange

arXiv:1204.2527



# Amplitudes III

- $\Delta_L^-$  and  $\Delta_R^-$  exchange

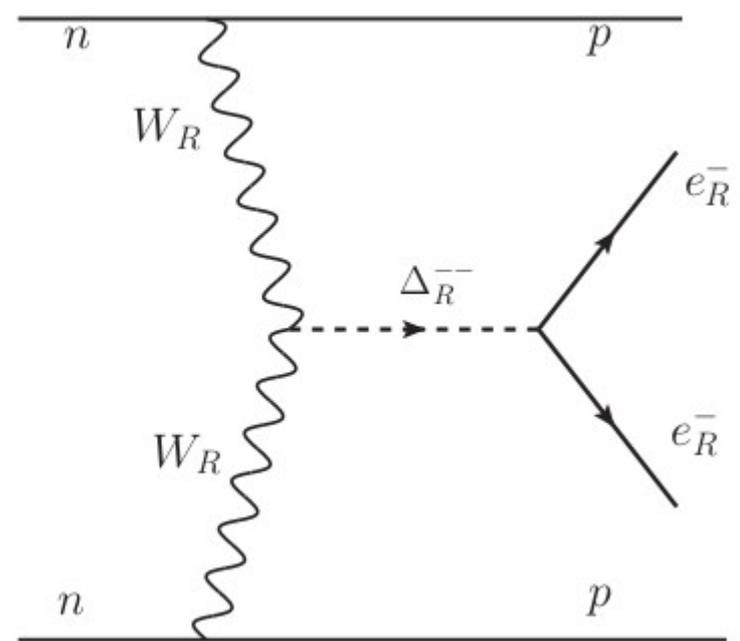
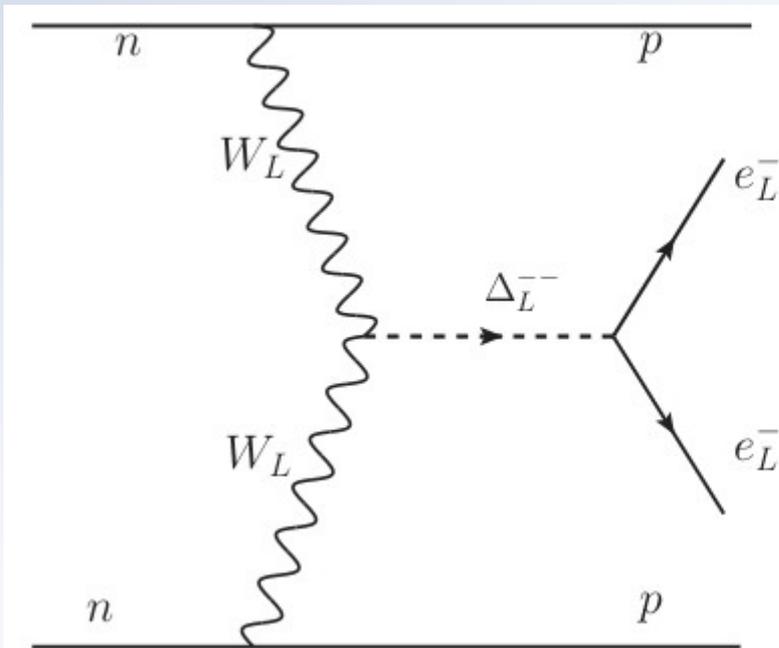


$$A_{\Delta_L}^{LL} \propto G_F^2 \sum_i^3 \frac{(U_L)_{ei}^2 m_i}{M_{\Delta_L}^2}$$

# Amplitudes III

- $\Delta_L^-$  and  $\Delta_R^-$  exchange

arXiv:1204.2527



$$A_{\Delta_L}^{LL} \propto G_F^2 \sum_i^3 \frac{(U_L)_{ei}^2 m_i}{M_{\Delta_L}^2}$$

$$A_{\Delta_R}^{RR} \propto G_F^2 \frac{M_{W_R}^4}{M_{W_L}^4} \sum_i^3 \frac{(U_R^*)_{ei}^2 M_i}{M_{\Delta_R}^2}$$

## New Physics Parameters

$$T_{1/2}^{-1} = G_x(Q, Z) |\mathcal{M}_x(A, Z) \eta_x|^2$$

- NME and Phase Factors can be found in the literature

# New Physics Parameters

$$T_{1/2}^{-1} = G_x(Q, Z) |\mathcal{M}_x(A, Z) \eta_x|^2$$

- NME and Phase Factors can be found in the literature
- Define  $\eta_x$ :

$$\eta_\nu^{LL} = \sum_i^3 \frac{(U_L)_{ei}^2 m_i}{m_e} = \frac{m_\nu^{ee}}{m_e}$$
$$\eta_N^{LL} = \sum_i^3 \frac{(T)_{ei}^2 m_p}{M_i} = \frac{m_p}{m_N^{ee}}$$

# New Physics Parameters

$$T_{1/2}^{-1} = G_x(Q, Z) |\mathcal{M}_x(A, Z) \eta_x|^2$$

- NME and Phase Factors can be found in the literature
- Define  $\eta_x$ :

$$\eta_\nu^{LL} = \sum_i^3 \frac{(U_L)_{ei}^2 m_i}{m_e} = \frac{m_\nu^{ee}}{m_e}$$

$$\eta_N^{LL} = \sum_i^3 \frac{(T)_{ei}^2 m_p}{M_i} = \frac{m_p}{m_N^{ee}}$$

- From limits in  $T_{1/2}$  can find limits in  $m_\nu^{ee}$  and  $m_N^{ee}$

# Limits

- For standard mechanism

Isotope	Half-life $T_{1/2}^{0\nu\beta\beta}$ (Yrs)	$m_\nu^{ee}$ (eV)	$\frac{1}{m_N^{ee}}$ (GeV <sup>-1</sup> )
<sup>76</sup> Ge	$3.0 \times 10^{25}$	0.29 – 0.74	$(0.97 – 1.72) \times 10^{-8}$
<sup>136</sup> Xe	$1.9 \times 10^{25}$	0.25 – 0.62	$(1.18 – 1.24) \times 10^{-8}$

# Limits

- For standard mechanism

Isotope	Half-life $T_{1/2}^{0\nu\beta\beta}$ (Yrs)	$m_\nu^{ee}$ (eV)	$\frac{1}{m_N^{ee}}$ (GeV <sup>-1</sup> )
<sup>76</sup> Ge	$3.0 \times 10^{25}$	0.29 – 0.74	$(0.97 – 1.72) \times 10^{-8}$
<sup>136</sup> Xe	$1.9 \times 10^{25}$	0.25 – 0.62	$(1.18 – 1.24) \times 10^{-8}$

- Very heavy Right handed neutrinos  $M_N^{ee} \sim 10^8$  GeV

# Limits

- For standard mechanism

Isotope	Half-life $T_{1/2}^{0\nu\beta\beta}$ (Yrs)	$m_\nu^{ee}$ (eV)	$\frac{1}{m_N^{ee}}$ (GeV <sup>-1</sup> )
<sup>76</sup> Ge	$3.0 \times 10^{25}$	0.29 – 0.74	$(0.97 – 1.72) \times 10^{-8}$
<sup>136</sup> Xe	$1.9 \times 10^{25}$	0.25 – 0.62	$(1.18 – 1.24) \times 10^{-8}$

- Very heavy Right handed neutrinos  $M_N^{ee} \sim 10^8$  GeV
- Range due to uncertainties in NME's →  
Main limitation

# Limits

- For standard mechanism

Isotope	Half-life $T_{1/2}^{0\nu\beta\beta}$ (Yrs)	$m_\nu^{ee}$ (eV)	$\frac{1}{m_N^{ee}}$ (GeV <sup>-1</sup> )
<sup>76</sup> Ge	$3.0 \times 10^{25}$	0.29 – 0.74	$(0.97 – 1.72) \times 10^{-8}$
<sup>136</sup> Xe	$1.9 \times 10^{25}$	0.25 – 0.62	$(1.18 – 1.24) \times 10^{-8}$

- Very heavy Right handed neutrinos  $M_N^{ee} \sim 10^8$  GeV
- Range due to uncertainties in NME's →  
Main limitation
- Limits assume contribution for only one mechanism at a time

# Limits LRSM

- For LRSM

$$\eta_{\nu}^{RR} = \frac{M_{WL}^4}{M_{WR}^4} \sum_i^3 \frac{(S^*)_{ei}^2 m_i}{m_e},$$

$$\eta_{\nu}^{LR} = \frac{M_{WL}^2}{M_{WR}^2} \sum_i^3 (U_L)_{ei} (S^*)_{ei},$$

$$\eta_{\Delta L}^{LL} = \sum_i^3 \frac{(U_L)_{ei}^2 m_i m_e}{M_{\Delta L}^2},$$

$$\eta_N^{RR} = \frac{M_{WL}^4}{M_{WR}^4} \sum_i^3 \frac{(U_R^*)_{ei}^2 m_p}{M_i}$$

$$\eta_N^{LR} = \frac{M_{WL}^2}{M_{WR}^2} \sum_i^3 \frac{(T)_{ei} (U_R^*)_{ei} m_p^2}{M_i^2}$$

$$\eta_{\Delta R}^{RR} = \frac{M_{WL}^4}{M_{WR}^4} \sum_i^3 \frac{(U_R^*)_{ei}^2 M_i m_p}{M_{\Delta R}^2}$$

# Limits LRSM

Mechanism	New physics parameters	$^{76}\text{Ge}$ Limit	$^{136}\text{Xe}$ Limit
$\mathcal{A}_\nu^{RR}$	$\sum_i^3 \frac{(S^*)_{ei}^2 m_i}{M_{WR}^4}$	$(0.70 - 1.77) \times 10^{-17} \text{ GeV}^{-3}$	$(0.60 - 1.49) \times 10^{-17} \text{ GeV}^{-3}$
$\mathcal{A}_N^{RR}$	$\sum_i^3 \frac{(U_R^*)_{ei}^2}{M_{WR}^4 M_i}$	$(2.32 - 4.12) \times 10^{-16} \text{ GeV}^{-5}$	$(2.83 - 2.97) \times 10^{-16} \text{ GeV}^{-5}$
$\mathcal{A}_\nu^{LR}$	$\sum_i^3 \frac{(U_L)_{ei}(S^*)_{ei}}{M_{WR}^2}$	$(1.54 - 3.32) \times 10^{-10} \text{ GeV}^{-2}$	$(1.18 - 1.50) \times 10^{-10} \text{ GeV}^{-2}$
$\mathcal{A}_N^{LR}$	$\sum_i^3 \frac{(T)_{ei}(U_R^*)_{ei}}{M_{WR}^2 M_i^2}$	$(1.95 - 2.04) \times 10^{-12} \text{ GeV}^{-4}$	$(1.60 - 2.83) \times 10^{-12} \text{ GeV}^{-4}$
$\mathcal{A}_{\Delta L}^{LL}$	$\sum_i^3 \frac{(U_L)_{ei}^2 m_i}{M_{\Delta L}^2}$	$\sim 10^{-8} \text{ GeV}^{-1}$	$\sim 10^{-8} \text{ GeV}^{-1}$
$\mathcal{A}_{\Delta R}^{RR}$	$\sum_i^3 \frac{(U_R^*)_{ei}^2 M_i}{M_{\Delta R}^2 M_{WR}^4}$	$\sim (10^{-16} - 10^{-15}) \text{ GeV}^{-5}$	$\sim (10^{-16} - 10^{-15}) \text{ GeV}^{-5}$

# Limits LRSM

Mechanism	New physics parameters	$^{76}\text{Ge}$ Limit	$^{136}\text{Xe}$ Limit
$\mathcal{A}_\nu^{RR}$	$\sum_i^3 \frac{(S^*)_{ei}^2 m_i}{M_{WR}^4}$	$(0.70 - 1.77) \times 10^{-17} \text{ GeV}^{-3}$	$(0.60 - 1.49) \times 10^{-17} \text{ GeV}^{-3}$
$\mathcal{A}_N^{RR}$	$\sum_i^3 \frac{(U_R^*)_{ei}^2}{M_{WR}^4 M_i}$	$(2.32 - 4.12) \times 10^{-16} \text{ GeV}^{-5}$	$(2.83 - 2.97) \times 10^{-16} \text{ GeV}^{-5}$
$\mathcal{A}_\nu^{LR}$	$\sum_i^3 \frac{(U_L)_{ei}(S^*)_{ei}}{M_{WR}^2}$	$(1.54 - 3.32) \times 10^{-10} \text{ GeV}^{-2}$	$(1.18 - 1.50) \times 10^{-10} \text{ GeV}^{-2}$
$\mathcal{A}_N^{LR}$	$\sum_i^3 \frac{(T)_{ei}(U_R^*)_{ei}}{M_{WR}^2 M_i^2}$	$(1.95 - 2.04) \times 10^{-12} \text{ GeV}^{-4}$	$(1.60 - 2.83) \times 10^{-12} \text{ GeV}^{-4}$
$\mathcal{A}_{\Delta L}^{LL}$	$\sum_i^3 \frac{(U_L)_{ei}^2 m_i}{M_{\Delta L}^2}$	$\sim 10^{-8} \text{ GeV}^{-1}$	$\sim 10^{-8} \text{ GeV}^{-1}$
$\mathcal{A}_{\Delta R}^{RR}$	$\sum_i^3 \frac{(U_R^*)_{ei}^2 M_i}{M_{\Delta R}^2 M_{WR}^4}$	$\sim (10^{-16} - 10^{-15}) \text{ GeV}^{-5}$	$\sim (10^{-16} - 10^{-15}) \text{ GeV}^{-5}$

# Limits LRSM

Mechanism	New physics parameters	$^{76}\text{Ge}$ Limit	$^{136}\text{Xe}$ Limit
$\mathcal{A}_\nu^{RR}$	$\sum_i^3 \frac{(S^*)_{ei}^2 m_i}{M_{WR}^4}$	$(0.70 - 1.77) \times 10^{-17} \text{ GeV}^{-3}$	$(0.60 - 1.49) \times 10^{-17} \text{ GeV}^{-3}$
$\mathcal{A}_N^{RR}$	$\sum_i^3 \frac{(U_R^*)_{ei}^2}{M_{WR}^4 M_i}$	$(2.32 - 4.12) \times 10^{-16} \text{ GeV}^{-5}$	$(2.83 - 2.97) \times 10^{-16} \text{ GeV}^{-5}$
$\mathcal{A}_\nu^{LR}$	$\sum_i^3 \frac{(U_L)_{ei}(S^*)_{ei}}{M_{WR}^2}$	$(1.54 - 3.32) \times 10^{-10} \text{ GeV}^{-2}$	$(1.18 - 1.50) \times 10^{-10} \text{ GeV}^{-2}$
$\mathcal{A}_N^{LR}$	$\sum_i^3 \frac{(T)_{ei}(U_R^*)_{ei}}{M_{WR}^2 M_i^2}$	$(1.95 - 2.04) \times 10^{-12} \text{ GeV}^{-4}$	$(1.60 - 2.83) \times 10^{-12} \text{ GeV}^{-4}$
$\mathcal{A}_{\Delta L}^{LL}$	$\sum_i^3 \frac{(U_L)_{ei}^2 m_i}{M_{\Delta L}^2}$	$\sim 10^{-8} \text{ GeV}^{-1}$	$\sim 10^{-8} \text{ GeV}^{-1}$
$\mathcal{A}_{\Delta R}^{RR}$	$\sum_i^3 \frac{(U_R^*)_{ei}^2 M_i}{M_{\Delta R}^2 M_{WR}^4}$	$\sim (10^{-16} - 10^{-15}) \text{ GeV}^{-5}$	$\sim (10^{-16} - 10^{-15}) \text{ GeV}^{-5}$

- Can set bounds on masses and mixing of new particles

# Limitations

---

## Limitations

- This is only a qualitative analysis which sets limits in new physics parameters

# Limitations

- This is only a qualitative analysis which sets limits in new physics parameters
- Some mechanisms are highly suppressed and it is important to consider relative magnitude of the different mechanisms

# Limitations

- This is only a qualitative analysis which sets limits in new physics parameters
- Some mechanisms are highly suppressed and it is important to consider relative magnitude of the different mechanisms
- Interference and cancellation between different mechanisms might occur

# Limitations

- This is only a qualitative analysis which sets limits in new physics parameters
- Some mechanisms are highly suppressed and it is important to consider relative magnitude of the different mechanisms
- Interference and cancellation between different mechanisms might occur
- There are many other possible contributions to  $0\nu\beta\beta$ . Real challenge is to identify the underlying mechanism that drives the decay

## Going Forward

- Katrin → Half live of  $6 \times 10^{27}$  Years

## Going Forward

- Katrin → Half live of  $6 \times 10^{27}$  Years
- Improvement of NME's uncertainties

## Going Forward

- Katrin → Half live of  $6 \times 10^{27}$  Years
- Improvement of NME's uncertainties
- Study of  $0\nu\beta\beta$  for different isotopes

## Going Forward

- Katrin → Half live of  $6 \times 10^{27}$  Years
- Improvement of NME's uncertainties
- Study of  $0\nu\beta\beta$  for different isotopes
- Study of  $e^-$  angular distribution at SuperNEMO

## Going Forward

- Katrin → Half live of  $6 \times 10^{27}$  Years
- Improvement of NME's uncertainties
- Study of  $0\nu\beta\beta$  for different isotopes
- Study of  $e^-$  angular distribution at SuperNEMO
- Study of Lepton Number Violation processes and search for new gauge bosons at colliders

# Conclusion

---

## Conclusion

- $0\nu\beta\beta$  will confirm the Majorana nature of neutrinos

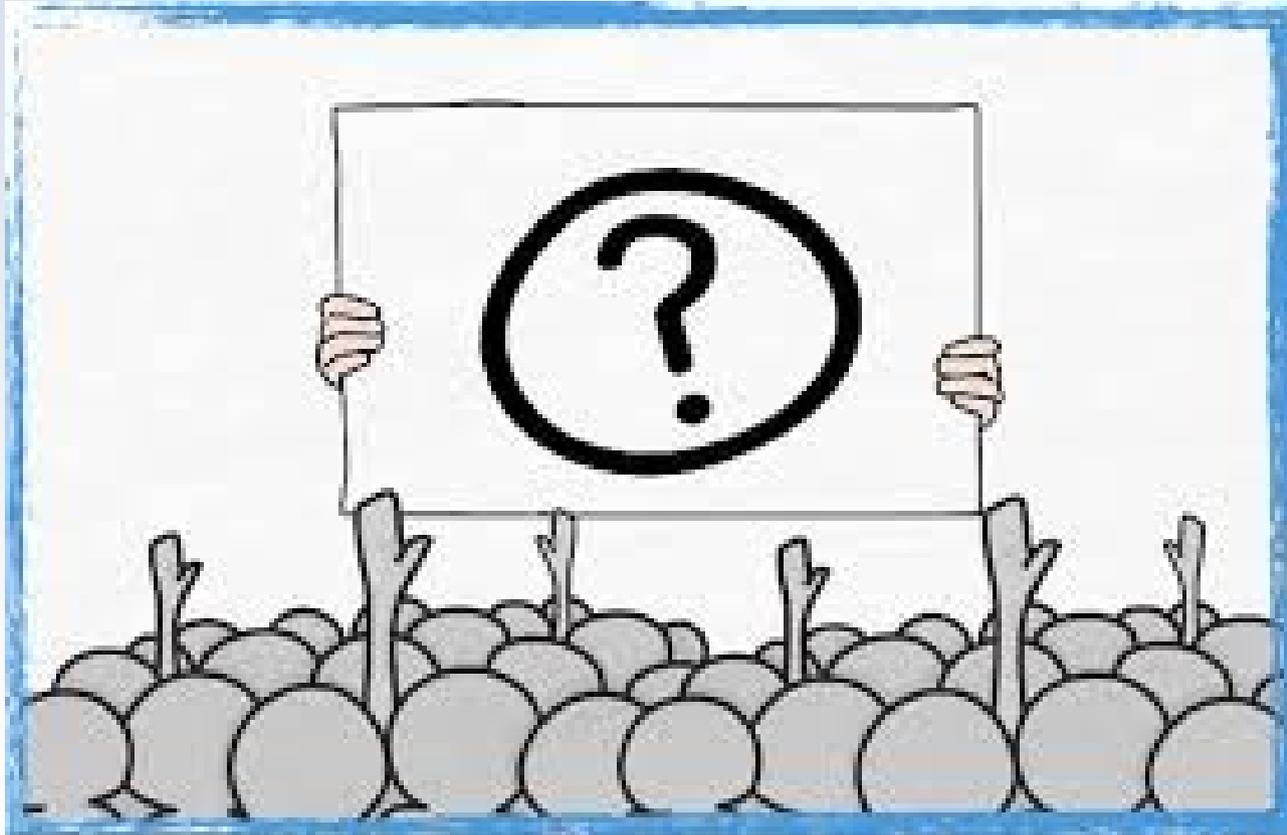
## Conclusion

- $0\nu\beta\beta$  will confirm the Majorana nature of neutrinos
- Can shine some light into BSM physics models and set bounds in new physics parameters

## Conclusion

- $0\nu\beta\beta$  will confirm the Majorana nature of neutrinos
- Can shine some light into BSM physics models and set bounds in new physics parameters
- The LRSM is an attractive extension of the SM which could be falsified at the LHC and has a clear signature in  $0\nu\beta\beta$  experiments

# Questions



- Questions??