

DF



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# FACETS OF ( $d, ^2\text{He}$ ) CHARGE-EXCHANGE REACTIONS: from few-body physics to astrophysics to double beta decay

astrophysics !

double-beta decay

NN-studies

halo nuclei

stretched states

# (d, $^2\text{He}$ ) research potential

Astrophysics

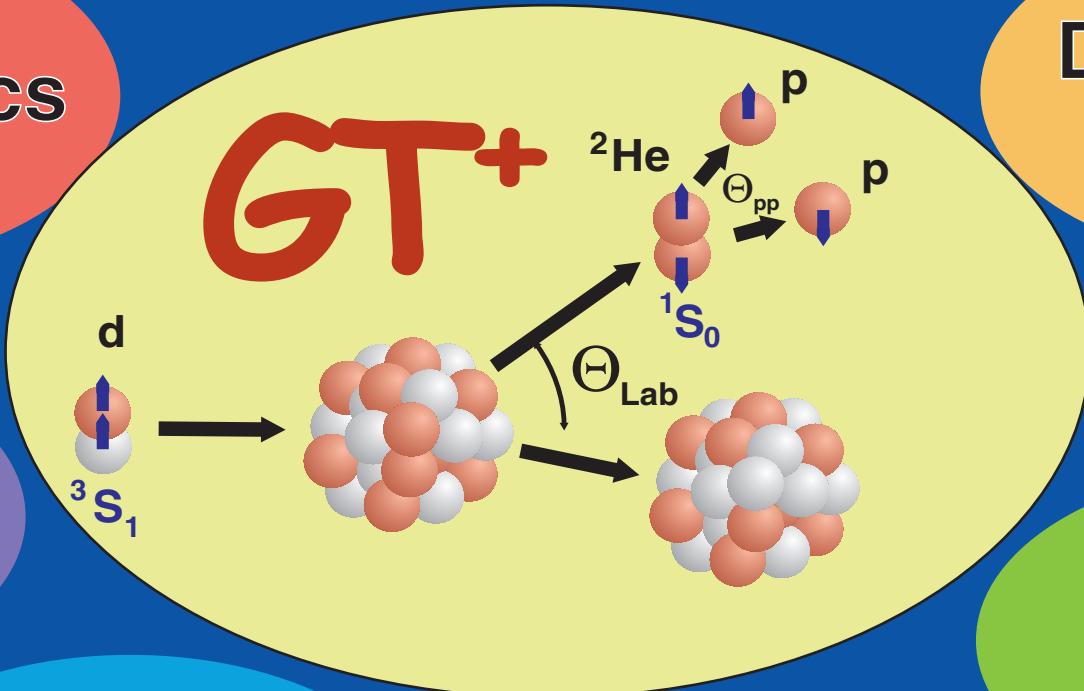
Isospin symmetry

Quantum entanglement

Halo nuclei

Double beta decay

Few-body ( $a_{nn}$ )

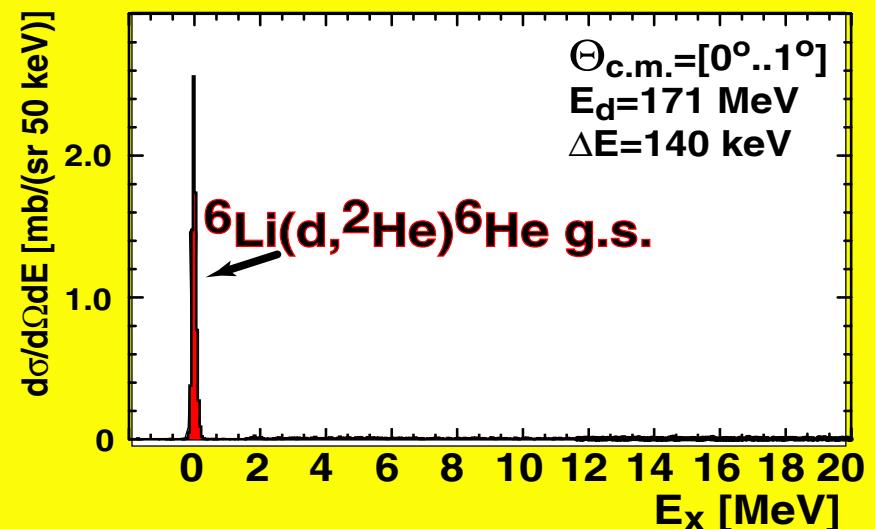
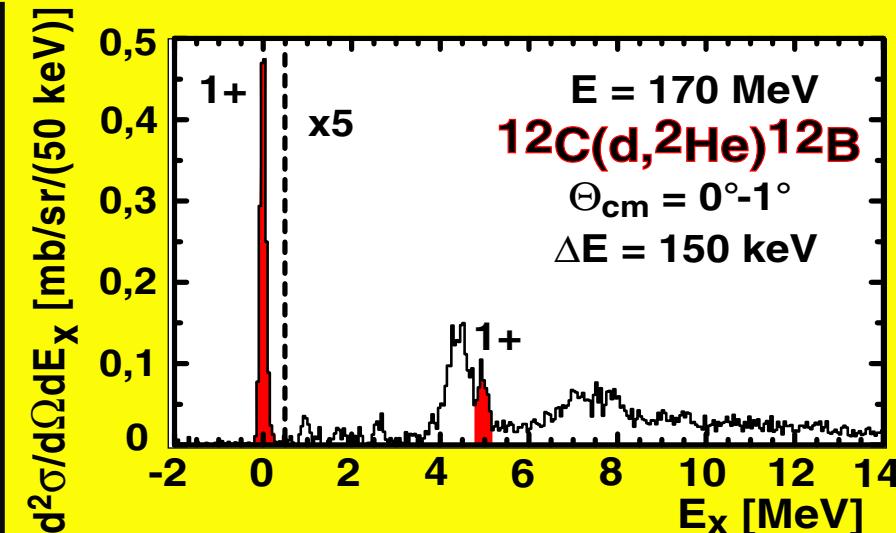
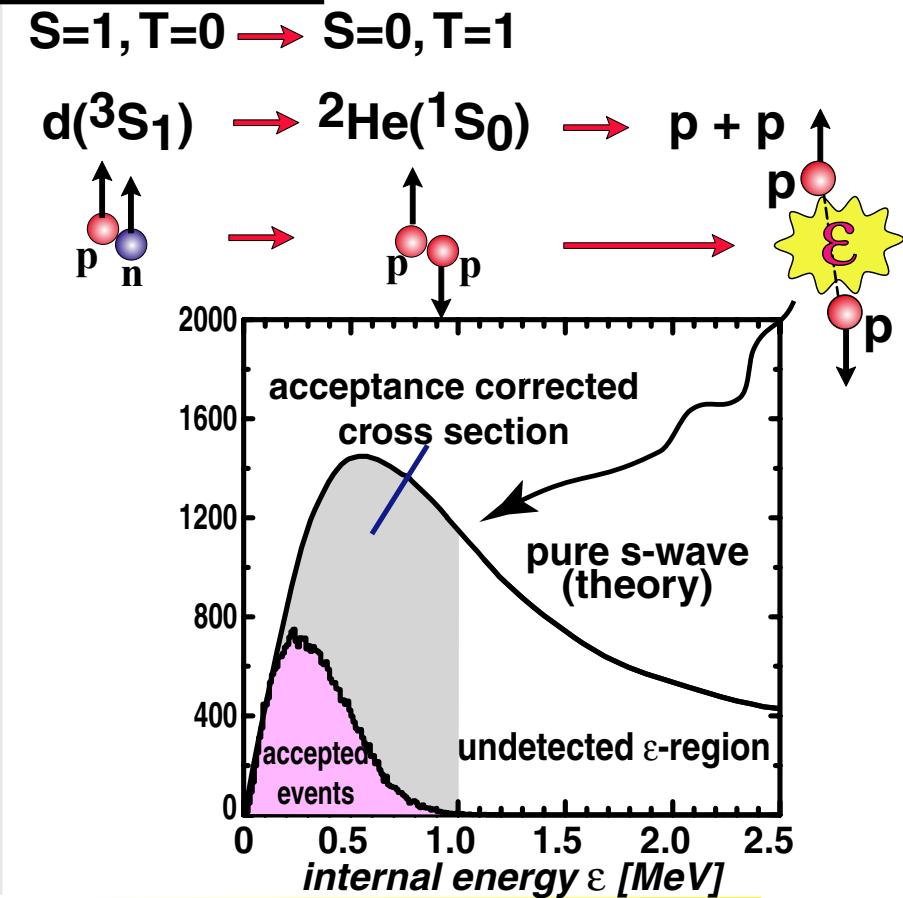


# The ( $d, {}^2\text{He}$ ) reaction

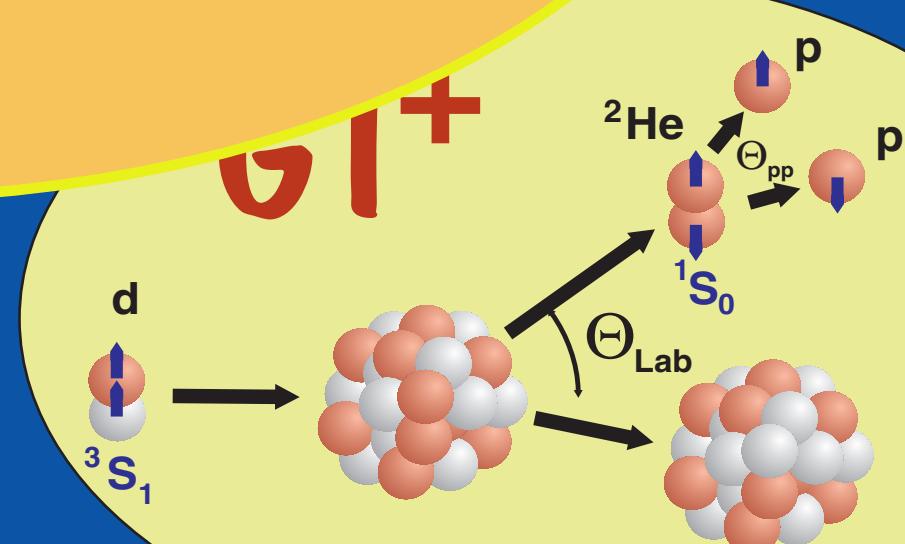
- 1)- reaction mechanism forces a spin-flip and an isospin-flip !  
 $\Delta S=1, \Delta T=1$  perfect GT filter
- 2)- coincident detection of two protons from  ${}^2\text{He}$  decay  
 → background-free spectra but need large accptnc spectrometer
- 3)- contributions from higher p-p partial waves? Dont' worry!!

## Alternatives:

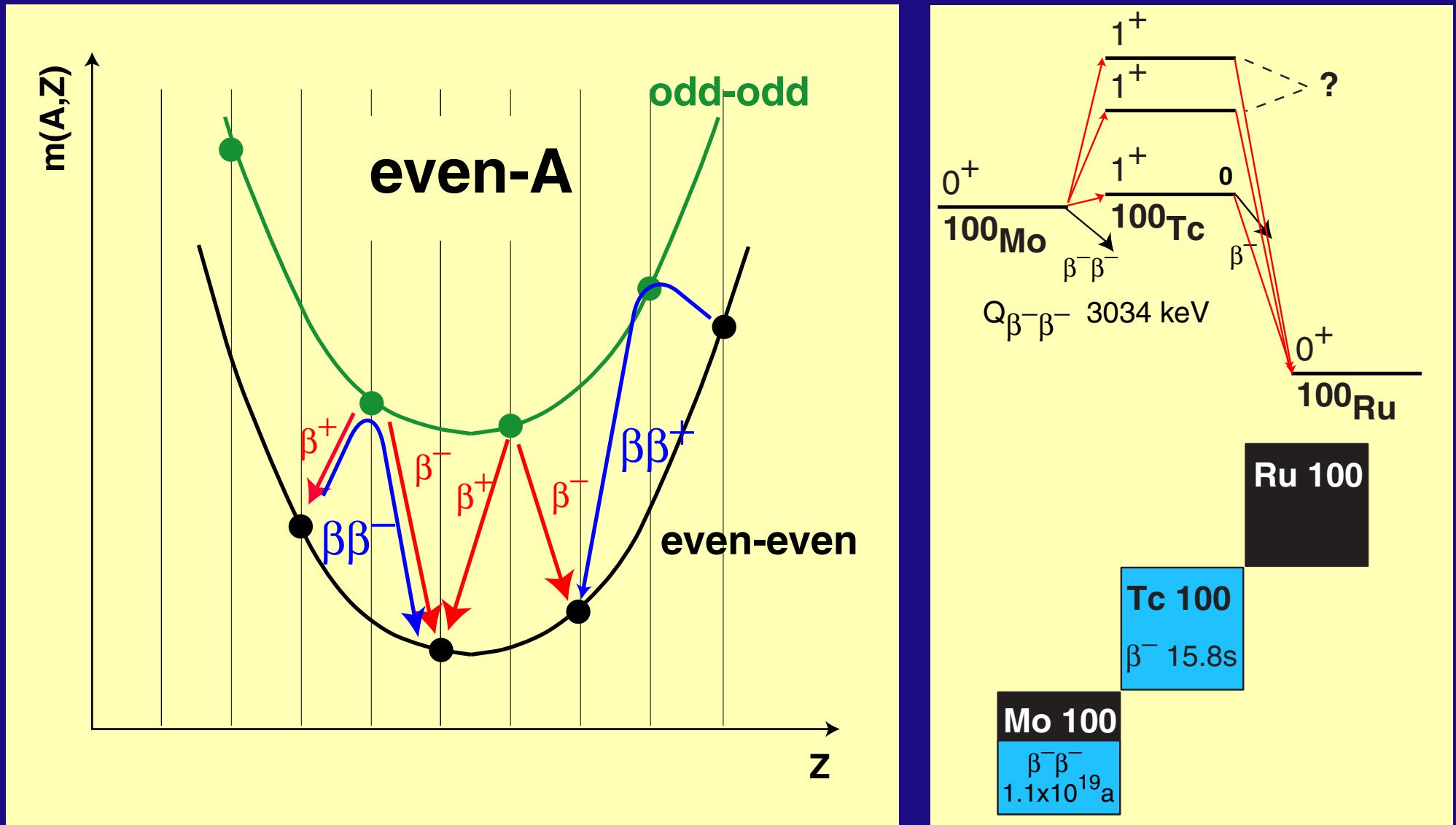
- (n,p) resolution?? Fermi transition  
 (t, ${}^3\text{He}$ ) triton beam?? Fermi transition  
 (HI,HI) resolution?? reaction mechanism??



# Double beta decay

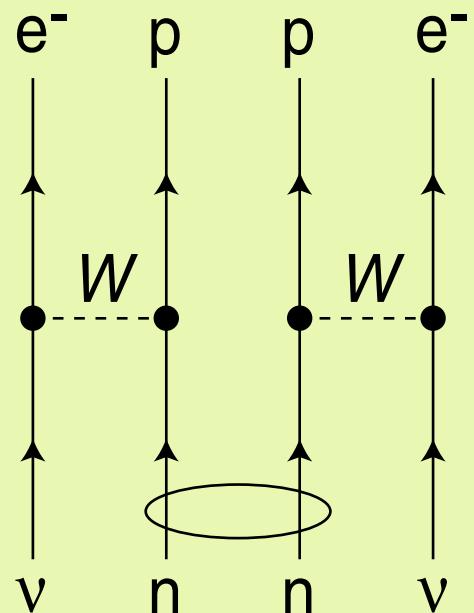


# Nuclear double beta decay



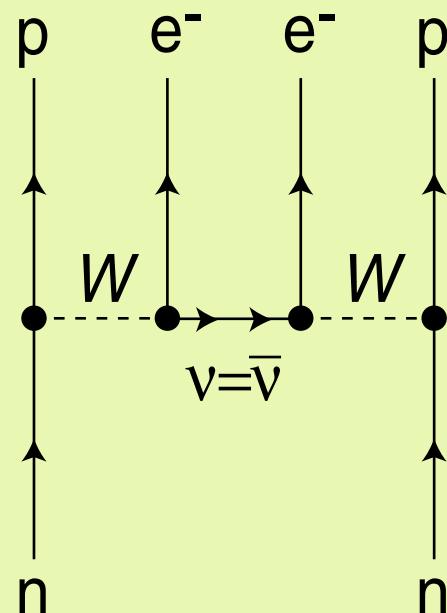
# Important $\beta\beta$ decay modes

( $2\nu\beta\beta$ )



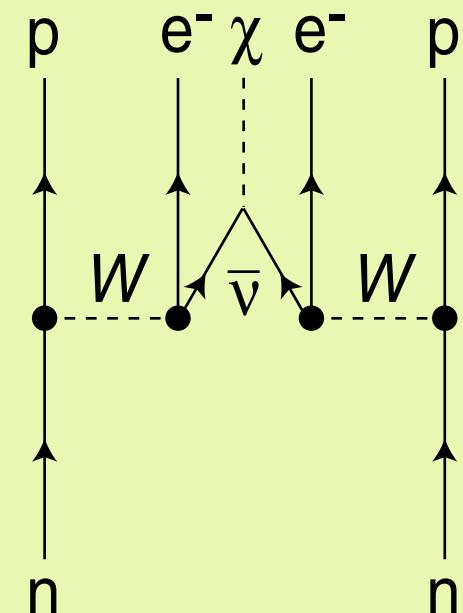
Dirac-Decay

( $0\nu\beta\beta$ )



Majorana-Decay

( $0\nu\chi\beta\beta$ )



Majoron-Decay

# 0νββ-decay: half-life & neutrino mass

$$\left[ T_{1/2}^{0\nu} (0^+ \rightarrow 0^+) \right]^{-1} = G^{0\nu} (E_0, Z) \left| M_{GT}^{0\nu} - \frac{g_V^2}{g_A^2} M_F^{0\nu} \right|^2 \langle m_\nu \rangle^2$$

measure!	look up	nuclear structure	ν mass
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$$M_{GT}^{0\nu} = \left\langle f \left| \sum_{lk} \boldsymbol{\sigma}_l \cdot \boldsymbol{\sigma}_k \tau_l^+ \tau_k^+ H(r_{lk}, \bar{A}) \right| i \right\rangle$$

$$M_F^{0\nu} = \left\langle f \left| \sum_{lk} \tau_l^+ \tau_k^+ H(r_{lk}, \bar{A}) \right| i \right\rangle$$

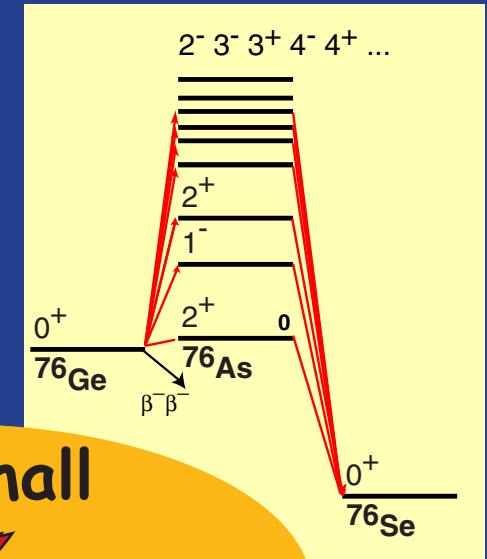
Neutrino potential  
(ν's don't escape  
from nucleus)

# 0νββ-decay: half-life & neutrino mass

$$M_{GT}^{0\nu} = \left\langle f \left| \sum_{lk} \sigma_l \cdot \sigma_k \tau_l^+ \tau_k^+ H(r_{lk}, \bar{A}) \right| i \right\rangle$$

$$M_F^{0\nu} = \left\langle f \left| \sum_{lk} \tau_l^+ \tau_k^+ H(r_{lk}, \bar{A}) \right| i \right\rangle$$

Neutrino potential



Expand expression with  $H(r, A)$

Many higher multipoles contribute!

$r_{lk}$  small  
q large ( $0.5 \text{ fm}^{-1}$ )

$$[t_{1/2}^{(0\nu)}]^{-1} = G^{(0\nu)} | M_{GT}^{(0\nu)} |^2 \langle m_\nu \rangle^2 + \text{Fermi contribution}$$

$$= G^{(0\nu)} \left| \sum_m \frac{\langle 0_{\text{g.s.}}^{(f)} || O_{\sigma\tau^-}(r, S, L) || J_m^\pi \rangle \langle J_m^\pi || O_{\sigma\tau^-}(r, S, L) || 0_{\text{g.s.}}^{(i)} \rangle}{1/2 Q_{\beta\beta}(0_{\text{g.s.}}^{(f)}) + E(J_m^\pi) - E_0} \right|^2 \langle m_\nu \rangle^2 + \text{Fermi contribution}$$

Easier case:

$2\nu\beta\beta$

# Half-lives & Matrix elements

Half life:

$$[t_{1/2}^{(2\nu)}]^{-1} = G^{(2\nu)} |M_{DGT}^{(2\nu)}|^2$$

Phase space + coupling constants:  $G \sim Q^{-11}$

$\beta\beta$  matrix element:

$$M_{DGT} = \sum_m \frac{\langle 0_{g.s.}^{(f)} || \sigma \tau^- || 1_m^+ \rangle \langle 1_m^+ || \sigma \tau^- || 0_{g.s.}^{(i)} \rangle}{1/2 Q_{\beta\beta}(0_{g.s.}^{(f)}) + E(1_m^+) - M_i}$$

All  $1^+$  levels must be considered!

Approximation:

$$M_{DGT} \approx \frac{M_S |M_{S'}|}{\Delta S}$$

$M_S$ : Single beta decay matrix elements

$\Delta S$ : Energy denominator

holds if

- only one strong  $1^+$  intermediate state
- further excited states weak or  $E_x$  high

# Measurement of $M_{DGT}^{(2v)}$ thru hadronic probes

$$M_{DGT} = \sum_m \frac{\langle 0_{g.s.}^{(f)} || \sigma\tau^- || 1_m^+ \rangle \langle 1_m^+ || \sigma\tau^- || 0_{g.s.}^{(i)} \rangle}{1/2 Q_{\beta\beta}(0_{g.s.}^{(f)}) + E(1_m^+) - M_i}$$
$$= \sum_m \frac{M_m^{GT+} M_m^{GT-}}{1/2 Q_{\beta\beta}(0_{g.s.}^{(f)}) + E(1_m^+) - M_i}$$

Measure  $B(GT+)$  through  $(n,p)$ -type reactions

Measure  $B(GT-)$  through  $(p,n)$ -type reactions

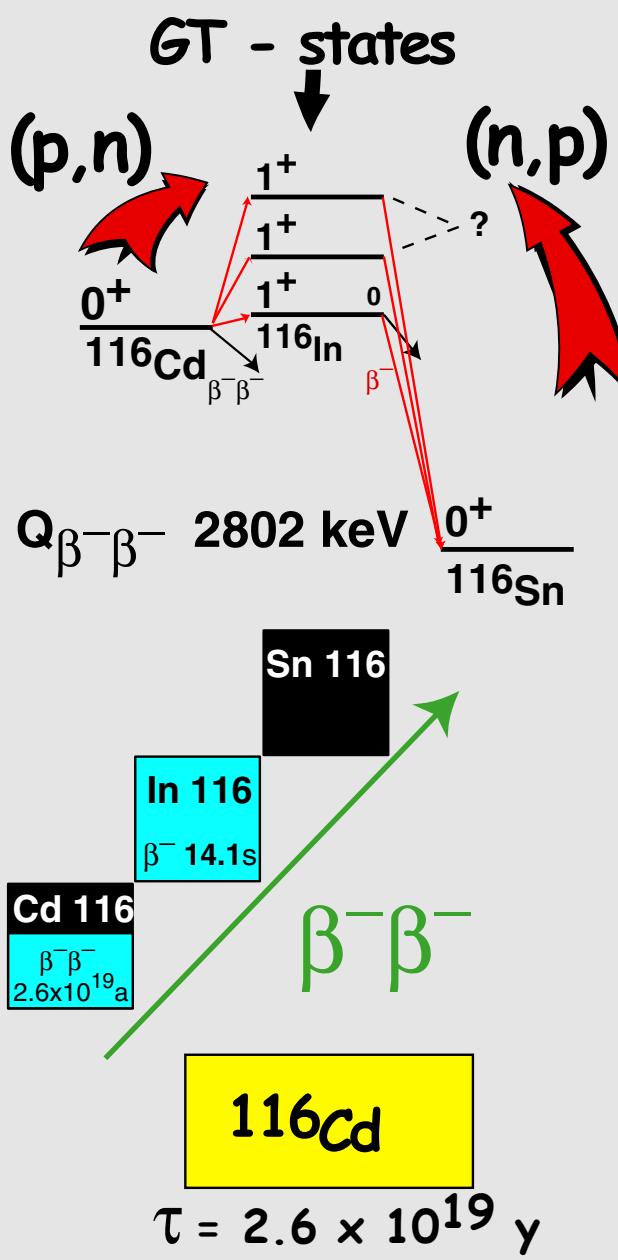
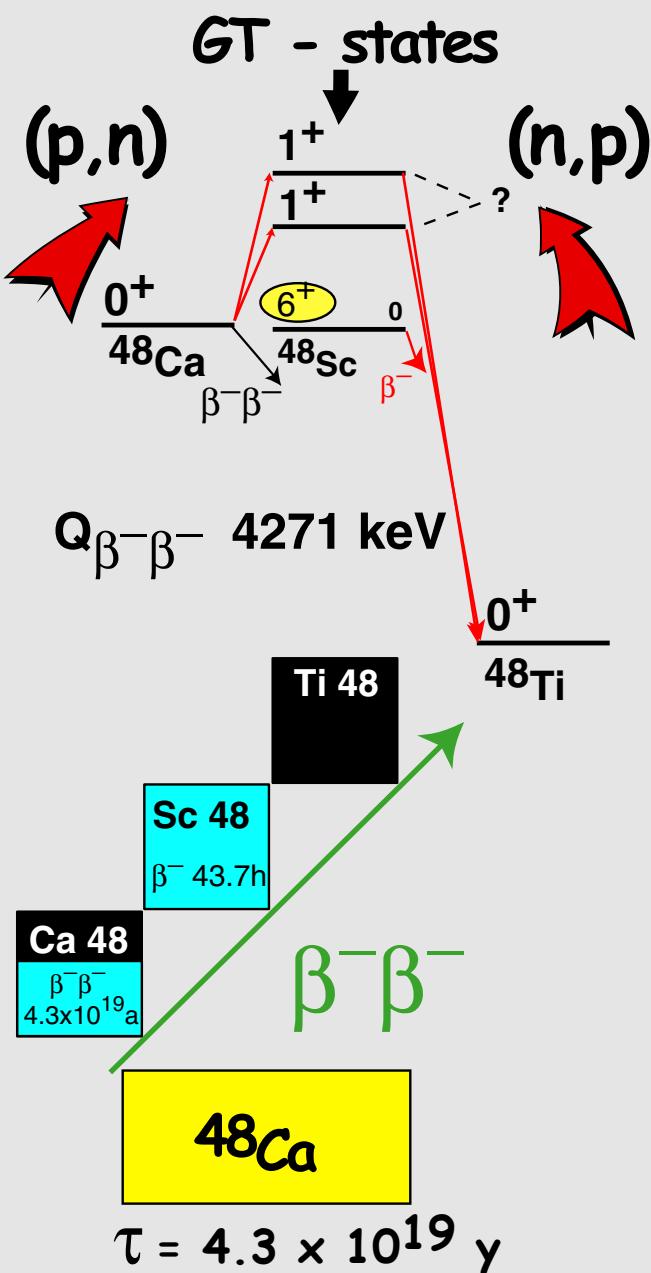
$$B(GT) = \frac{1}{2J_i + 1} | M(GT) |^2$$

- Phase cannot be measured
- Simple relation  $\sigma \leftrightarrow B(GT)$
- Little model dependence

$$B(GT) = \hat{\sigma}(GT) \frac{d\sigma(q=0)}{d\Omega}$$

forward  
angles

# The $2\nu$ double- $\beta$ decay



$\tau$  from counting experiments and as 2nd order weak process ( $\beta^- \rightarrow \beta^-$ ) !!!

Half life:

$$[t_{1/2}]^{-1} = G^{(2\nu)} | M_{DGT} |^2$$

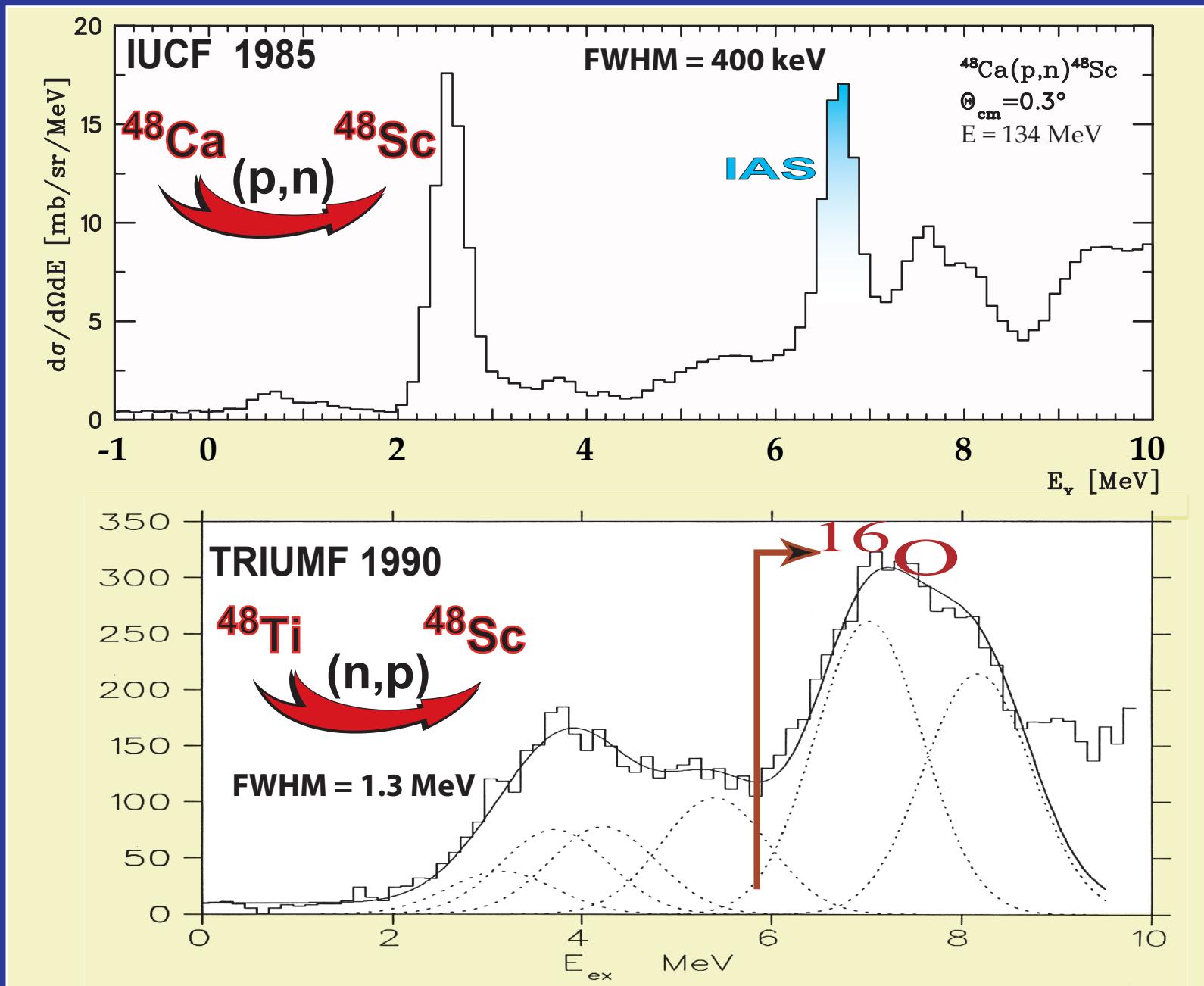
$M_{DGT} =$

$$\sum_m \frac{\langle 0_{g.s.}^{(f)} || \sigma \tau^- || 1_m^+ \rangle \langle 1_m^+ || \sigma \tau^- || 0_{g.s.}^{(i)} \rangle}{1/2 Q_{\beta\beta}(0_{g.s.}^{(f)}) + E(1_m^+) - E_0}$$

$$G^{(2\nu)} \sim (Q_{\beta\beta})^{11}$$

matrix elements available thru  $(p,n)$  and  $(n,p)$  type reactions

# $^{48}\text{Ca} - ^{48}\text{Sc} - ^{48}\text{Ti}$

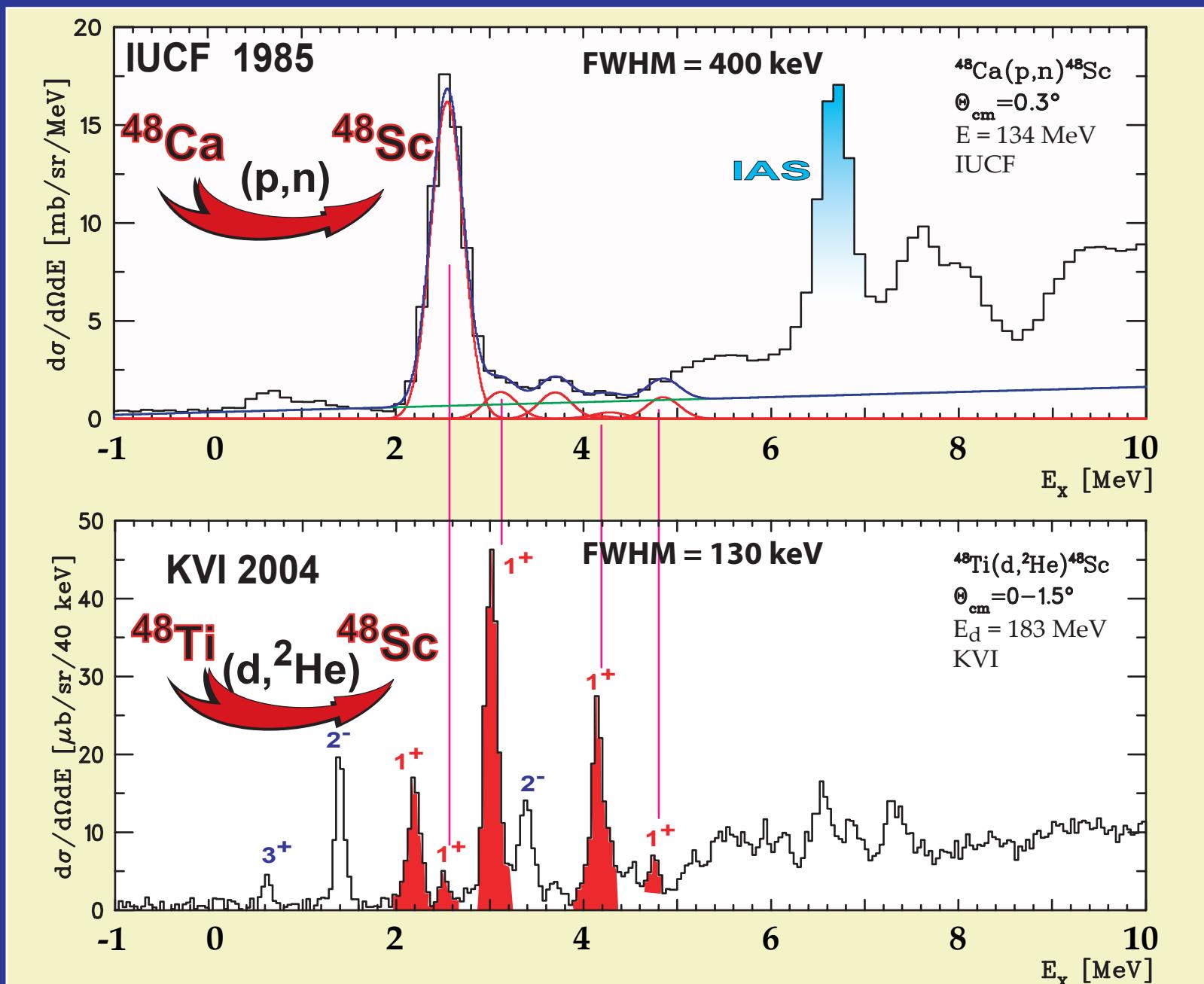


(p,n)

(n,p)

How to connect  
these states ??

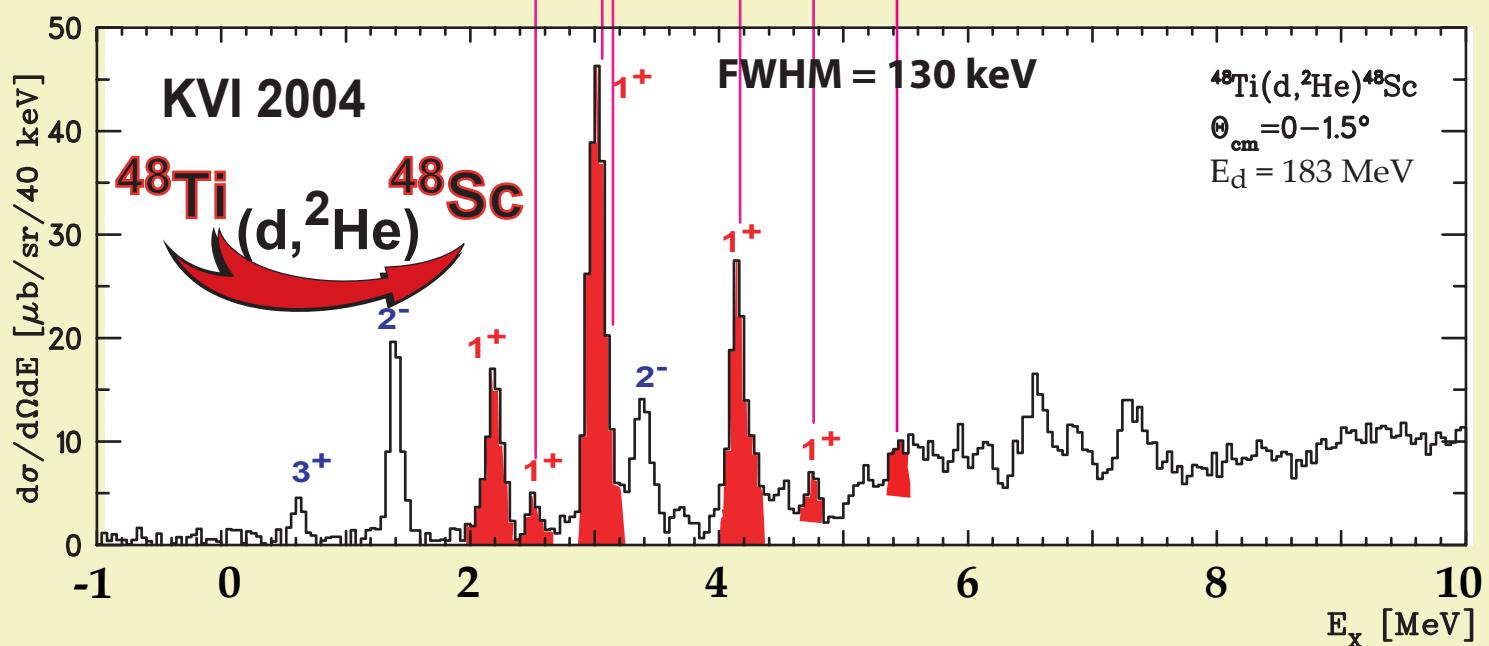
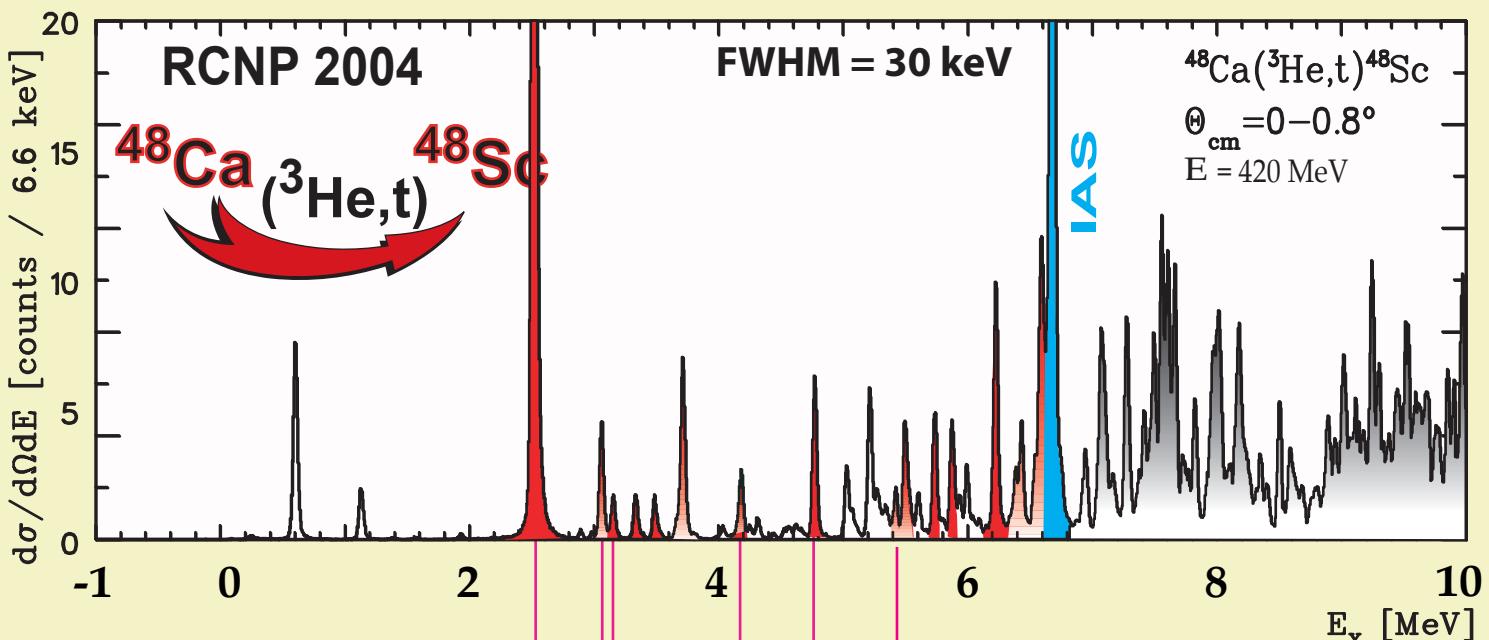
# $^{48}\text{Ca} - ^{48}\text{Sc} - ^{48}\text{Ti}$



(p,n)

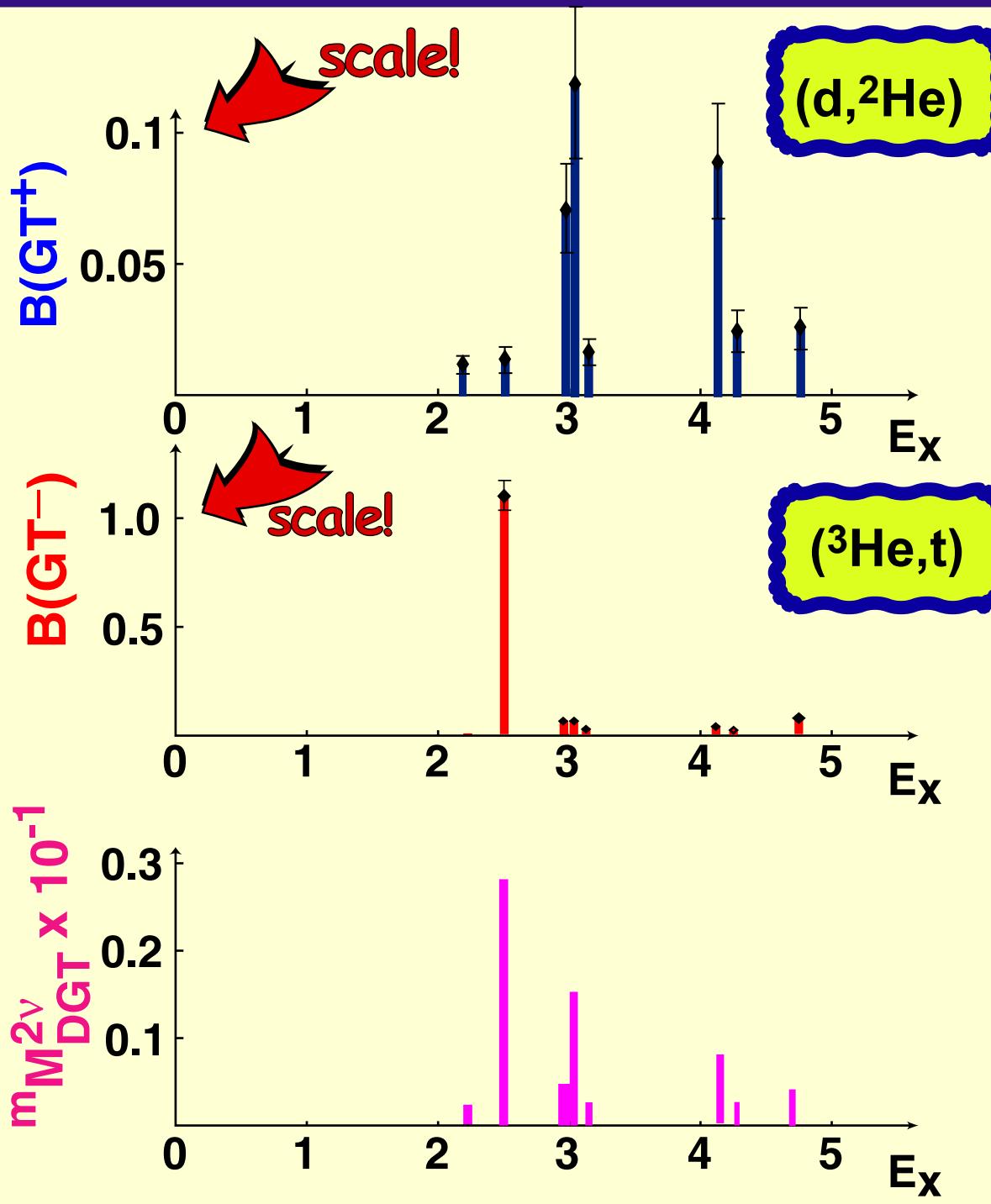
(d, $^2\text{He}$ )

# $^{48}\text{Ca} - ^{48}\text{Sc} - ^{48}\text{Ti}$



( ${}^3\text{He}, t$ )

(d,  ${}^2\text{He}$ )



## Experimental matrix elements

$$M_{DGT} = \sum_m m M_{DGT}/E_m$$

$$= 0.074 \pm 0.0155$$

all positive

$$T_{1/2} = (1.66 \pm 0.31) \times 10^{19} \text{ yr}$$

Compare to counting exp't:

$$T_{1/2} = (4.3 \pm 2.5) \times 10^{19} \text{ yr}$$

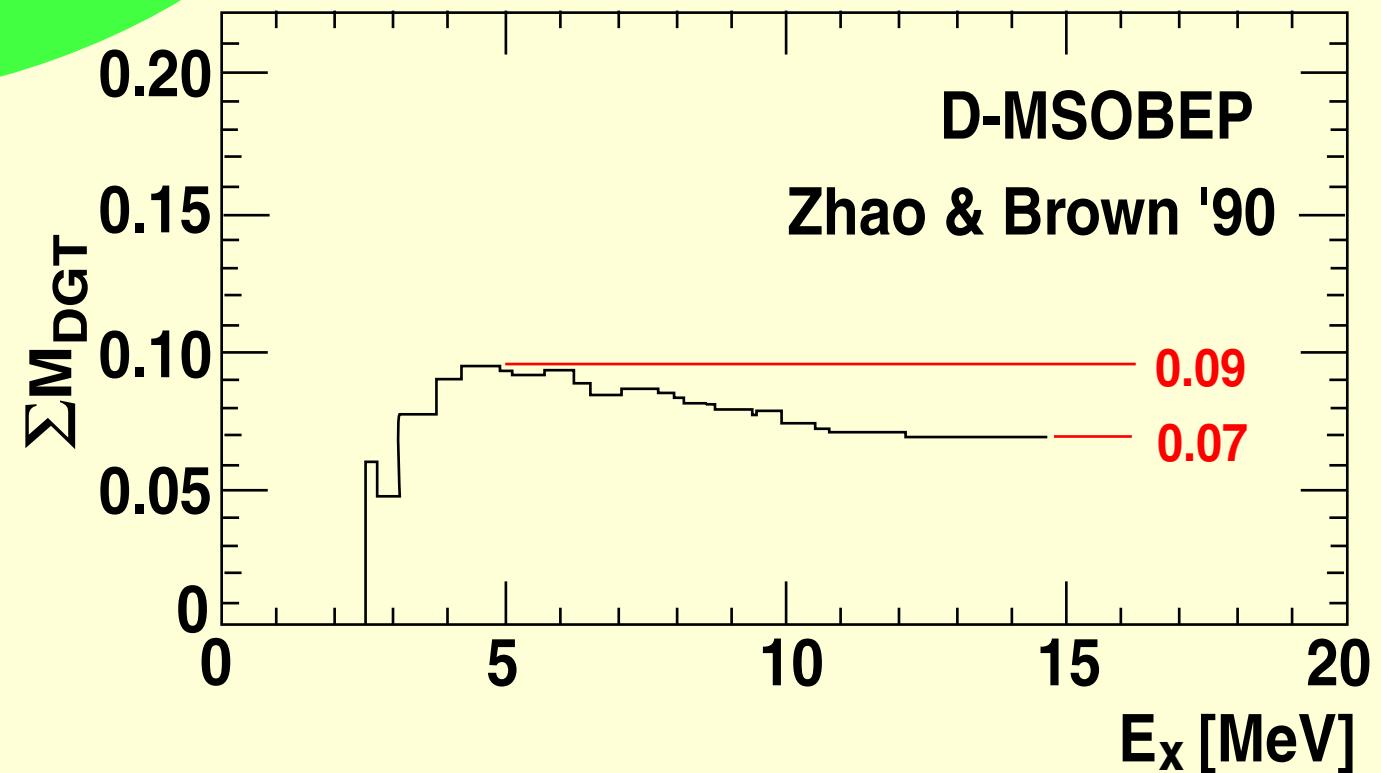
# Higher lying states ( $E_x > 5$ MeV)

SHELL MODEL

Reduction of  
matrix element  
by ~ 25%



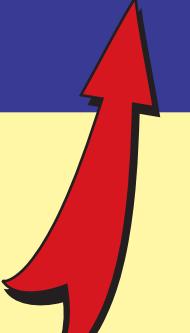
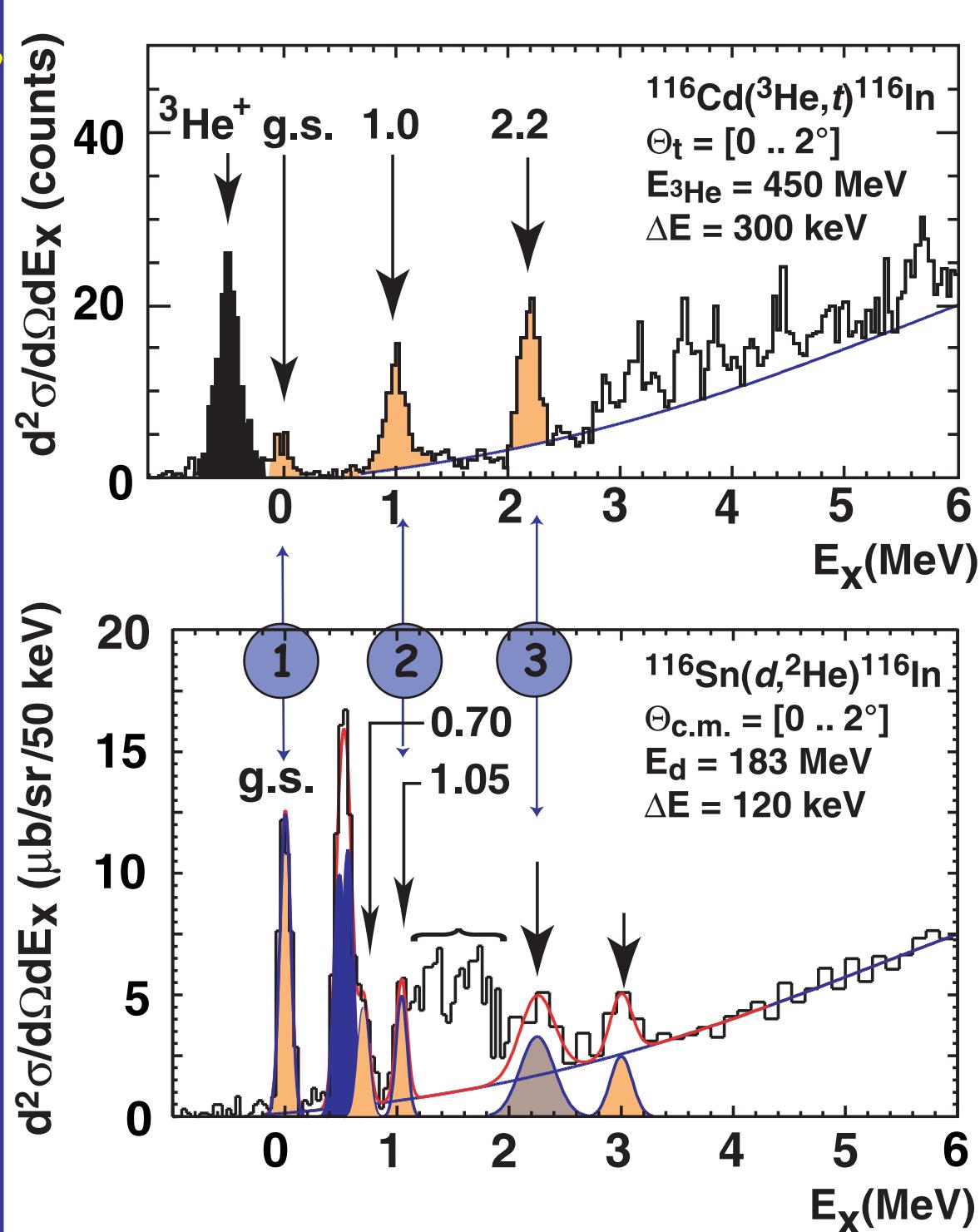
$$T_{1/2} = 3.1 \times 10^{19} \text{ yr}$$



# $^{116}\text{Cd}$ $2\nu\beta\beta$ decay

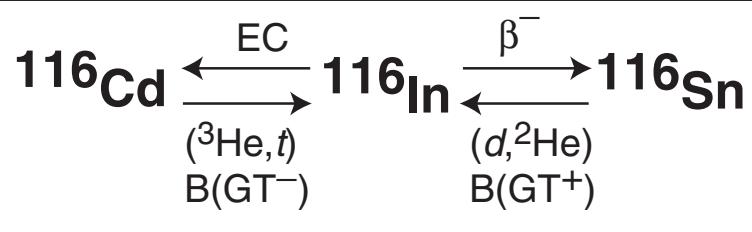
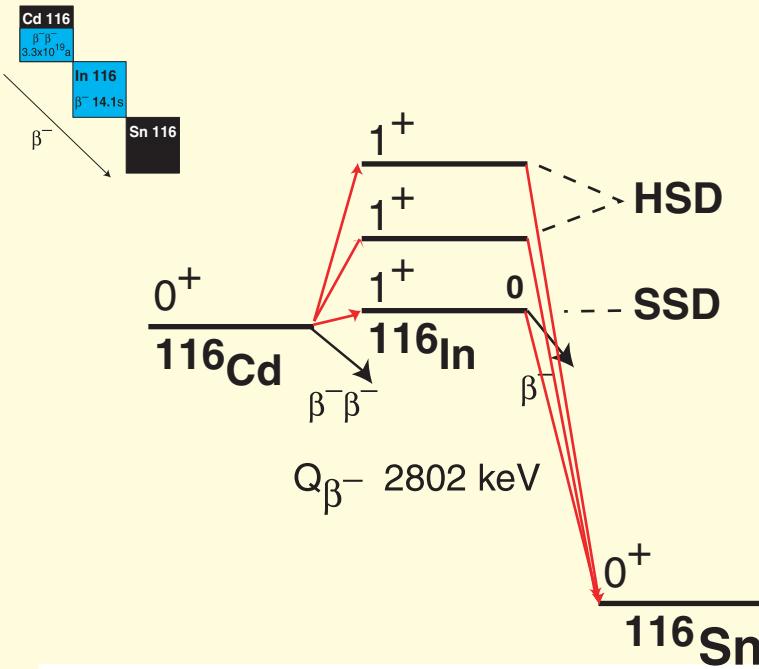
	$B(GT^-)$	$B(GT^+)$	$M_{DGT}^m$	$\Sigma M_{DGT}$ running sum
1	0.032	0.256	0.025	0.025
2	0.12	0.11	0.020	0.045
3	0.17	0.07	0.013	0.058

Matrix element from  
counting experiment:  
 $\Sigma M_{DGT} = 0.064 \pm 0.007$

# Single state dominance and its oddities

## the conjecture



## the oddity

Case	B(GT <sup>-</sup> )	B(GT <sup>+</sup> )	M(DGT)	$T_{1/2}^{(2v)} [10^{19} \text{ y}]$
direct	-	-	0.06	3.3
$(^3\text{He}, t)/\beta^-$	0.032	0.256	0.025	22
EC/ $\beta^-$	0.47	0.256	0.09	1.5
theory	1.165	0.065	0.07	2.4
$(^3\text{He}, t)/(d, ^2\text{He})$	0.061*	1.09*	0.05	4.0





# Study of Gamow-Teller transitions to intermediate states in $\beta\beta$ -decay

Status:

Finished:

$^{48}\text{Ca}$

$^{116}\text{Cd}$

Proposed  
for 2004:

$^{76}\text{Ge}$

$^{96}\text{Zr}$

$^{100}\text{Mo}$

$^{150}\text{Nd}$

2005-2007:

$^{82}\text{Se}$

$^{128,130}\text{Te}$

$^{136}\text{Xe}$



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# double beta decay -- Majorana neutrino and atomic/nuclear state degeneracy

## $2\nu \beta\beta$ decay rate:

$Q$  dependence :

$Z$  dependence:

nuclear strucure:

$Q^{10}$  (ordinary phase space)

high- $Z$  quenching (Coulomb effect)

GT-quenching, Pauli-blocking at high  $A$   
( $N-Z$  large)

state mismatch (re:  $^{48}\text{Ca}!!$ )

## $0\nu \beta\beta$ decay rate:

above arguments are weakened

but:

$$\sim m_\nu^2$$

theory:

needs the  $2\nu \beta\beta$  decay for calculating  
the  $0\nu$  variant

experiment:

needs to detect sum energy of the  
2 electrons !

are there other double beta decay variants ??

# $\beta^+\beta^+$ decay

# $\beta^+EC$ decay

# ECEC decay

$2\nu \beta^+\beta^+$  decay rate:

Q dependence :

$(Q - 4m_0c^2)^{10}$  (ordinary phase space)  
energy penalty!!!

Z dependence:

high-Z enhancement (Coulomb effect)

nuclear strucure:

much reduced Pauli-blocking at high A

(N-Z small)

$2\nu$  or  $0\nu$  experiment:

very ! low isotopic abundances  
vanishing decay rates

4 gamma's plus two beta's

not much different

except energy penalty is only  $2m_0c^2$

$\beta^+EC$  case:

hopeless

# ECEC decay

2v ECEC decay rate:

Q dependence : moderate (ordinary phase space)  
no energy penalty!!!

Z dependence:  $\sim (\alpha Z)^6$  electron density!!

nuclear strucure: much reduced Pauli-blocking at high A  
(N-Z small)

## hopeless

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can never detect the X-rays (Auger-electrons)  
estimated lifetime  $\sim 10^{30-32}$  y

0v ECEC decay rates: need an extra photon to carry away energy

(0v $\gamma$  ECEC--X-ray )  $\sim (\alpha/Q)^2$  i.e. need low Q !!

$\sim (\alpha Z)^6$  electron density

no phase space factor

estimated lifetime  $\sim 10^{30-32}$  y

however: there is a resonant enhancement if the Q-value  
matches the atomic excitation

# $\beta^+\beta^+$ decay nuclei(6)

transition	Q-value	final states in daughter	isotopic abundance
$^{78}\text{Kr} \rightarrow ^{78}\text{Se}$	2.868	many	0.35%
$^{96}\text{Ru} \rightarrow ^{96}\text{Mo}$	2.725	many	5.52%
$^{106}\text{Cd} \rightarrow ^{106}\text{Pd}$	2.770	many	1.25%
$^{124}\text{Xe} \rightarrow ^{124}\text{Te}$	2.865	many	0.10%
$^{130}\text{Ba} \rightarrow ^{130}\text{Xe}$	2.610	many	0.11%
$^{136}\text{Ce} \rightarrow ^{136}\text{Ba}$	2.400	many	0.20%

energy penalty  $\beta^+\beta^+$  :  $4m_0c^2 = 2.044 \text{ MeV}$

energy penalty  $\beta^+ EC$  :  $2m_0c^2 = 1.022 \text{ MeV}$

energy penalty  $EC EC$ : none

decay rate  $\Gamma(\beta^+\beta^+) \sim (Q - 4m_0c^2)^{10}$

# $\beta^+ EC$ decay nuclei (13)

transition	Q-value	final states in daughter	isotopic abundance
$^{58}\text{Ni} \rightarrow ^{58}\text{Fe}$	1.924	g.s. , 0.810 (2+) , 1.674 (2+)	68.08%
$^{64}\text{Zn} \rightarrow ^{64}\text{Ni}$	1.097	g.s.	48.60%
$^{74}\text{Se} \rightarrow ^{74}\text{Ge}$	1.209	g.s. , 0.595 (2+) , 1.204 (2+)	0.89%
$^{84}\text{Sr} \rightarrow ^{84}\text{Kr}$	1.789	g.s. , 0.881 (2+)	0.56%
$^{92}\text{Mo} \rightarrow ^{92}\text{Zr}$	1.649	g.s., 0.934 (2+) , 1.382 (0+) , 1.495 (4+)	14.84%
$^{112}\text{Sn} \rightarrow ^{112}\text{Cd}$	1.923	many	0.97%
$^{120}\text{Te} \rightarrow ^{120}\text{Sn}$	1.703	g.s. , 1.171 (2+)	0.09%
$^{144}\text{Sm} \rightarrow ^{144}\text{Nd}$	1.781	g.s. , 0.696 (3-) , 1.314 (4+) , 1.510 (3-) , 1.561 (2+)	3.10%
$^{150}\text{Gd} \rightarrow ^{150}\text{Sm}$	1.290	many	$\alpha$ -decay $10^6$ y
$^{156}\text{Dy} \rightarrow ^{156}\text{Gd}$	2.010	many	0.06%
$^{162}\text{Er} \rightarrow ^{162}\text{Dy}$	1.845	many	0.14%
$^{168}\text{Yb} \rightarrow ^{168}\text{Er}$	1.422	many	0.13%
$^{174}\text{Hf} \rightarrow ^{174}\text{Yb}$	1.102	many	0.16%

energy penalty  $\beta^+ EC : 2m_0 c^2 = 1.022$  MeV

energy penalty  $EC EC$ : none -- decay rate  $\Gamma(EC EC) \sim (\alpha Z)^6$

# ECEC decay nuclei (12)

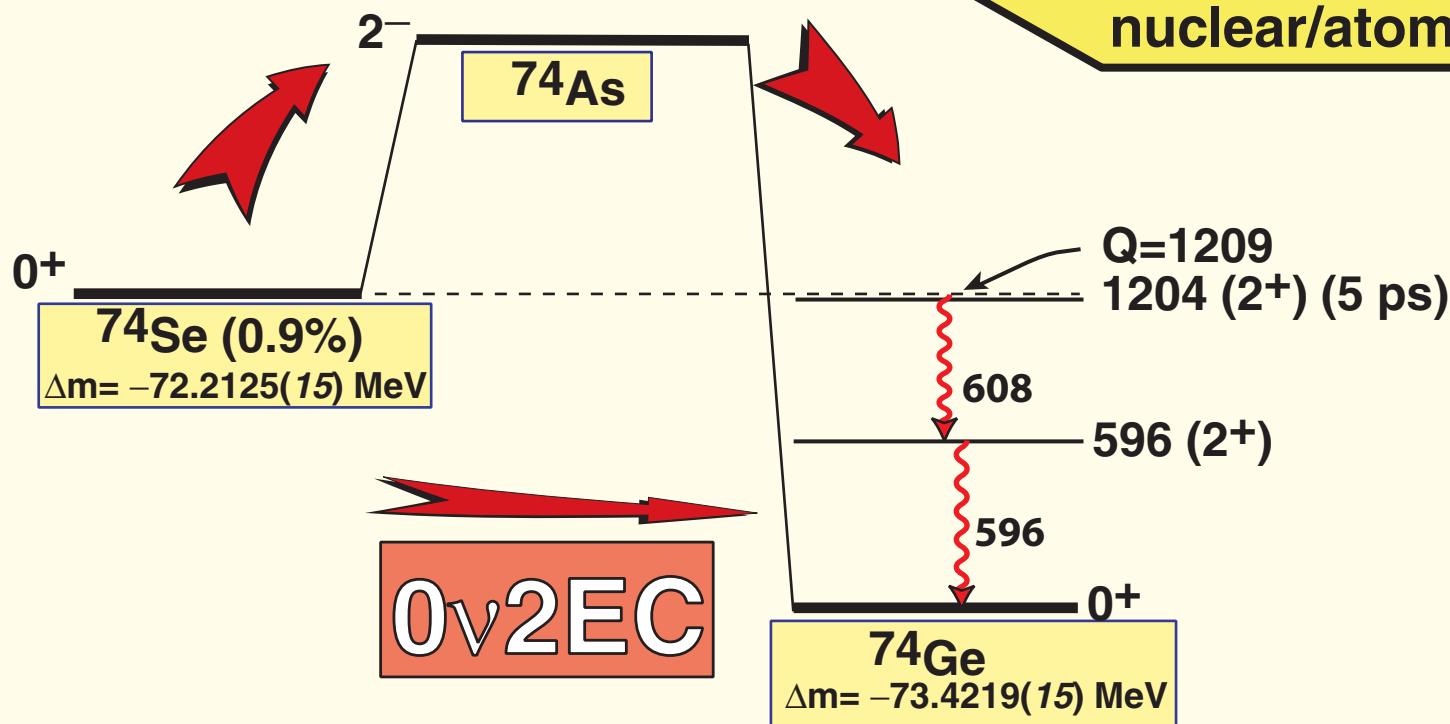
<i>transition</i>	<i>Q-value</i>	<i>final states in daughter</i>	<i>isotopic abundance</i>
$^{36}\text{Ar} \rightarrow ^{36}\text{S}$	0.434	g.s. (0+)	0.34%
$^{40}\text{Ca} \rightarrow ^{40}\text{Ar}$	0.193	g.s. (0+)	96.94%
$^{54}\text{Fe} \rightarrow ^{54}\text{Cr}$	0.681	g.s. (0+)	5.80%
$^{108}\text{Cd} \rightarrow ^{108}\text{Pd}$	0.259	g.s. (0+)	0.89%
$^{126}\text{Xe} \rightarrow ^{126}\text{Te}$	0.897	g.s. (0+) . 0.666 (2+)	0.09%
$^{132}\text{Ba} \rightarrow ^{132}\text{Xe}$	0.841	g.s. (0+) . 0.668 (2+)	0.10%
$^{138}\text{Ce} \rightarrow ^{138}\text{Ba}$	0.693	g.s. (0+)	71.70%
$^{152}\text{Gd} \rightarrow ^{152}\text{Sm}$	0.056	g.s. (0+)	0.20%
$^{158}\text{Dy} \rightarrow ^{158}\text{Gd}$	0.284	g.s. (0+) . 0.079 (2+) . 0.261 (4+)	0.10%
$^{164}\text{Er} \rightarrow ^{164}\text{Dy}$	0.024	g.s. (0+)	1.61%
$^{180}\text{W} \rightarrow ^{180}\text{Hf}$	0.145	g.s. (0+) . 0.093 (2+)	0.13%
$^{196}\text{Hg} \rightarrow ^{196}\text{Pt}$	0.820	g.s. (0+) . 0.355 (2+) . 0.689 (2+)	0.15%

energy penalty EC EC: none

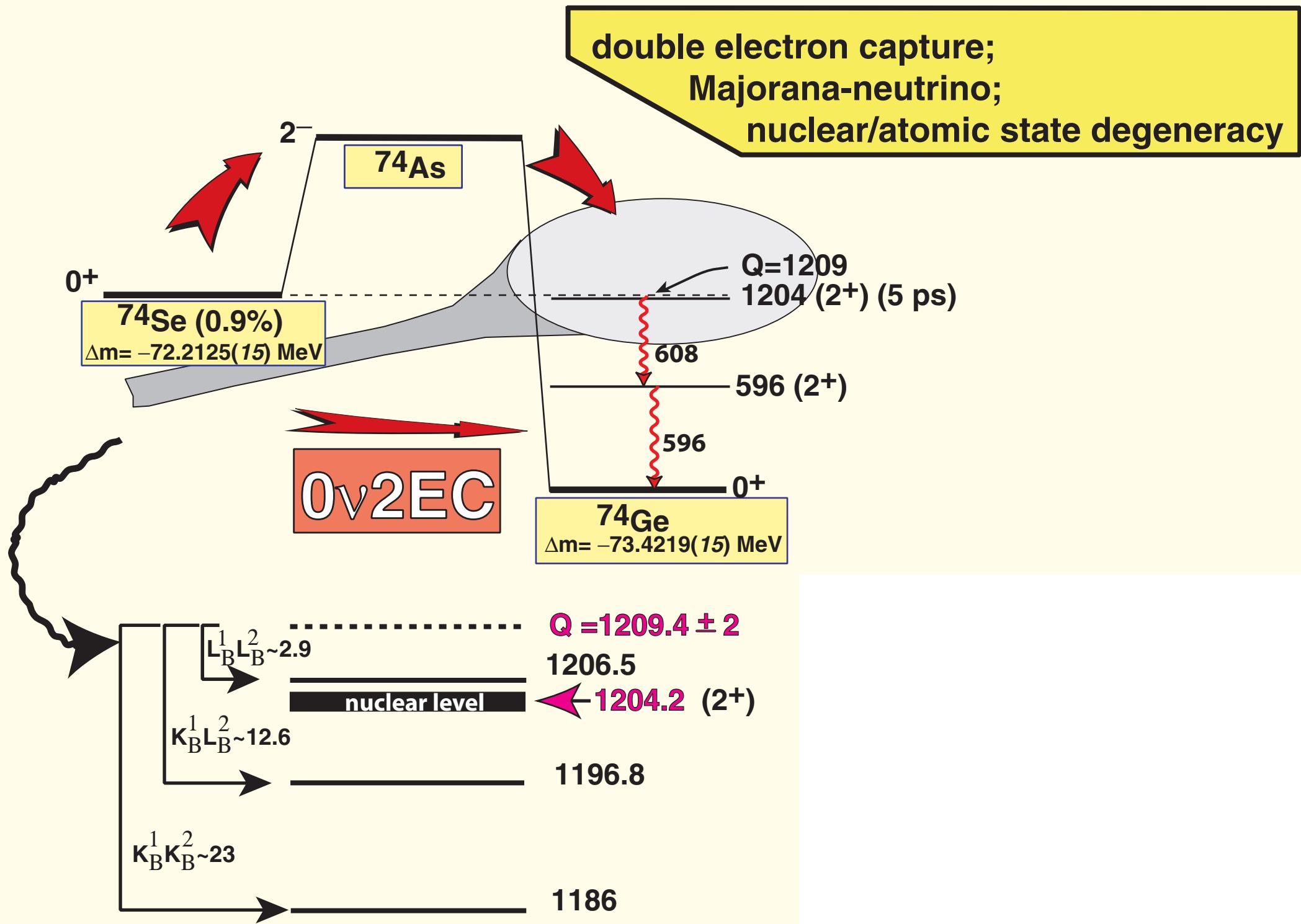
decay rate  $\Gamma \sim (\alpha Z)^6 (\alpha/Q)^2$  (e-density x photon propagator)

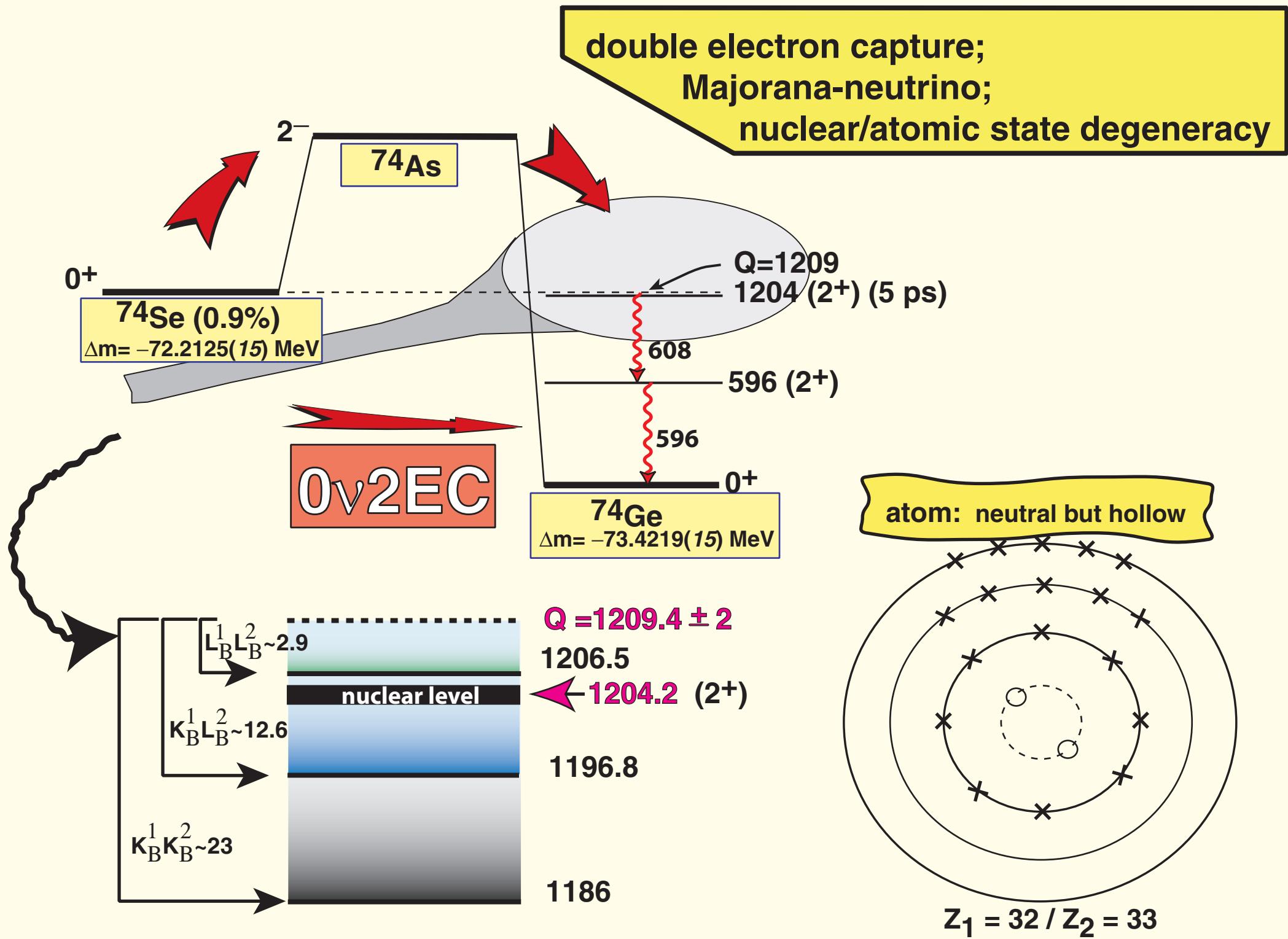
**state degeneracy and the  
special case of  $^{74}\text{Se}$  for a  
Majorana neutrino signature**

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double electron capture;  
 Majorana-neutrino;  
 nuclear/atomic state degeneracy





# **$^{74}\text{Se}$ and Majorana neutrino signature**

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estimated lifetime for  $0\nu\text{ECEC}$  decay:

$10^{20} - 10^{25} \text{ yr possible}$   
 $(0.2 - 10^4 \text{ cts/day/kmol}) \times |\text{m}_\nu(\text{eV})|^2$

need desperately theory support !!

(present calculations from S. Wycech to be refined)

# The case of $^{48}\text{Ca}$

$$T_{1/2} \sim G \sim Q^{-11}$$

$^{48}\text{Ca}$

$$\begin{aligned} \frac{T}{Q} &= 4.27 \text{ MeV} \\ T_{1/2} &= 4.3 \times 10^{19} \text{ yr} \end{aligned}$$

$^{116}\text{Cd}$

$$\begin{aligned} \frac{T}{Q} &= 2.80 \text{ MeV} \\ T_{1/2} &= 3.8 \times 10^{19} \text{ yr} \end{aligned}$$

NUCLEAR STRUCTURE

Shell-model calculations (Caurier&Poves, 1990)

