

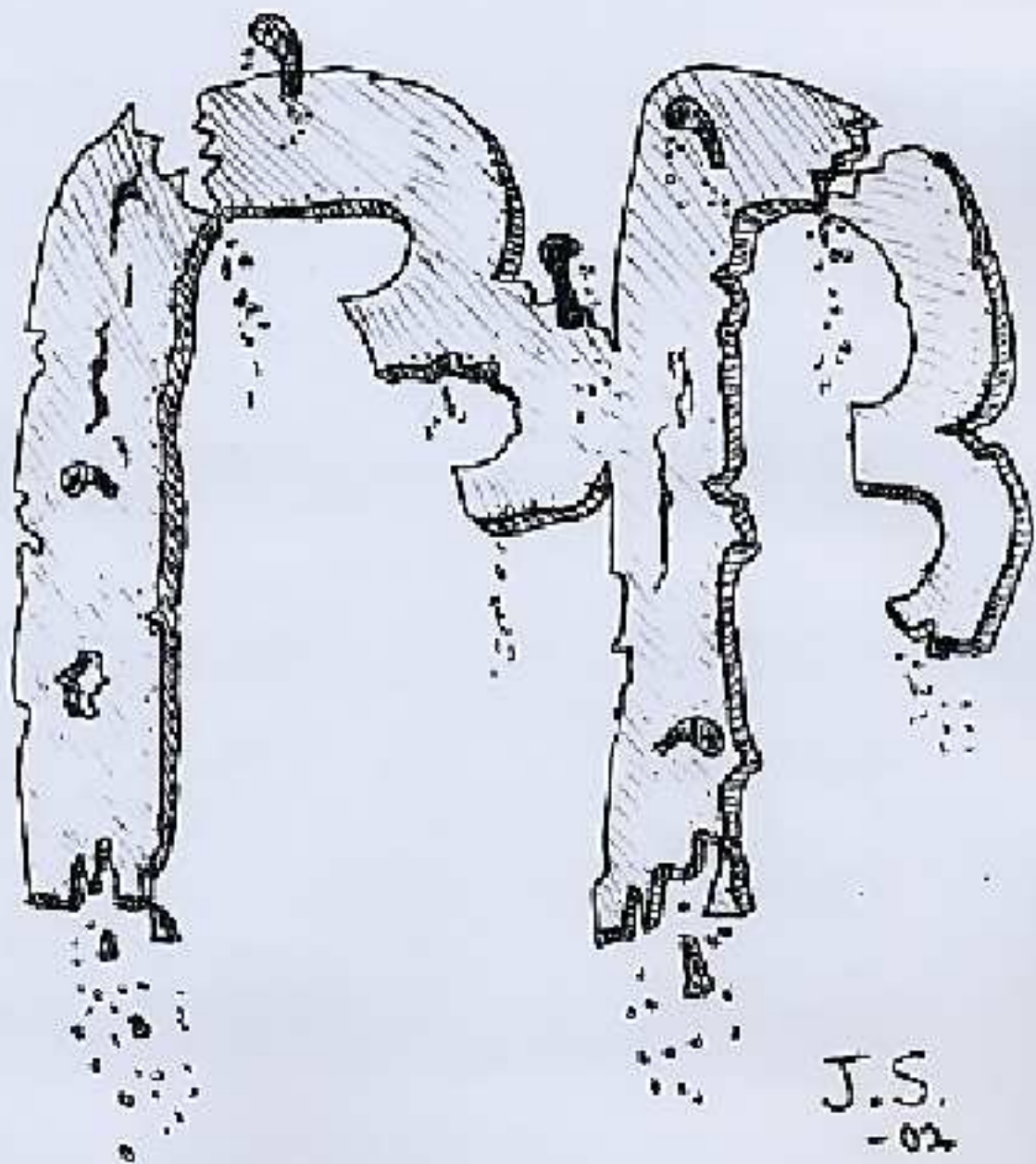
WAYS TO PROBE DOUBLE BETA MATRIX ELEMENTS BY INDEPENDENT DATA

Jouni Suhonen

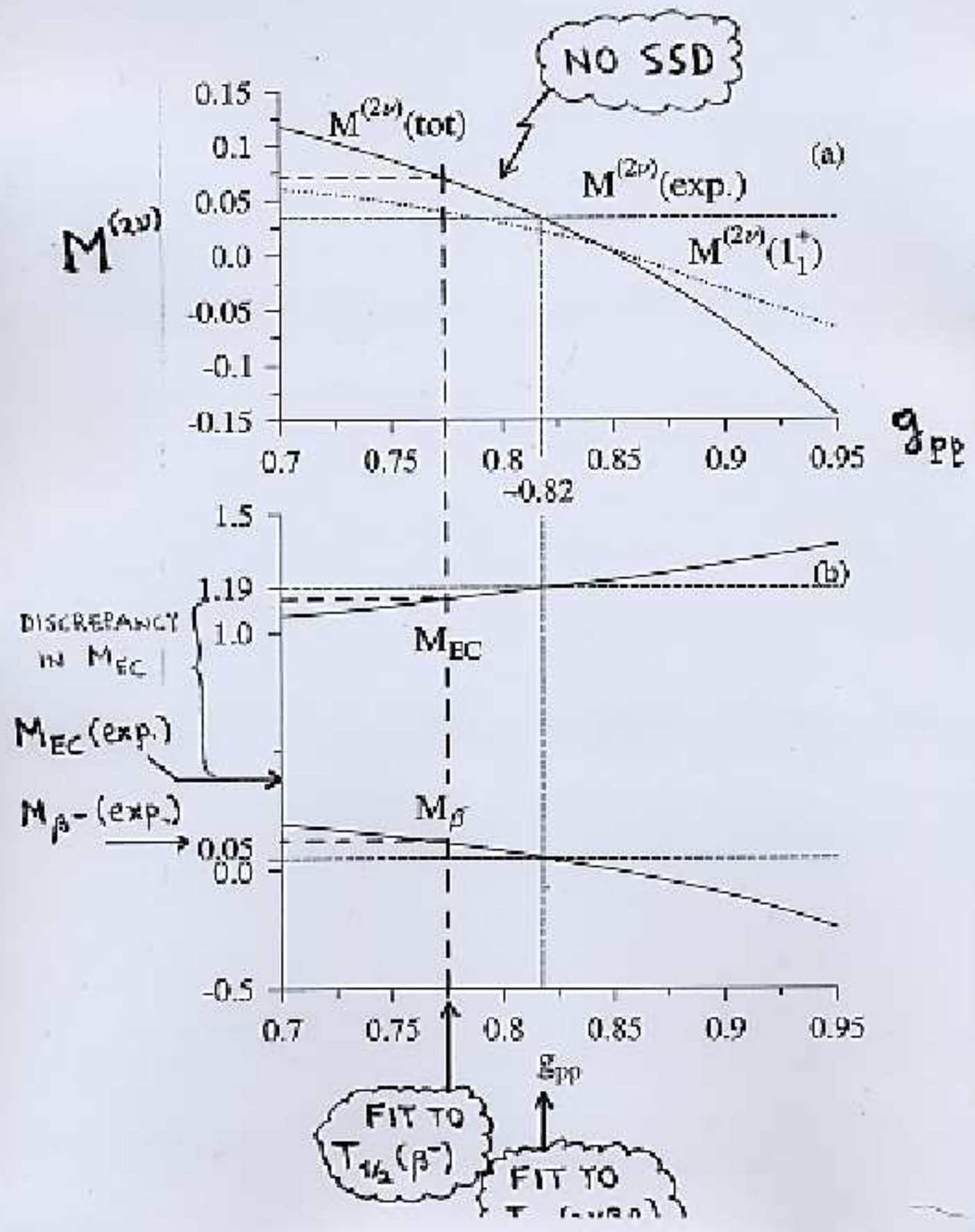
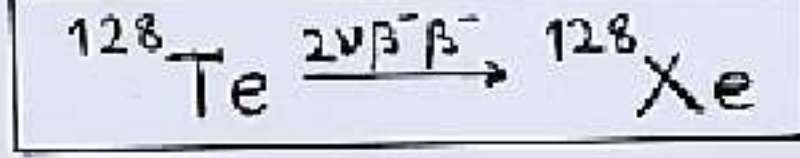
IPPP Durham, England, May 23-24,
2005

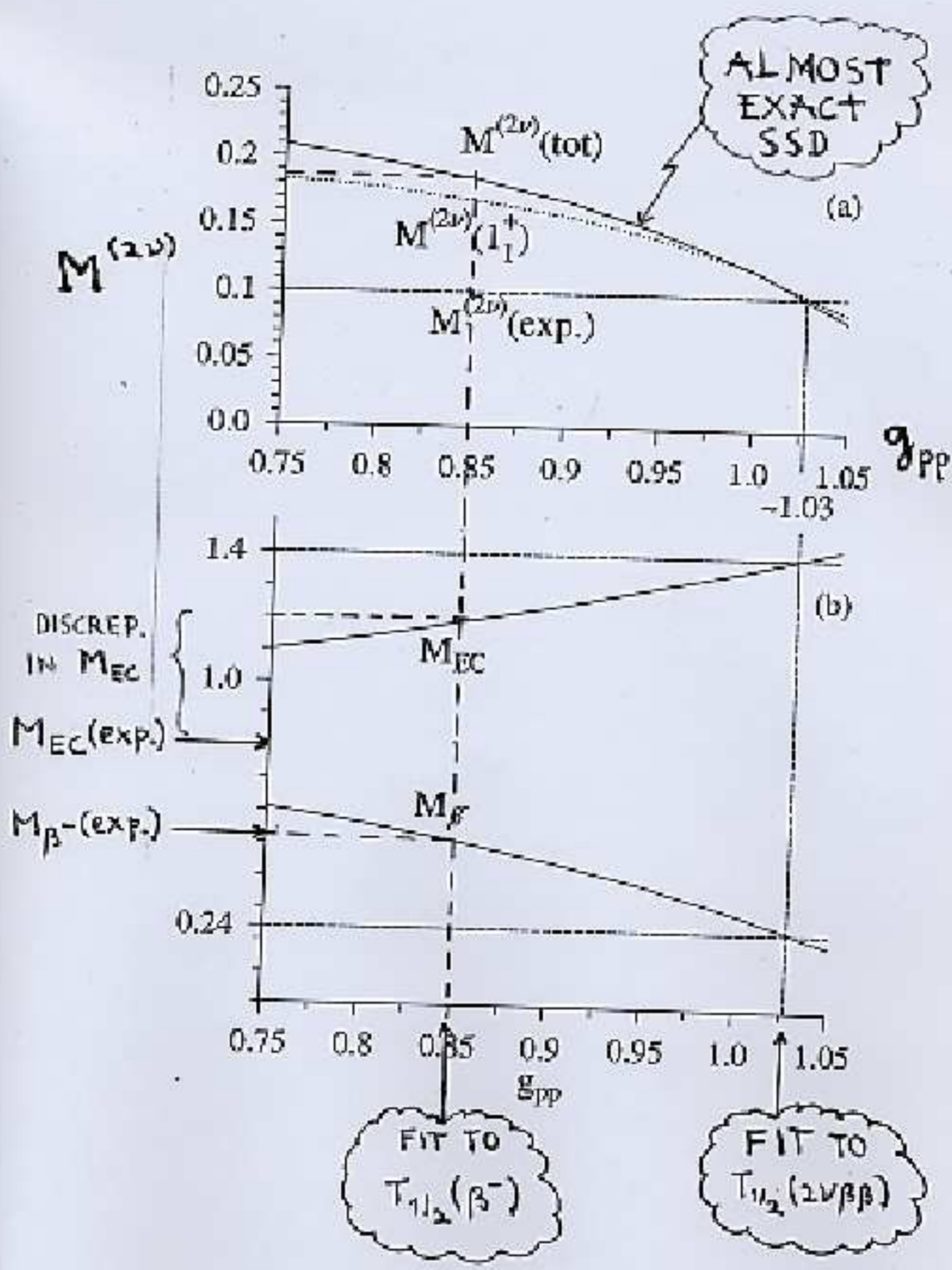
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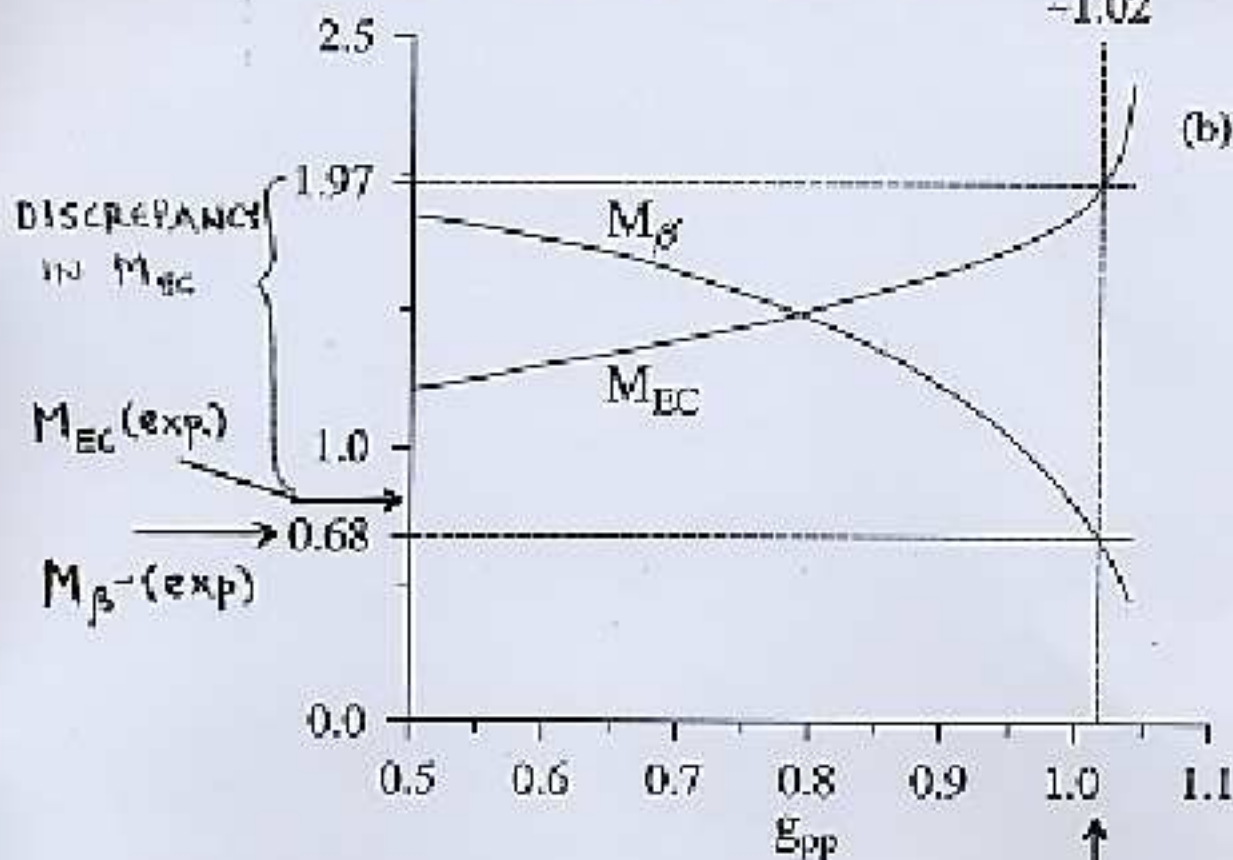
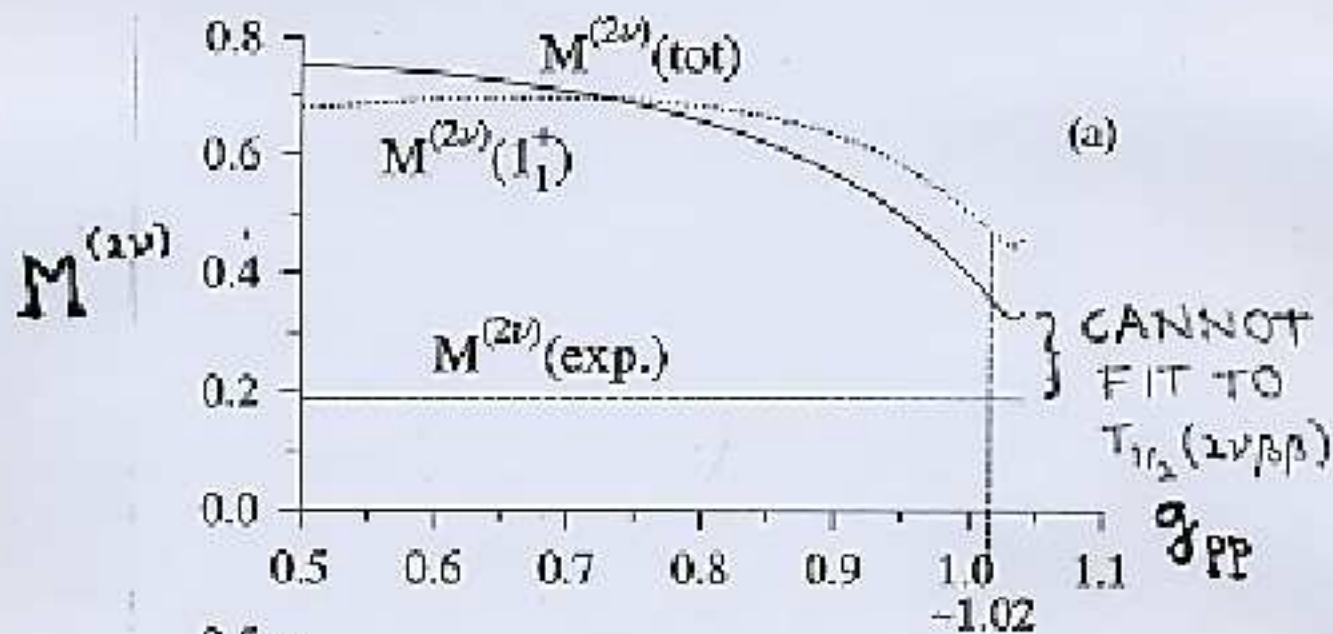
- $0\nu\beta\beta$ m.e.'s from $2\nu\beta\beta$ -
or β -decay data?
- Decomposition of the $0\nu\beta\beta$
m.e. $\leftrightarrow$ dedicated experiments
↑
change exchange
[μ^- capture]

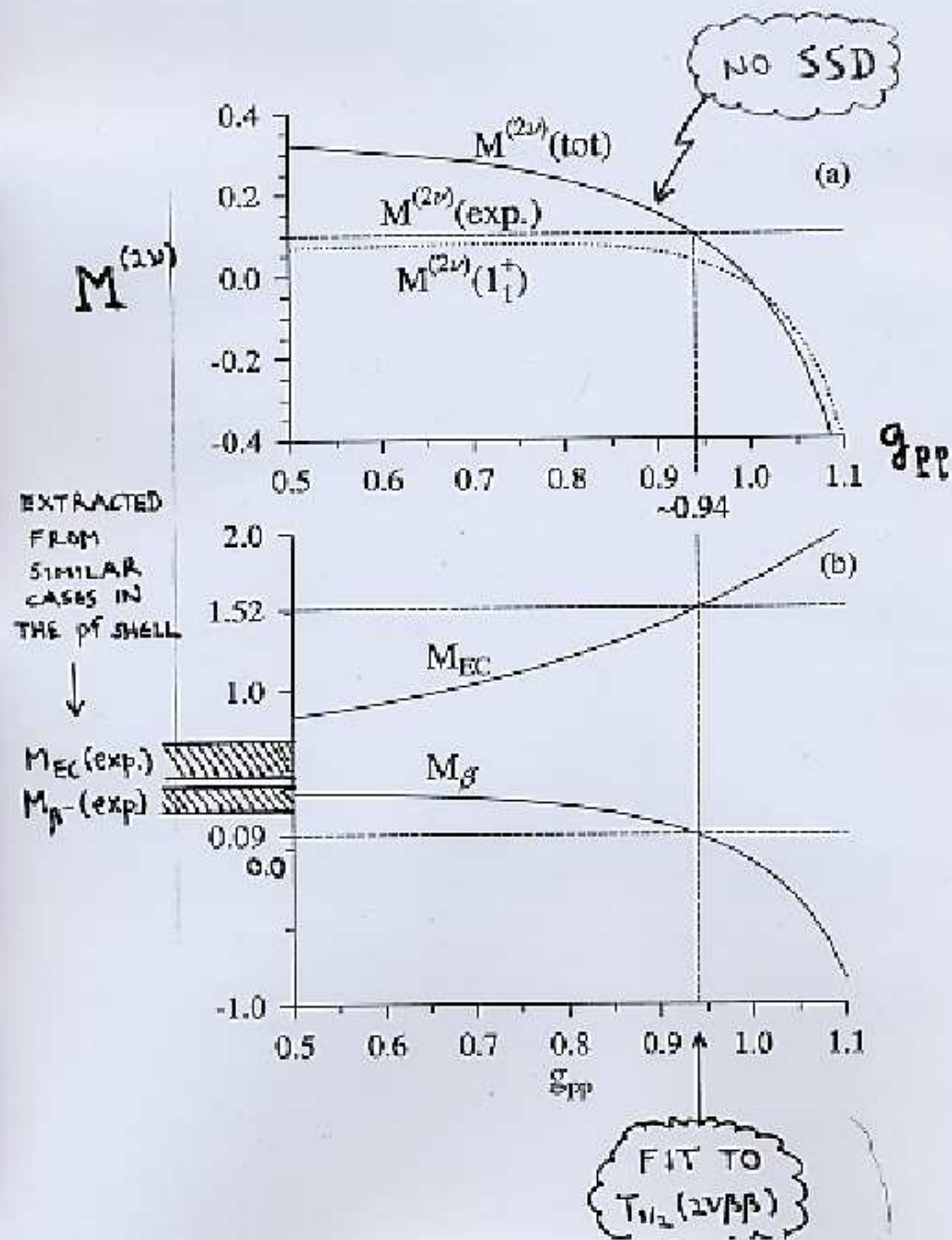
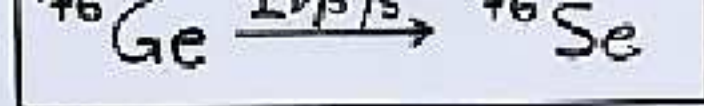


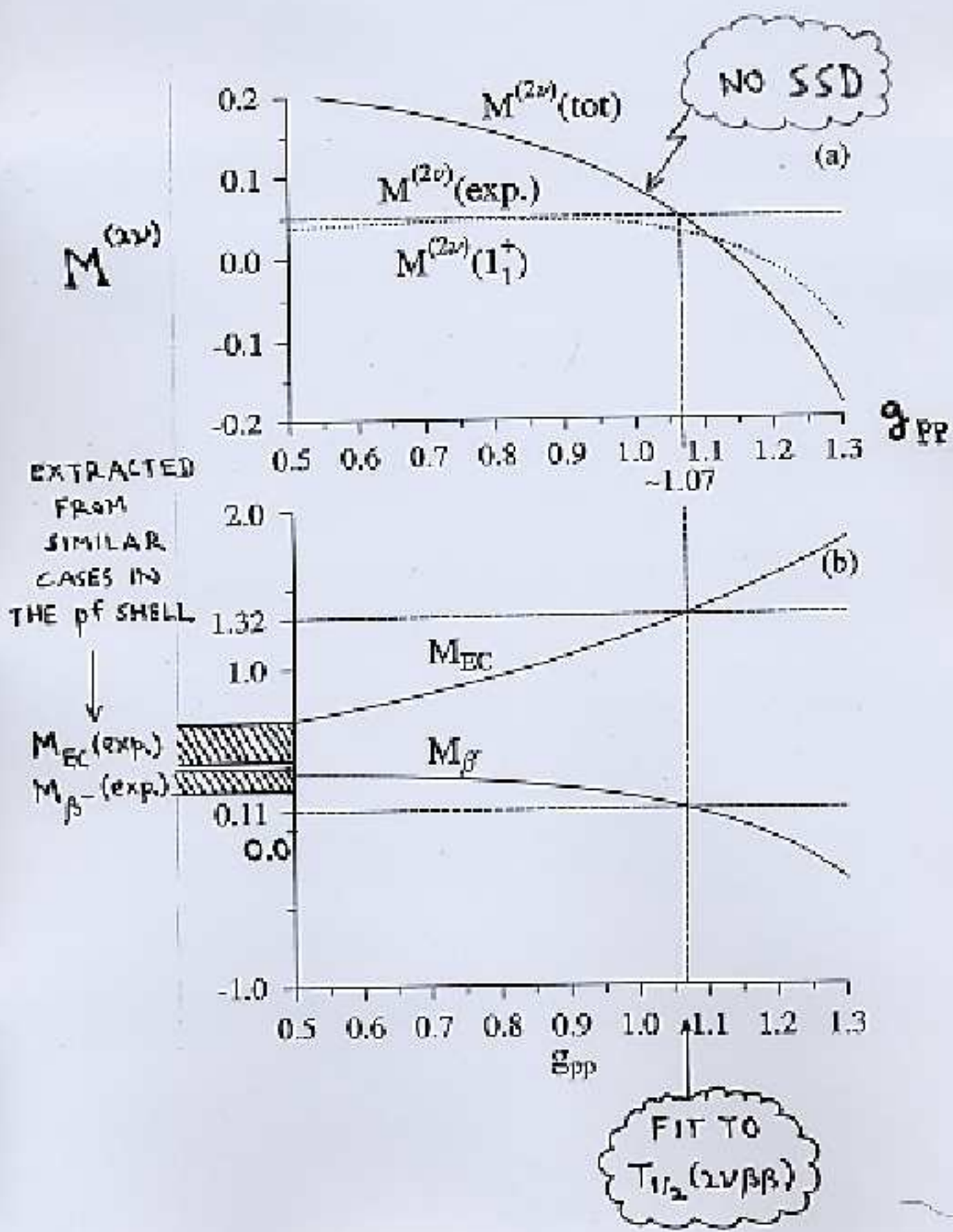
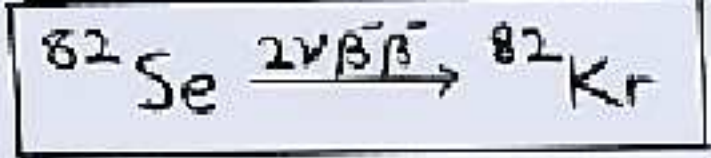
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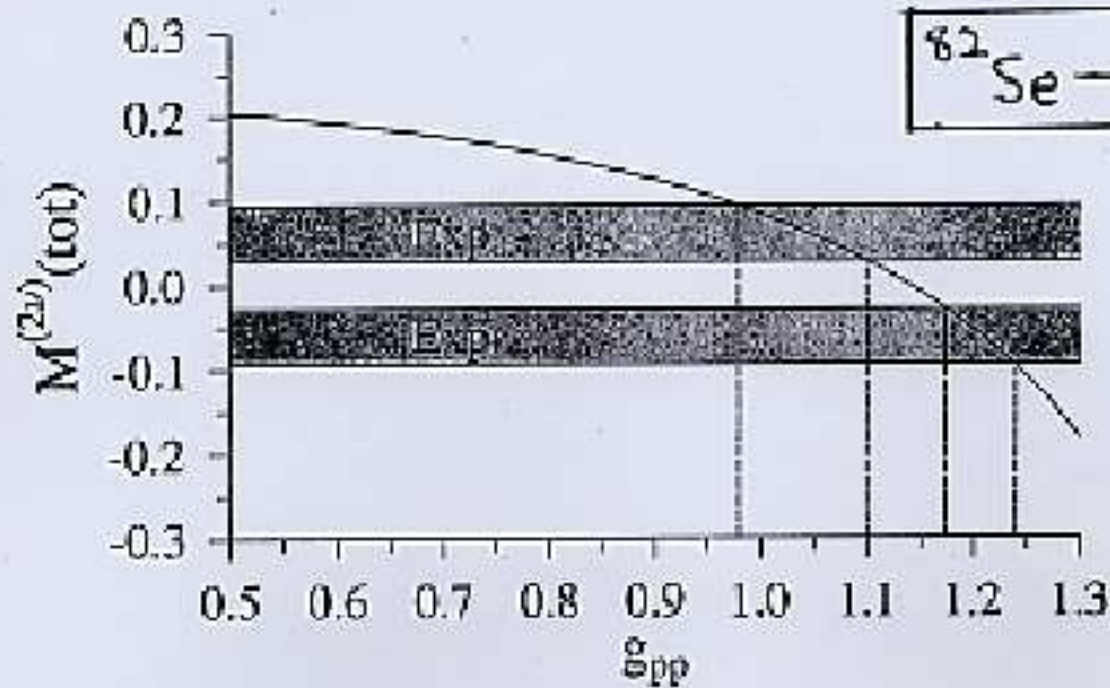
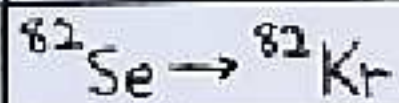
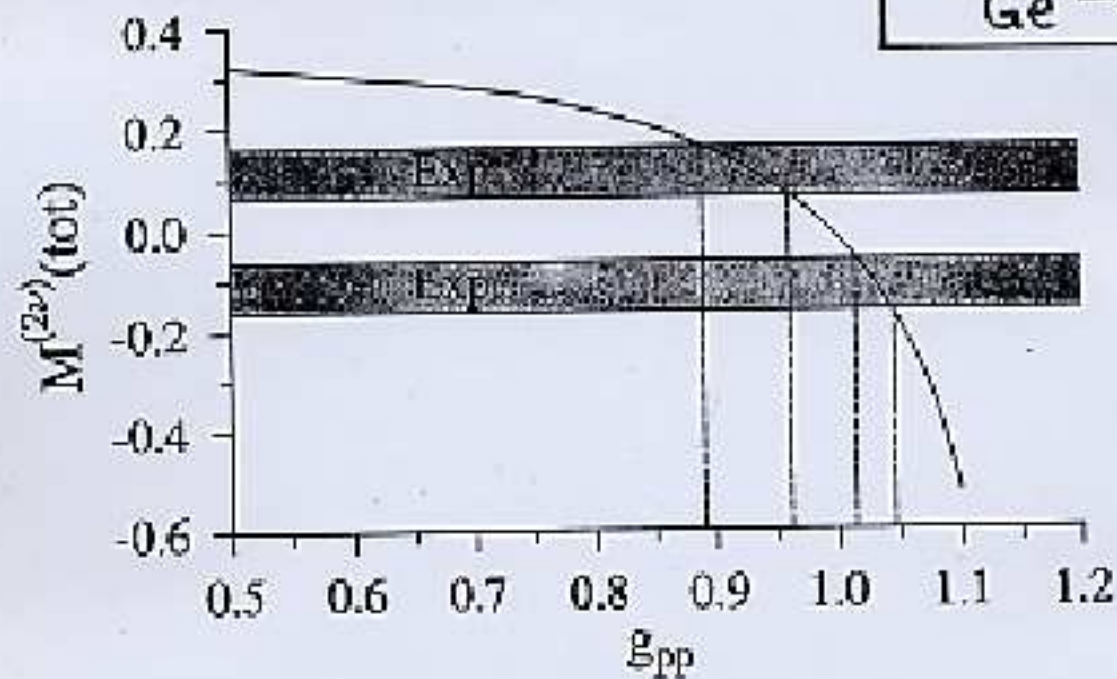
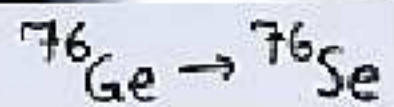


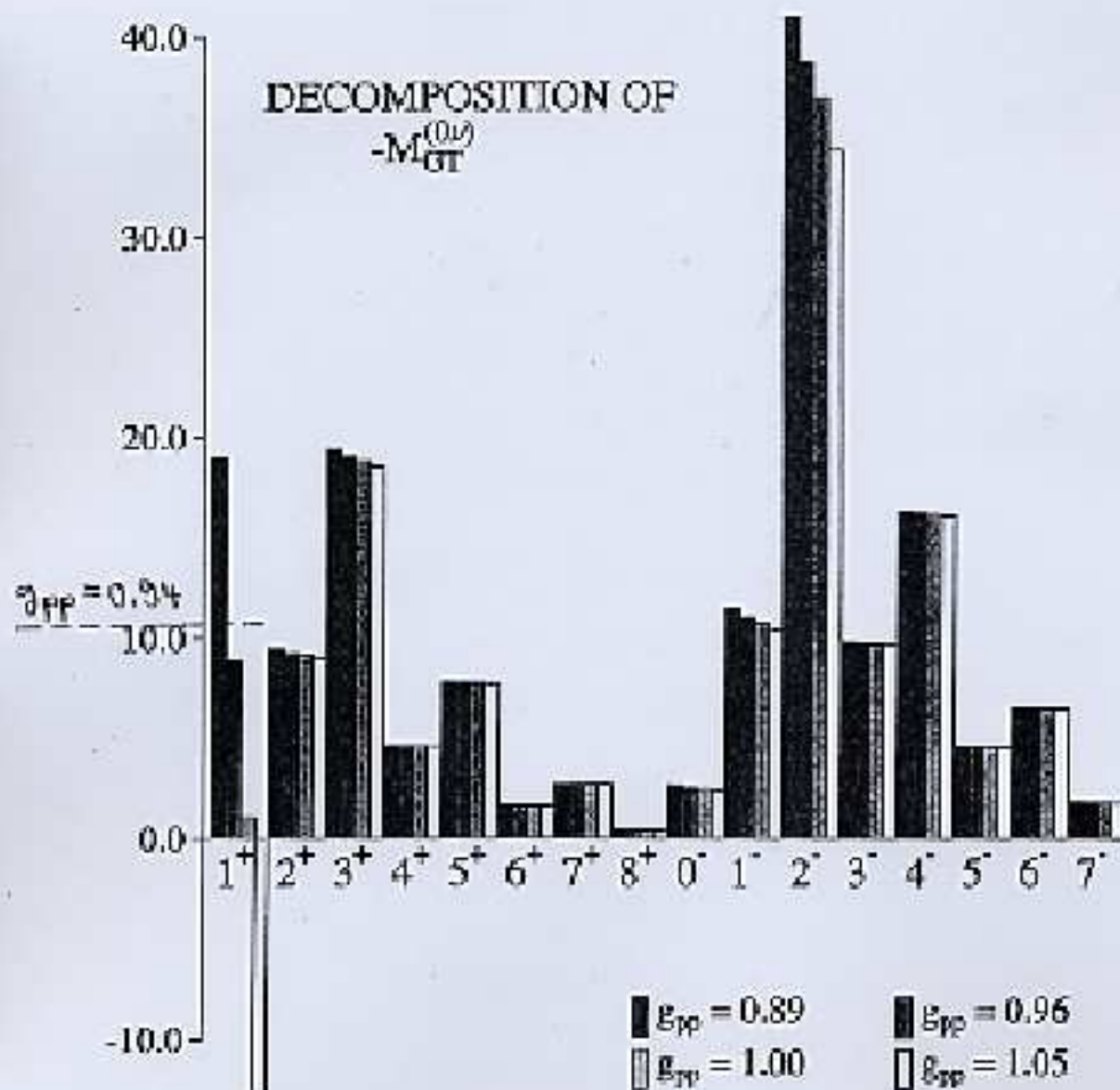




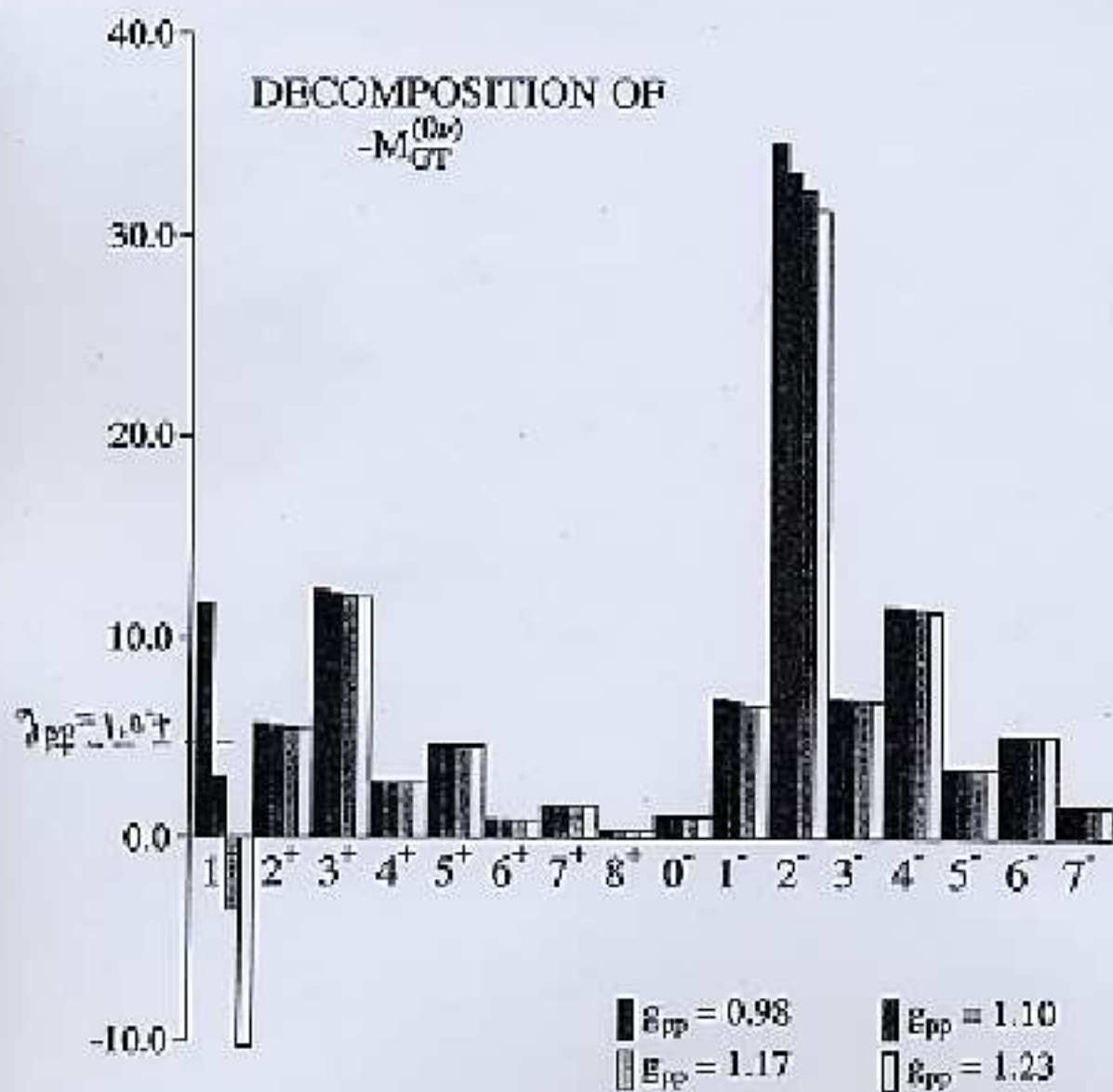
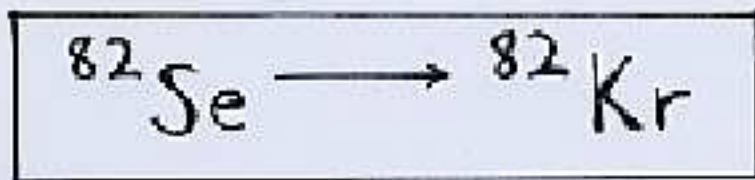
β Decay vs. $2\nu\beta\beta$ Decay in Fitting g_{pp}

Point	Fit to β^- and/or EC decay(s)	Fit to $2\nu\beta^-\beta^-$ decay
1	One, two or more observables can be used for the fit (+)	Only one observable can be used for the fit (-)
2	Direct access to grass-root-level deficiencies of a nuclear model (+)	Two or more compensating errors may conspire to produce a good $2\nu\beta^-\beta^-$ decay rate (-)
3	The beta-decay properties better reproduced (+)	The $2\nu\beta^-\beta^-$ decay properties better reproduced (+)
4	Error limits from comparison of the experimental and computed $2\nu\beta^-\beta^-$ decay rate (+)	Advisable to check against data on β^- decays
5	Largely eliminates the model-space dependence of the computed $0\nu\beta^-\beta^-$ decay rates (+)	Largely eliminates the model-space dependence of the computed $0\nu\beta^-\beta^-$ decay rates (+)
6	Can be extended to study of forbidden contributions, e.g. 2^- , in $0\nu\beta^-\beta^-$ decay (+)	No access to a possible variation of g_{pp} from multipole to multipole (-)
7	Can access $\beta\beta$ decays where no $2\nu\beta\beta$ data exists, see Table 2 (+)	Can access $\beta\beta$ decays where no direct β -decay data exists (^{76}Ge and ^{82}Se) (+)
Balance	7 x (+)	3 x (+) and 3 x (-)

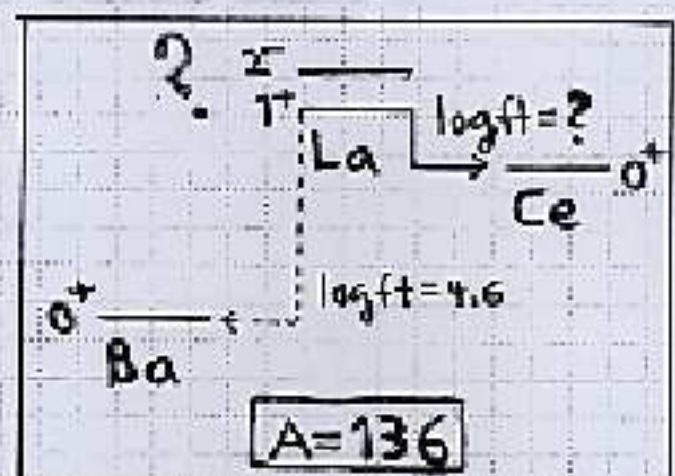
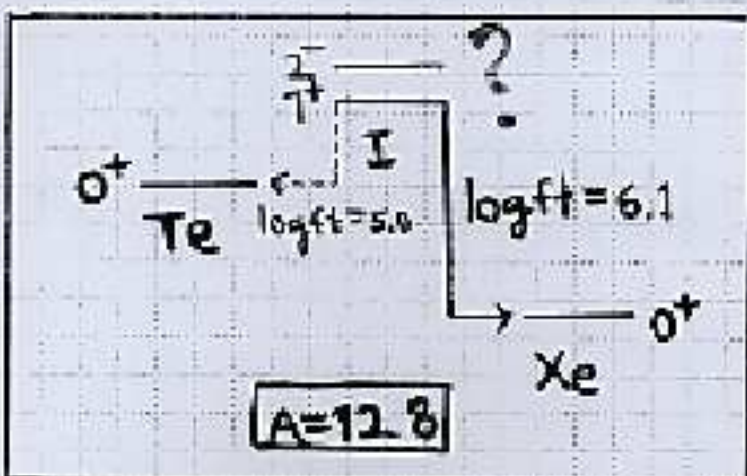
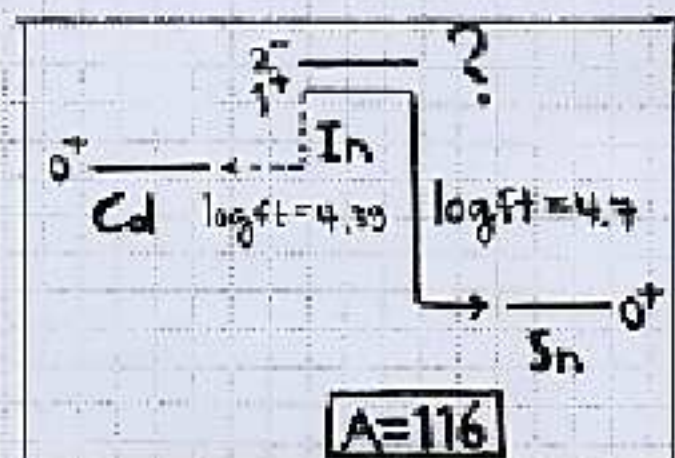
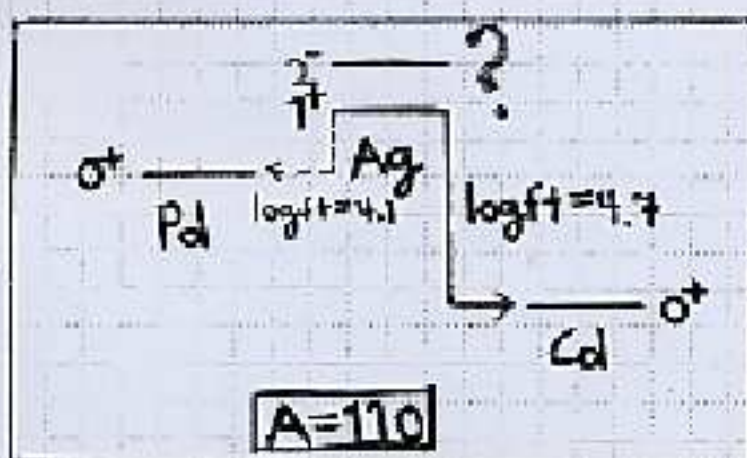
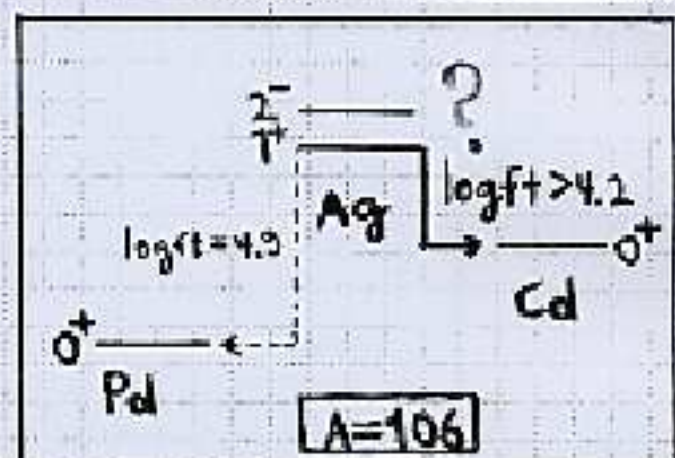
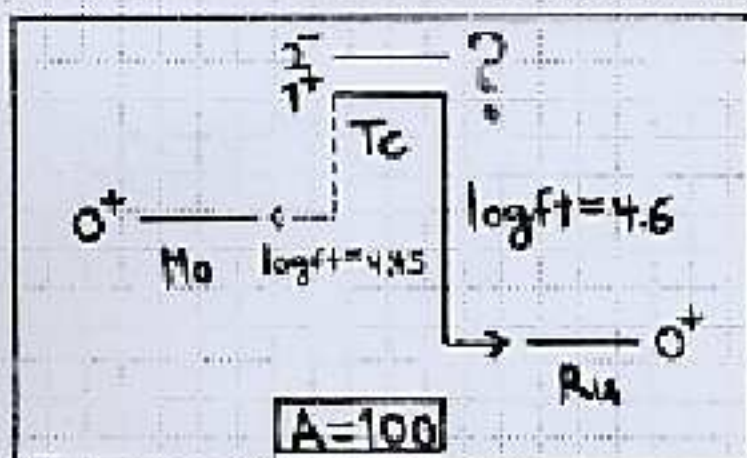
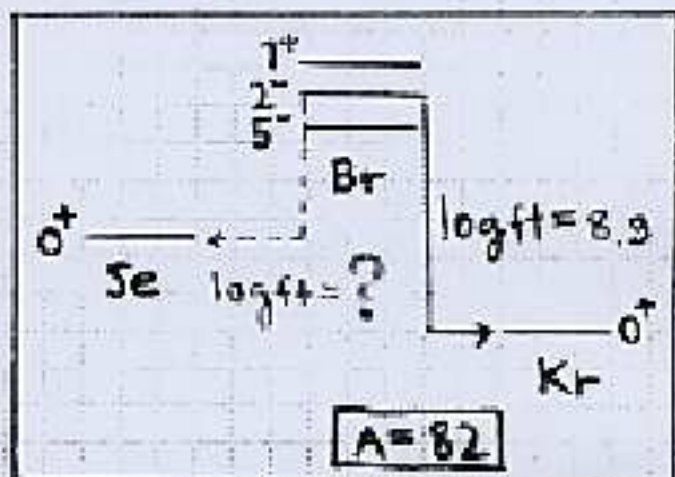
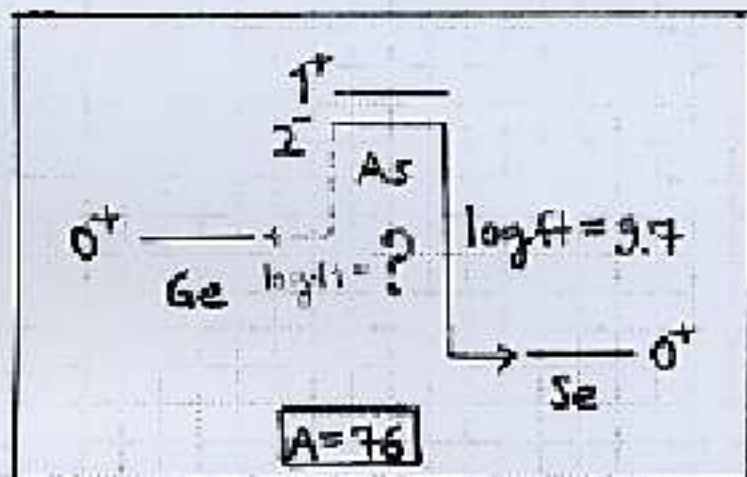


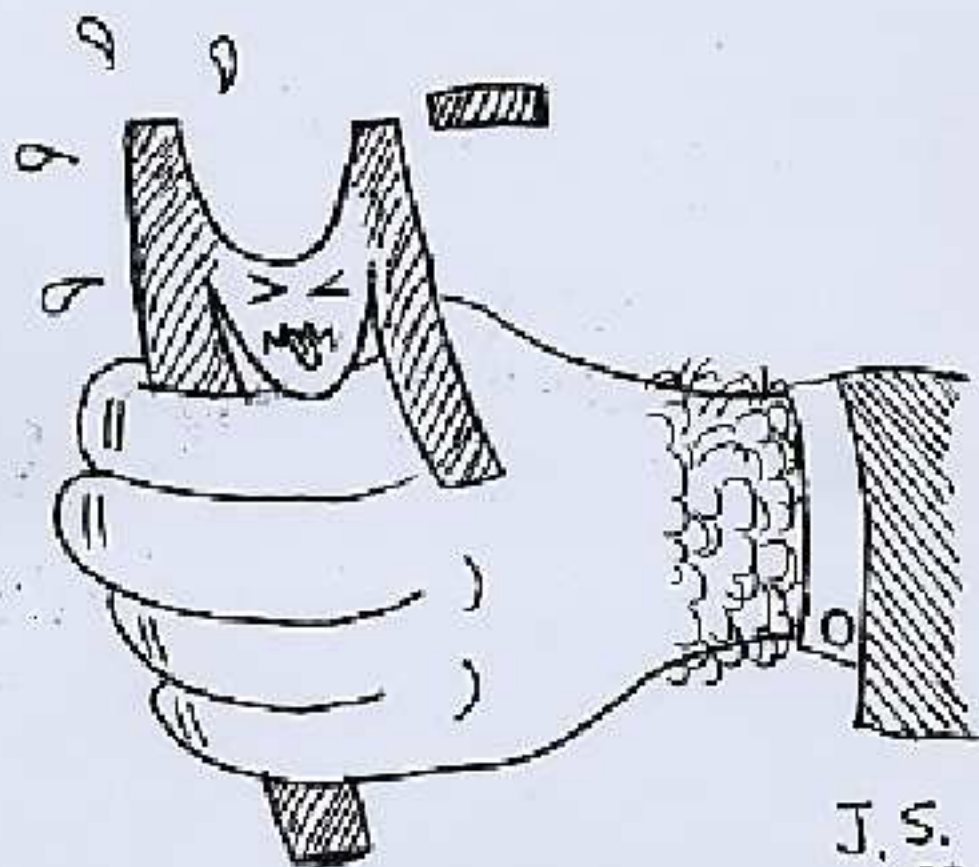


From fit to central value: $g_{\text{PP}} = 0.94$



From fit to central value: $g_{pp} = 1.07$





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MUONIC PROCESSES

The muon disappears from the atomic 1s orbit by

① MUON DECAY ($t_{1/2} \sim 10^{-6} s$):

$$\mu^- \rightarrow e^- + \bar{\nu}_e + \nu_\mu \quad (\text{OMD})$$

$$\mu^- \rightarrow e^- + \bar{\nu}_e + \nu_\mu + \gamma \quad (\text{RMD})$$

② CAPTURE by the nucleus:

$$\mu^- + (A, Z) \rightarrow (A, Z-1) + \nu_\mu \quad (\text{OMC})$$

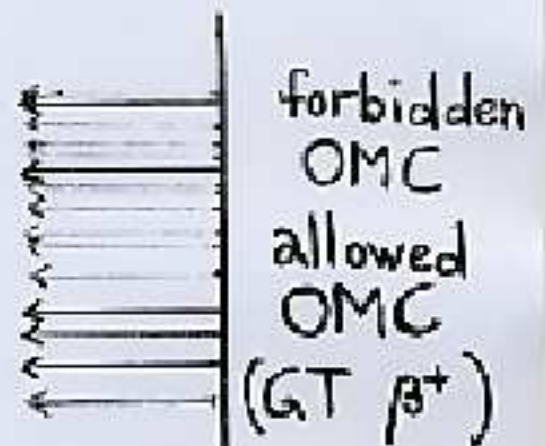
$$\mu^- + (A, Z) \rightarrow (A, Z-1) + \nu_\mu + \gamma \quad (\text{RMC})$$

③ CONVERSION processes:

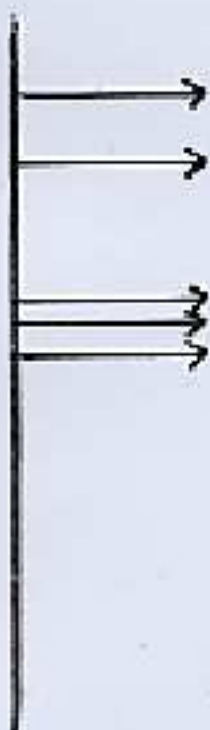
$$\mu^- + (A, Z) \rightarrow e^- + (A, Z)$$

$$\mu^- + (A, Z) \rightarrow e^+ + (A, Z-2)$$

(lepton-flavour violating)

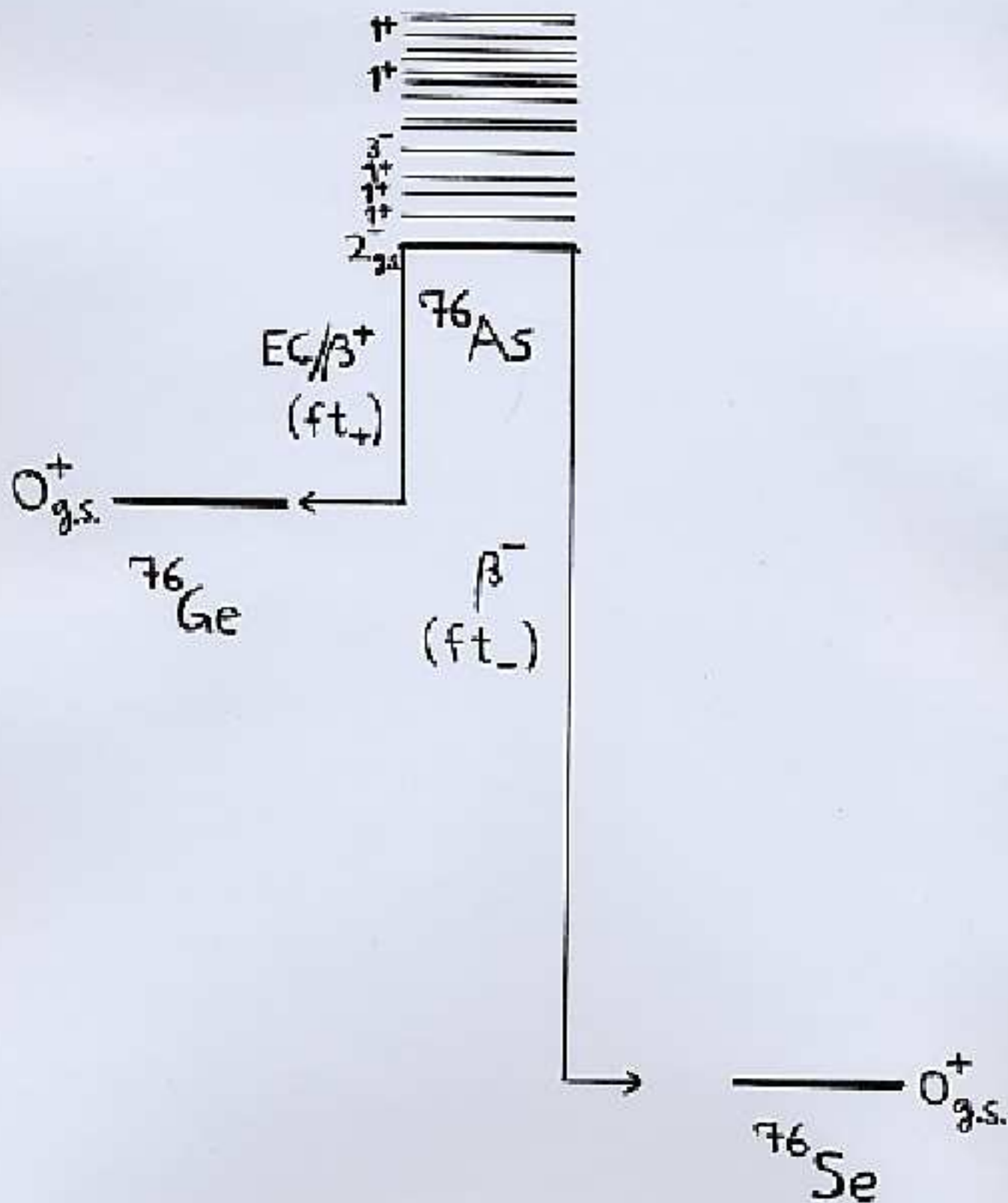


$$\beta/\beta_L^{(2\nu)}$$

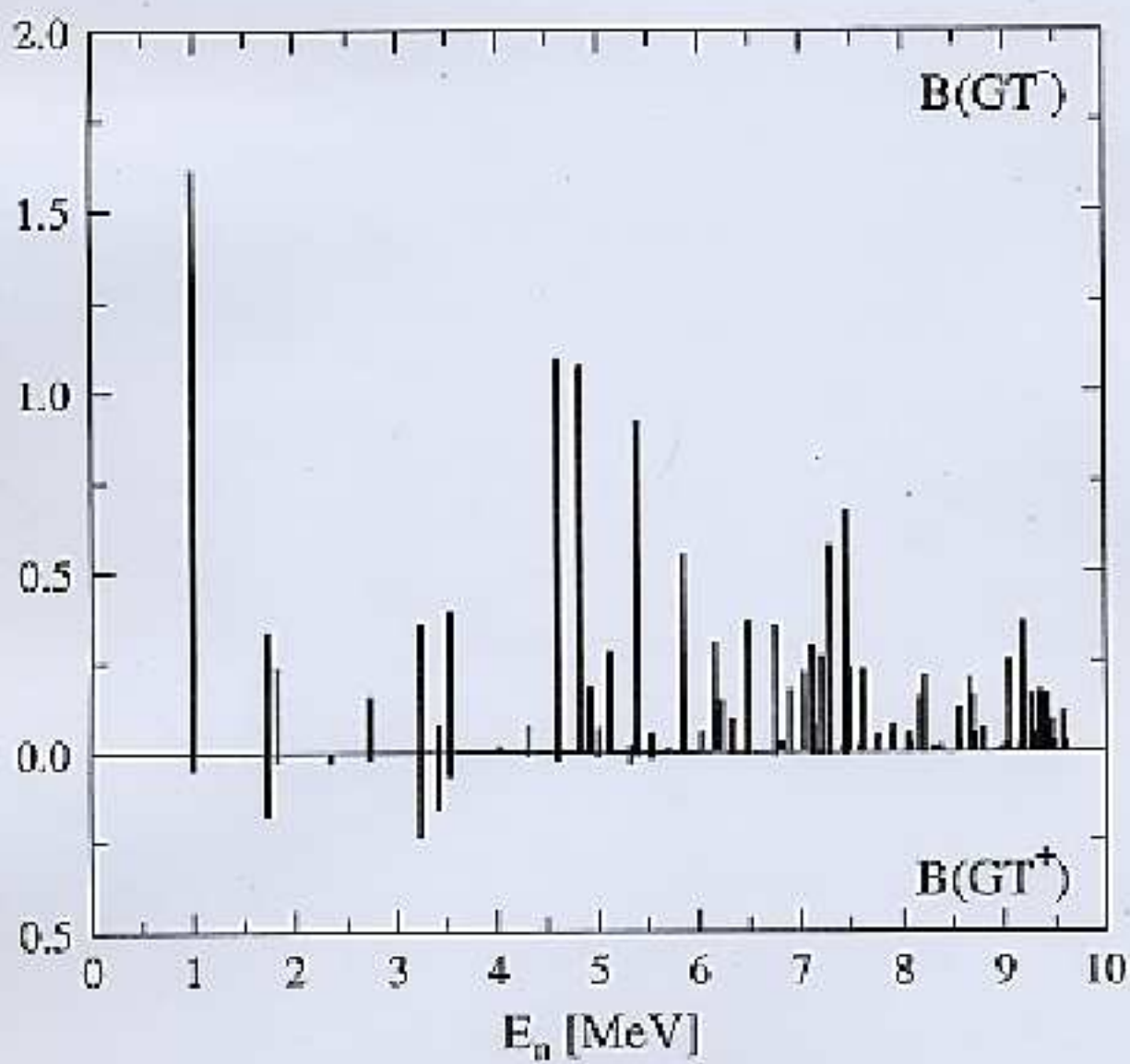


$$\beta/\beta_R^{(2\nu)}$$

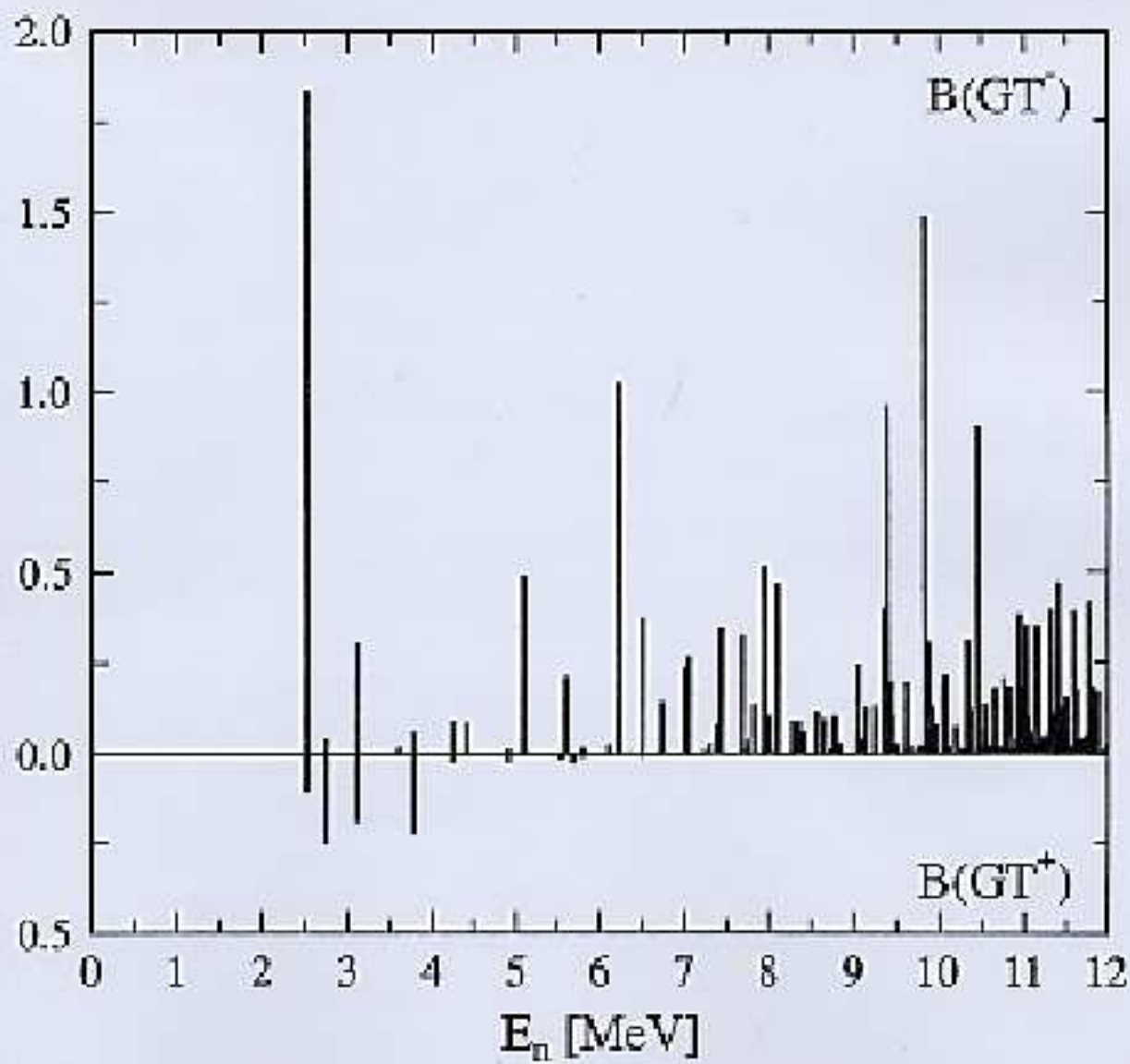




^{46}Ca (fpbp)

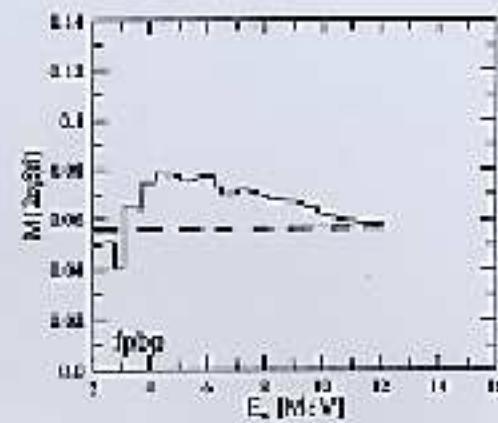
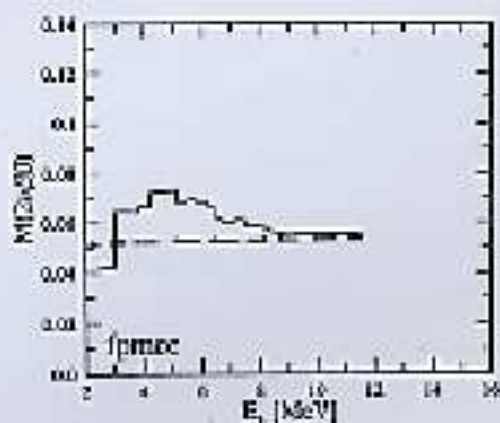
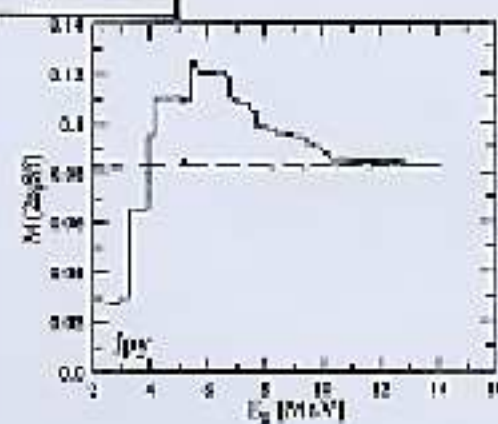
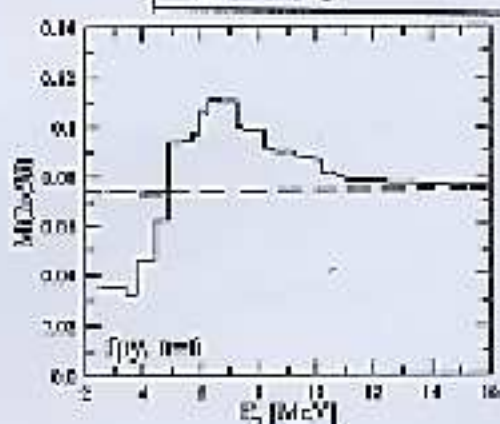
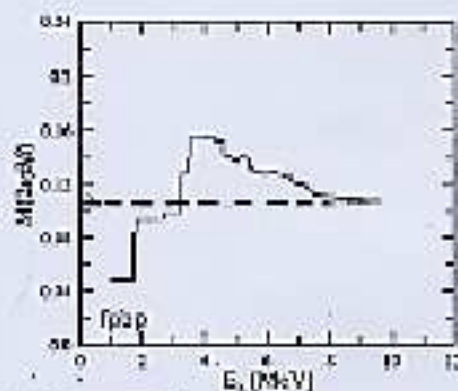
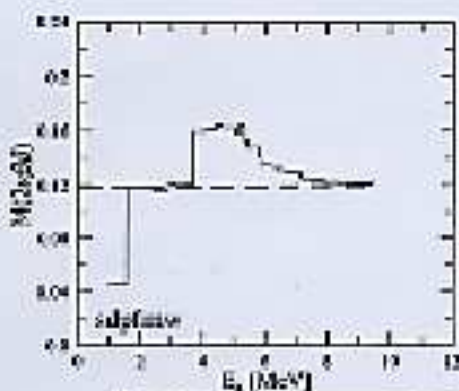
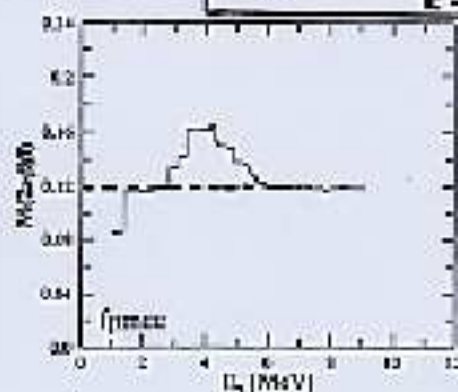
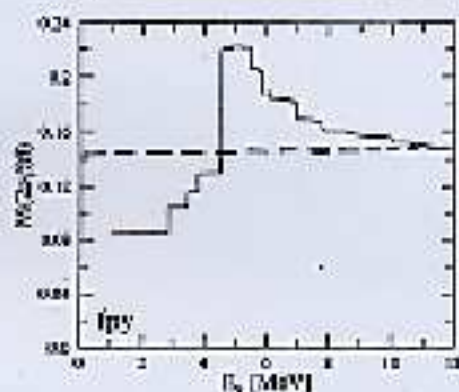


^{48}Ca (fpbp)

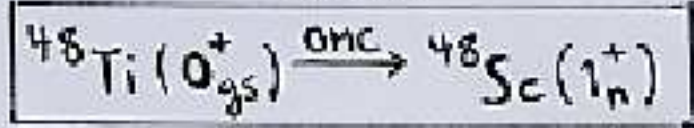
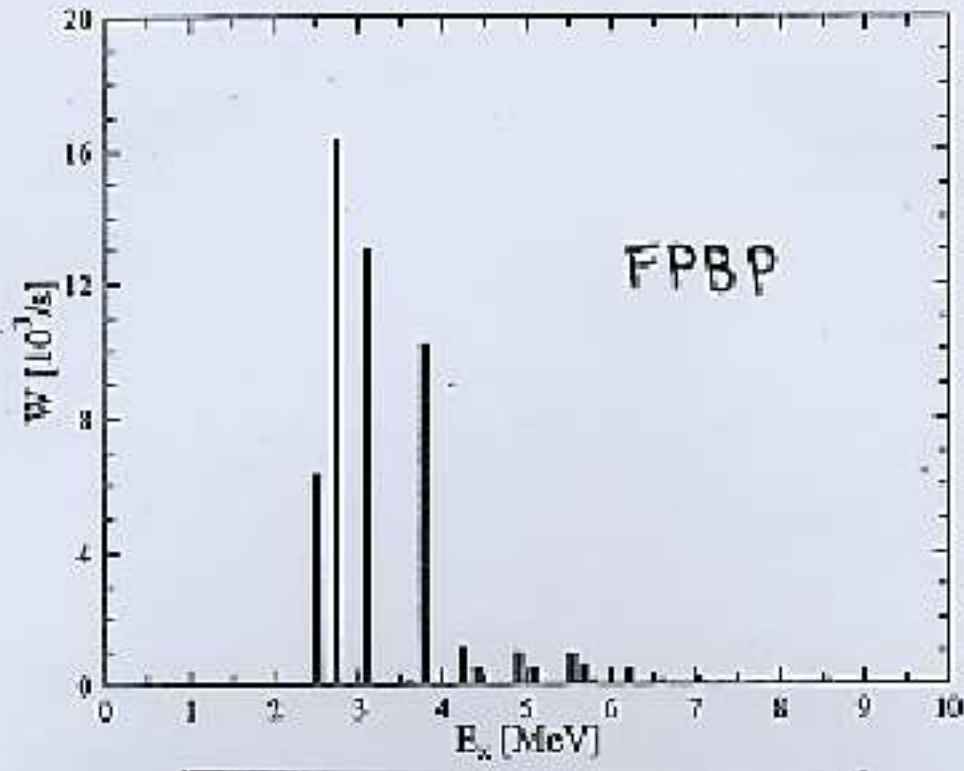
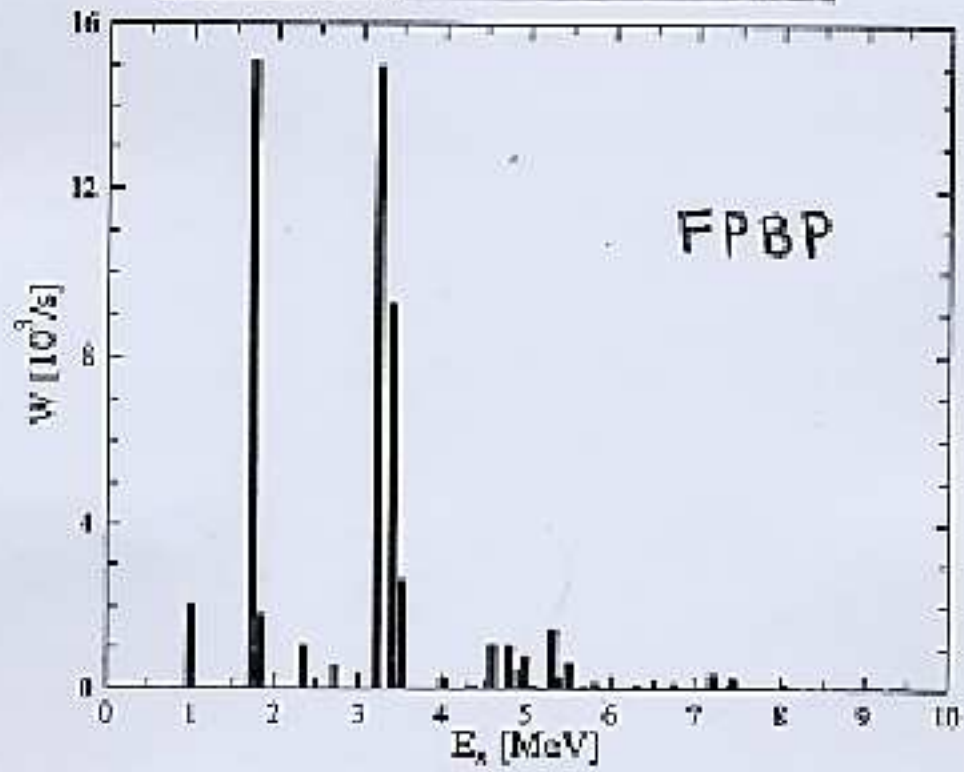
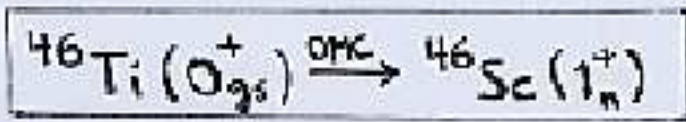




$$M_{GT}^{(0)}(E_x) = \sum \frac{M_L M_R}{E E_x \text{DENOM}(E_x)}$$

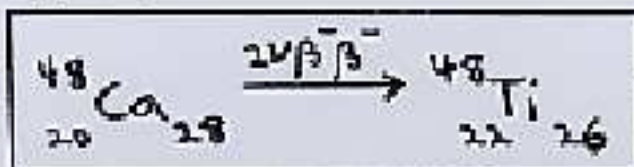
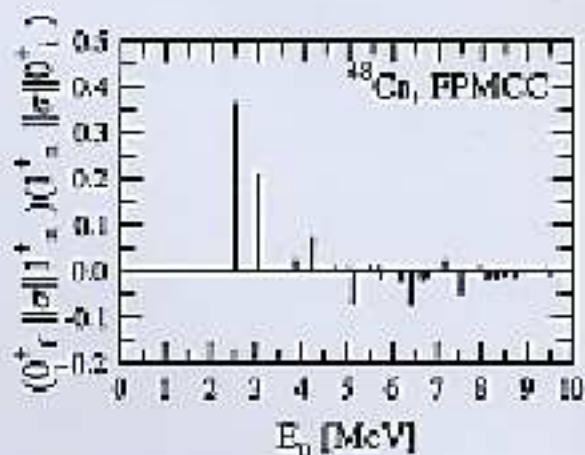
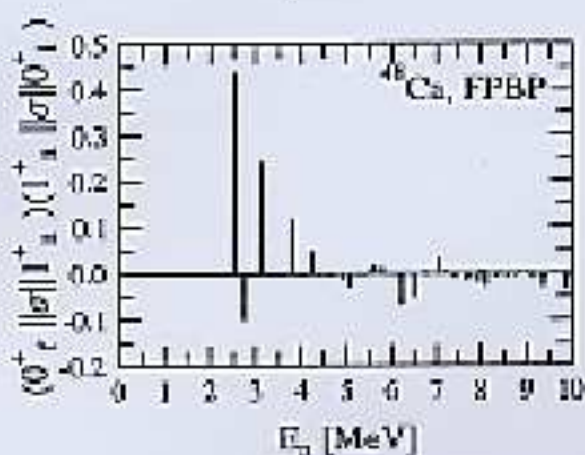
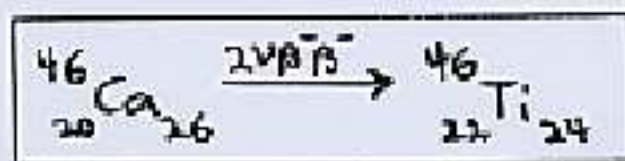
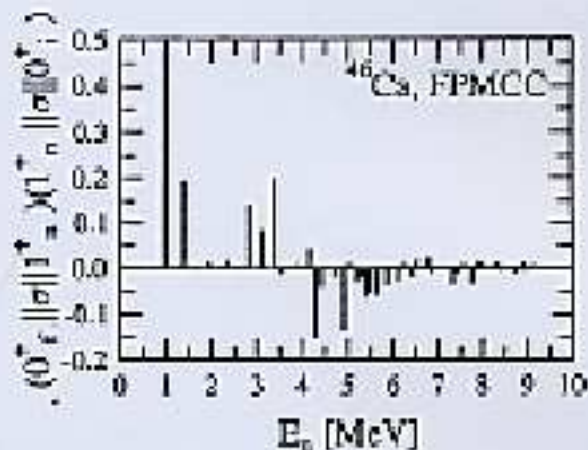
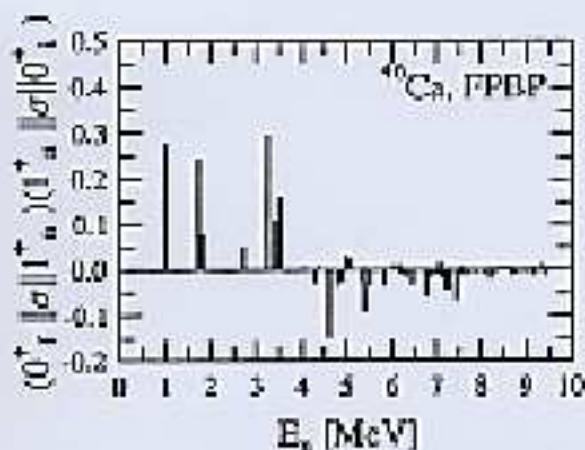


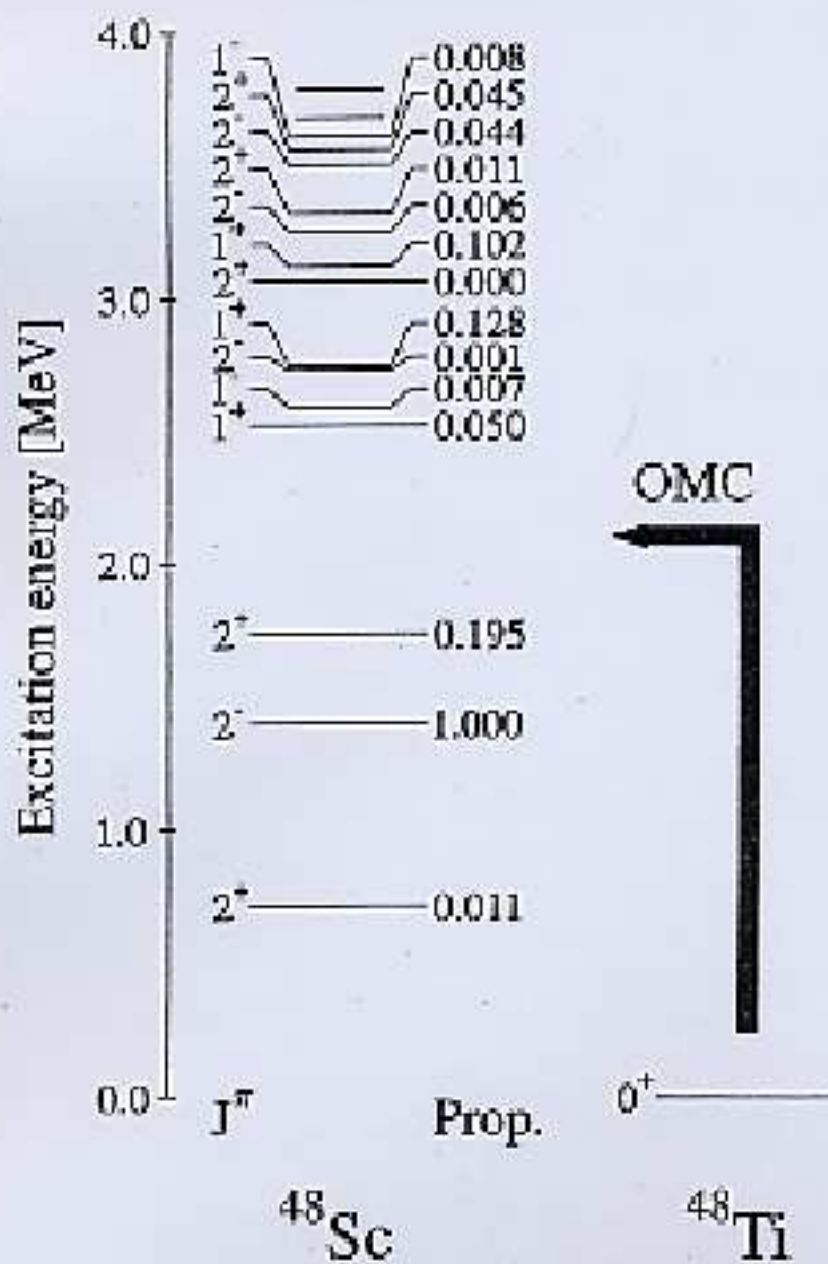
$$^{48}\text{Ca}: M_{GT}^{(0)}(\text{exp}) = (0.034 - 0.115) \text{MeV}^{-1}$$



$$M_{\text{DGT}}^{(2\nu)} = \sum_n \frac{(0_f^+ \| \delta \| 1_n^+) (1_n^+ \| \delta \| 0_i^+)}{n^+ (\frac{1}{2} Q_{pp} + E_n - M_i) / m_e + 1}$$

pf model space





CONCLUSIONS

- μ^- CAPTURE RATE TO THE 1_1^+ STATE
VERY SENSITIVE TO g_{pp} (pn p1e-p1e
int. action) $\leftrightarrow$ LINK IN THE $2\nu\beta\beta$
DECAY!
- GOOD POSSIBILITY TO CONSTRAIN THE WAVE
FUNCTIONS OF THE $\beta\beta$ INTERMEDIATE NUCLEUS
- POSSIBILITY TO CONSTRAIN THE MAGNITUDE OF
THE INDUCED PSEUDO-SCALAR COMPONENT OF
THE HADRONIC CHARGED WEAK-INTERACTION
CURRENT
- PRESENT STATUS :
 - ANALYSIS OF THE PSI EXPERIMENT(S)
GOING ON
 - SHELL-MODEL CALCULATIONS FOR THE
SYSTEMS

