

W+3 jet production

— *signal or background* —

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W + 3 jets

- I. W + 3 jets measured at the Tevatron, but LO varies by more than a factor 2 for reasonable changes in scales

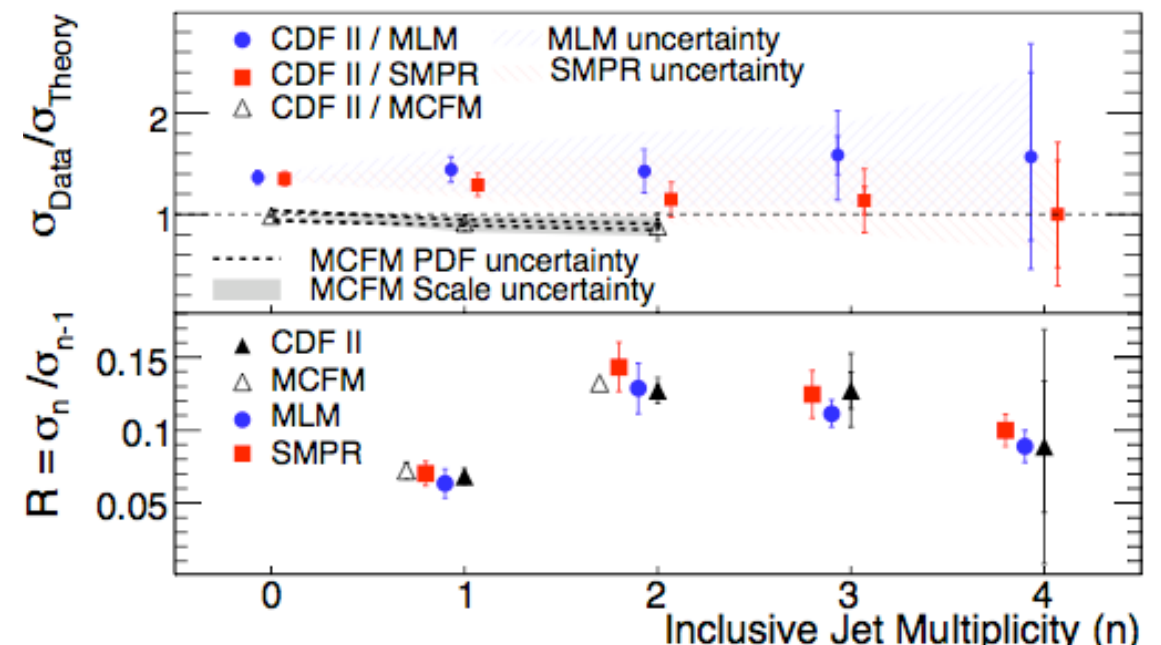
| | W^\pm , TeV | W^+ , LHC | W^- , LHC |
|--------------------------------|----------------|-----------------|-----------------|
| σ [pb], $\mu = 40$ GeV | 74.0 ± 0.2 | 783.1 ± 2.7 | 481.6 ± 1.4 |
| σ [pb], $\mu = 80$ GeV | 45.5 ± 0.1 | 515.1 ± 1.1 | 316.7 ± 0.7 |
| σ [pb], $\mu = 160$ GeV | 29.5 ± 0.1 | 353.5 ± 0.8 | 217.5 ± 0.5 |

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II. CDF data for W + n jets with n=1,2 is described **exceptionally well by NLO QCD**
 \Rightarrow verify this for 3 and more jets



$W + 3 \text{ jets}$

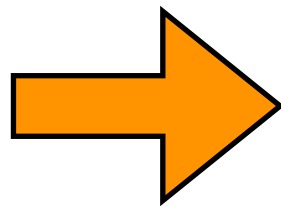
III. $W/Z + 3 \text{ jets}$ of interest at the LHC, as one of the backgrounds to
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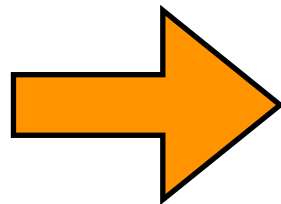
IV. Calculation highly non-trivial optimal testing ground

$$0 \rightarrow \bar{u} d g g g W^+$$



1203 + 104 Feynman diagrams

$$0 \rightarrow \bar{u} d \bar{Q} Q g W^+$$



258 + 18 Feynman diagrams

Generalized unitarity

I will not explain the method.

I will concentrate on applications & recent results

References:

- Ellis, Giele, Kunszt '07 [Unitarity in $D=4$]
- Giele, Kunszt, Melnikov '08 [Unitarity in $D\neq 4$]
- Giele & GZ '08 [All one-loop N -gluon amplitudes]
- Ellis, Giele, Melnikov, Kunszt '08 [Massive fermions, $ttgg$ amplitudes]
- Ellis, Giele, Melnikov, Kunszt, GZ '08 [$W+5p$ one-loop amplitudes]
- Ellis, Melnikov, GZ '09, Melnikov & GZ '09 [$W+3$ jets]

These papers heavily rely on previous work

- Bern, Dixon, Kosower '94 [Unitarity, oneloop from trees]
- Ossola, Pittau, Papadopoulos '06 [OPP]
- Britto, Cachazo, Feng '04 [Generalized cuts]
- [...]

The F90 Rocket program

Rocket science!

Eruca sativa = Rocket = roquette = arugula = rucola

Recursive unitarity calculation of one-loop amplitudes



So far computed one-loop amplitudes:

- ✓ N-gluons
- ✓ qq + N-gluons
- ✓ qq + W + N-gluons
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NB: N is a parameter in Rocket
In perspective, for gluons:

$N = 6 \Rightarrow 10860$ diags.

$N = 7 \Rightarrow 168925$ diags.

Successfully computed up to $N=20$

Cross-section calculation

- Consider the NLO **leading color approximation**, keep n_f dependence exact (important for beta function) but neglect $1/N_c^2$ terms
- Real radiation part:
 -  leading color tree level **V+6** parton amplitudes computed recursively
 -  we use **Catani-Seymour subtraction** terms modified to deal with the minimal set of color structures needed at leading color
- Real + virtual implemented in the **MCFM** parton level integrator

Full-color NLO calculation done by Berger et al. '09

Leading color adjustment

Define

$$\mathcal{R}_{\mathcal{O}} = \frac{\int \mathcal{O}(p) d\sigma_{LO}^{\text{FC}}(\mu, p)}{\int \mathcal{O}(p) d\sigma_{LO}^{\text{LC}}(\mu, p)}$$

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Other $\mathcal{O}(1\%)$ effects neglected:

- CKM set to unity $\Rightarrow \sim -1\%$
- W treated onshell $\Rightarrow \sim +1\%$

CDF cuts

$$p_{\perp,j} > 20\text{GeV} \quad p_{\perp,e} > 20\text{GeV} \quad E_{\perp,\text{miss}} > 30\text{GeV}$$

$$|\eta_e| < 1.1$$

$$M_{\perp,W} > 20\text{GeV}$$

$$\mu_0 = \sqrt{p_{\perp,W}^2 + M_W^2}$$

$$\mu = \mu_R = \mu_F = [\mu_0/2, 2\mu_0]$$

- PDFs: cteq6l1 and cteq6m
- CDF applies lepton-isolation cuts. This is a $O(10\%)$ effect. Lepton-isolation has been corrected for (would not have been needed ...)
No lepton isolation applied
- CDF uses JETCLU with $R = 0.4$, but this is **not infrared safe**, use a different jet-algorithm

Jet-algorithms

- CDF uses JETCLU which is not infrared safe
- NLO calculation with JETCLU not possible
- use e.g. SIScone and anti-kt algorithm which are IR safe
- can compare Leading order results for these algorithm (even if meaning of LO for JETCLU is questionable ...)

Leading order:

| Algorithm | R | $E_{\perp}^{\text{jet}} > 20 \text{ GeV}$ | $E_{\perp}^{\text{3rdjet}} > 25 \text{ GeV}$ |
|-------------------|-----|---|--|
| JETCLU | 0.4 | $1.845(2)^{+1.101(3)}_{-0.634(2)}$ | $1.008(1)^{+0.614(2)}_{-0.352(1)}$ |
| SIScone | 0.4 | $1.470(1)^{+0.765(1)}_{-0.560(1)}$ | $0.805(1)^{+0.493(1)}_{-0.281(1)}$ |
| anti- k_{\perp} | 0.4 | $1.850(1)^{+1.105(1)}_{-0.638(1)}$ | $1.010(1)^{+0.619(1)}_{-0.351(1)}$ |

SIScone: Salam & Soyez '07;
anti-kt: Cacciari, Salam, Soyez '08

At LO anti-kt $R=0.4$ is closer to JETCLU

Moral:

precision comparison with theory require that experiments use IR-safe algorithms

Cross-section at the Tevatron

$$\sigma_{W+3j}(p_{\perp,j} > 25 \text{ GeV}) = (0.84 \pm 0.24) \text{ pb}$$

CDF

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LO^{LC}

SIS.

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a-k_t

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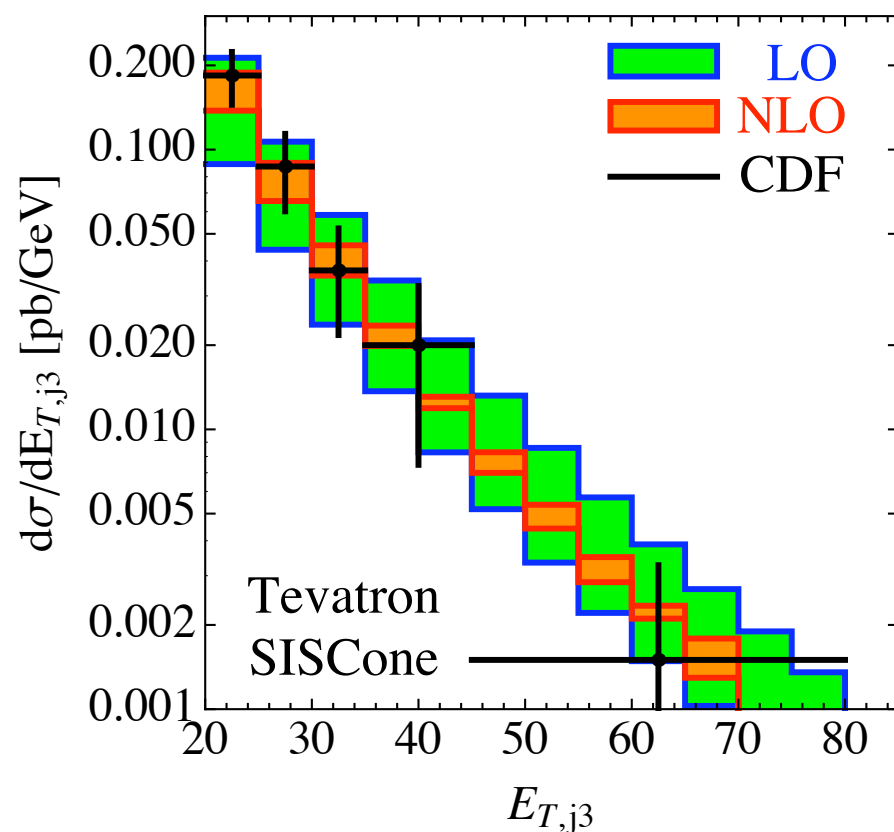
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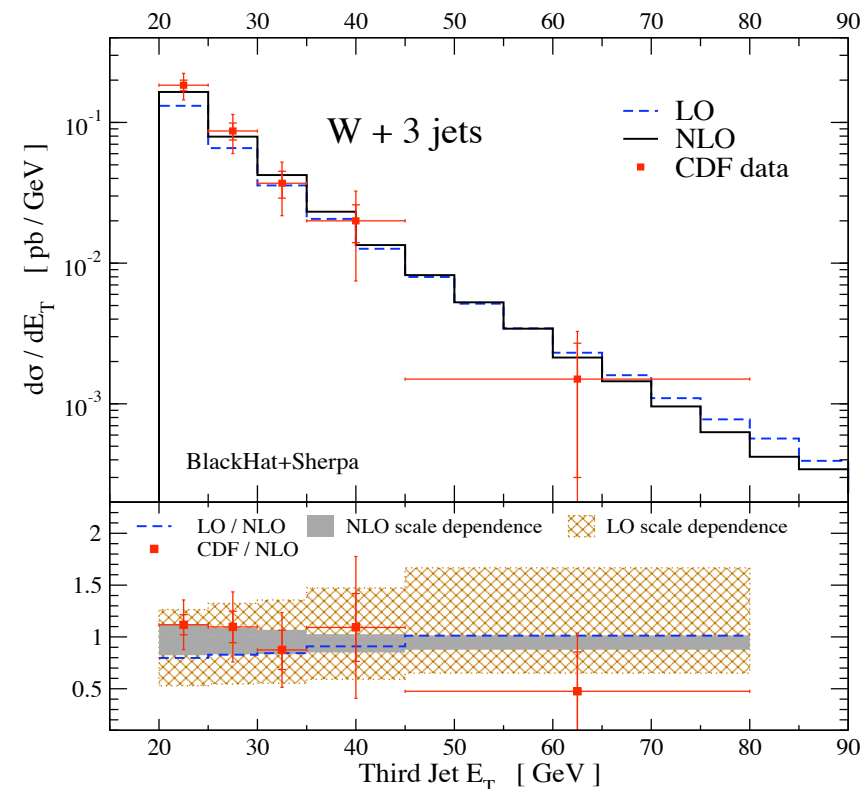
- ⇒ agreement between independent calculations to within 3%
- ⇒ leading color approximation works very well. After leading color adjustment procedure it is good to 3%
- ⇒ important (10% or more) differences due to different jet-algorithms.
High precision comparison impossible if using different algorithms

Tevatron: sample distribution: $E_{T,j3}$

***NB:** CDF \Rightarrow JetCLU VERSUS NLO Theory \Rightarrow SISCone*



Ellis et al '09 (LC)



Berger et al '09

- ☺ agreement with CDF data (within currently large errors)
- ☺ small $K=1.0-1.1$, reduced uncertainty: 50% (LO) \rightarrow 10% (NLO)
- ☺ first applications of new techniques to $2 \rightarrow 4$ LHC processes

Dual role of SM processes

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- primary signals (apply signal cuts)
- unwanted background (apply background cuts)

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Standard procedure

- study a given process with **signal cuts** \Rightarrow refine theoretical tools
- once good understanding of the process is achieved with signal cuts (e.g. low p_t region) **extrapolate to background cuts** region (e.g. high p_t)

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How reliable is this procedure ?

Purpose of background cuts: push into corners of phase-space the SM process, therefore the robustness of the procedure is not assured.
NLO QCD predictions for non-trivial processes can shed light on this.

$W^+ + 3 \text{ jets at the LHC}$

In the following: use highly non-trivial NLO calculation of $W^+ + 3 \text{ jets}$ to illustrate/study this issue

Signal-cut setup (inspired by CMS studies):

$$E_{\text{CM}} = 10 \text{ TeV}$$

$$E_{\perp, \text{jet}} = 30 \text{ GeV}$$

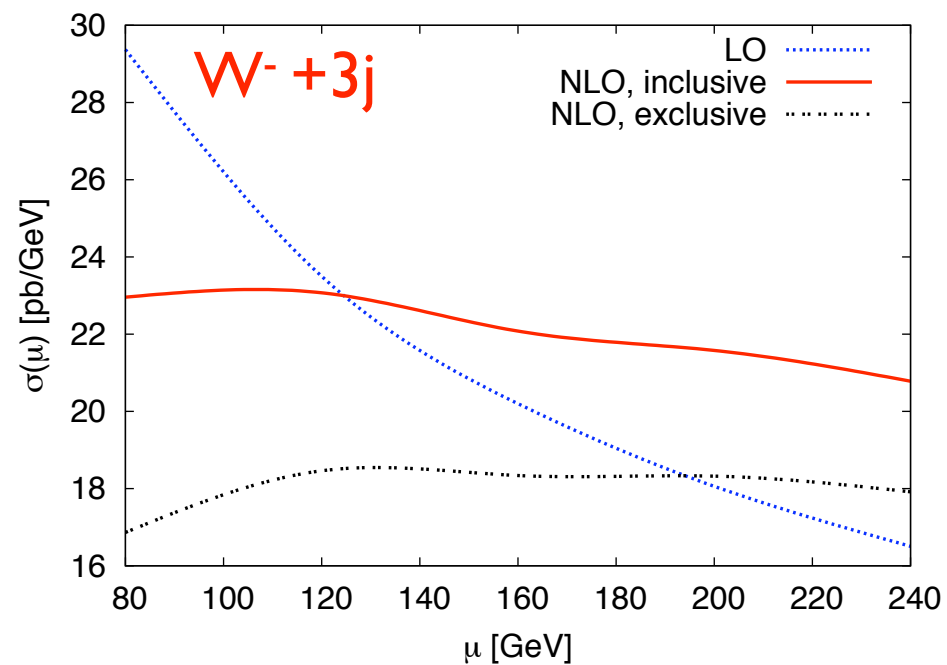
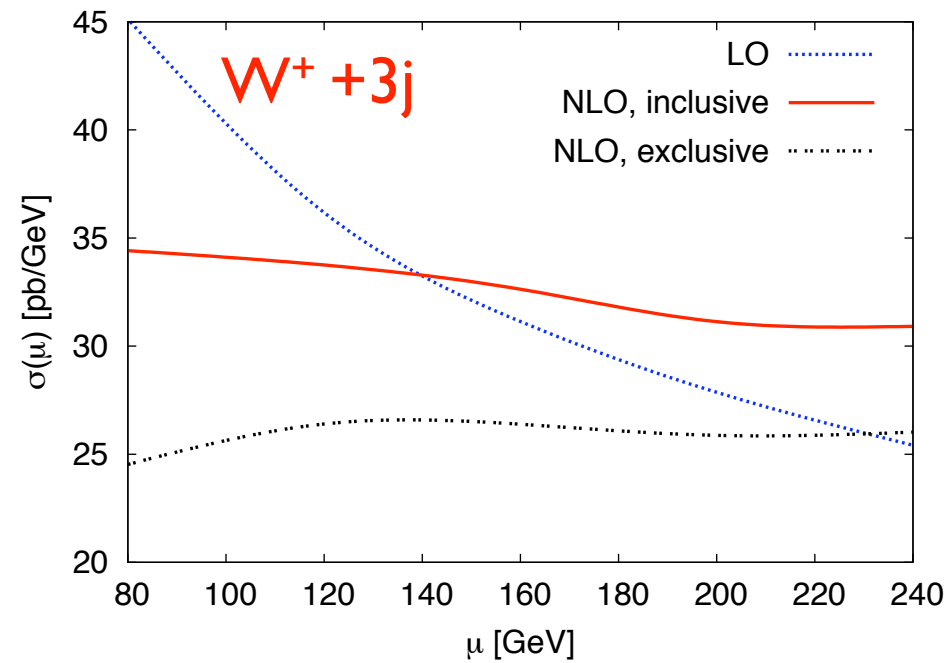
$$E_{\perp, e} = 20 \text{ GeV}$$

$$E_{\perp, \text{miss}} = 15 \text{ GeV} \quad M_{\perp, W} = 30 \text{ GeV} \quad |\eta_e| < 2.4 \quad |\eta_{\text{jet}}| < 3$$

$$\mu_0 = \sqrt{p_{\perp, W}^2 + M_W^2} \quad \mu = \mu_R = \mu_F = [\mu_0/2, 2\mu_0]$$

Jets: SIscone with $R = 0.5$; PDFs: cteq6l1/cteq6m

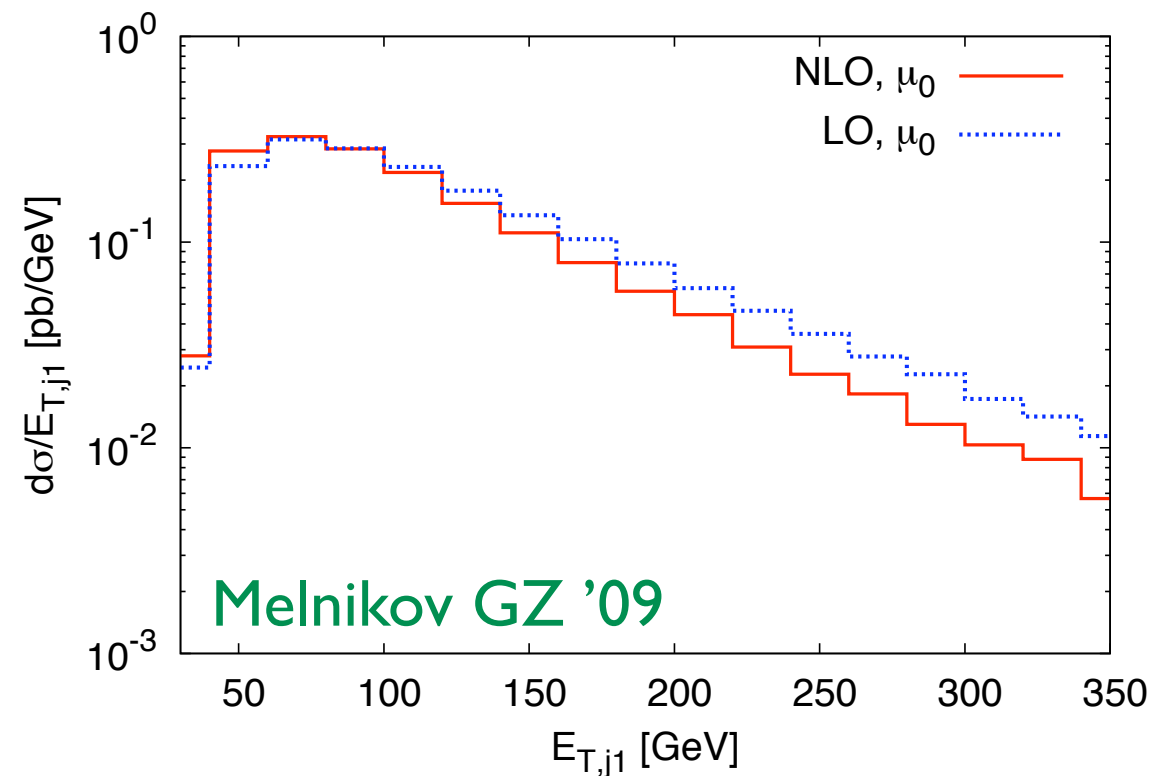
Scale dependence



Melnikov & GZ '09

- scale dependence considerably reduced at NLO (both inclusive and exclusive)
- NLO tends to reduce cross-section
- because of very large scale dependence of LO, quoting a K-factor not very meaningful

Sample transverse energy distribution



Renormalization and factorization scale set to

$$\mu_0 = \sqrt{p_{T,W}^2 + m_W^2}$$

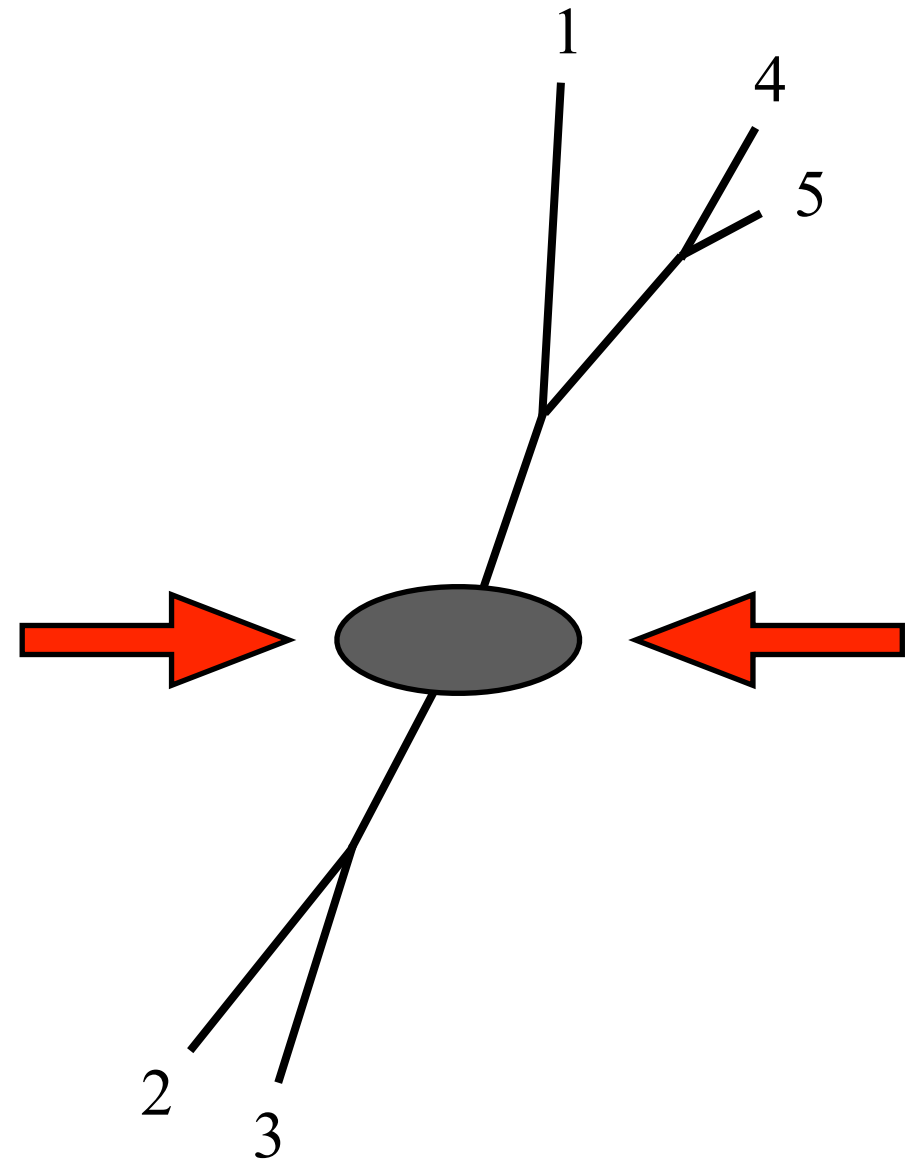
- with scale μ_0 : considerable change in shape between LO and NLO (extrapolation of LO from low p_t to high p_t would fail badly)
- but origin of the change in shape well understood: at high E_T , μ_0 is smaller than typical scales of the QCD branching \Rightarrow LO overshoots the result

Can one do a more sophisticated LO calculation?

Local (CKKW) scale

Local scale choice (CKKW):

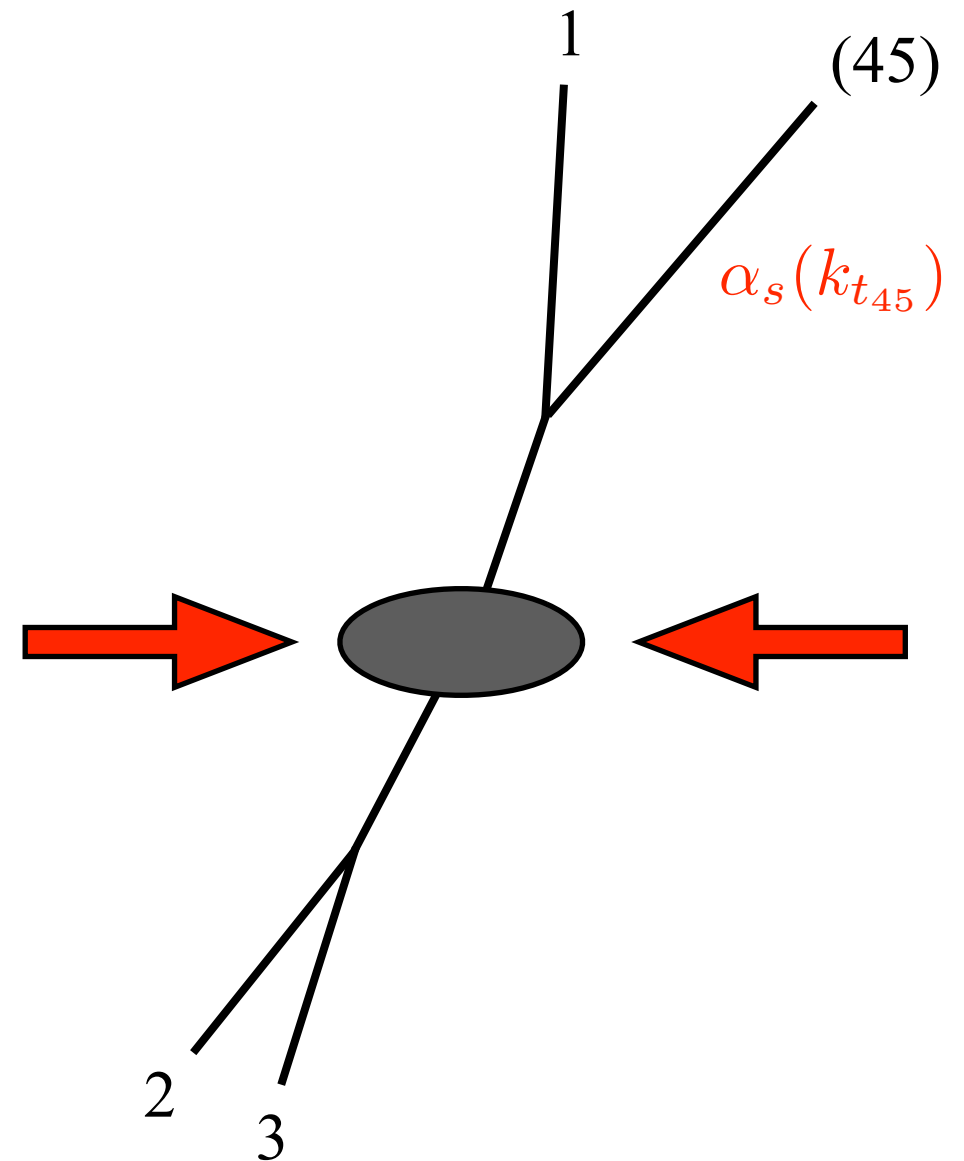
- given a partonic event reconstruct a branching history: cluster partons into jets using k_t -algorithm
- at each branching the scale in the coupling to set to the relative k_t of the daughter partons
- local scale = CKKW scale choice, but no Sudakov reweighting, no parton shower



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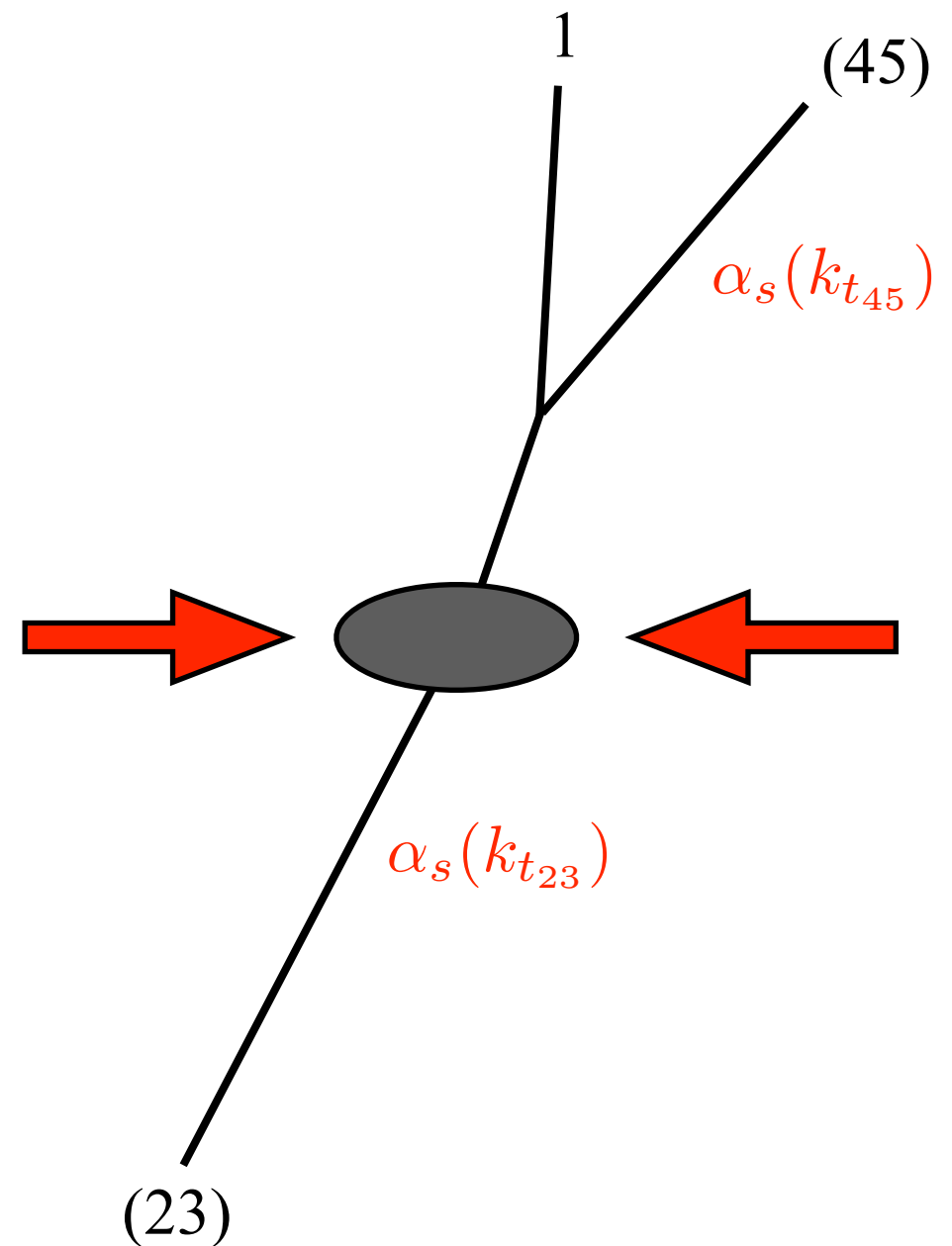
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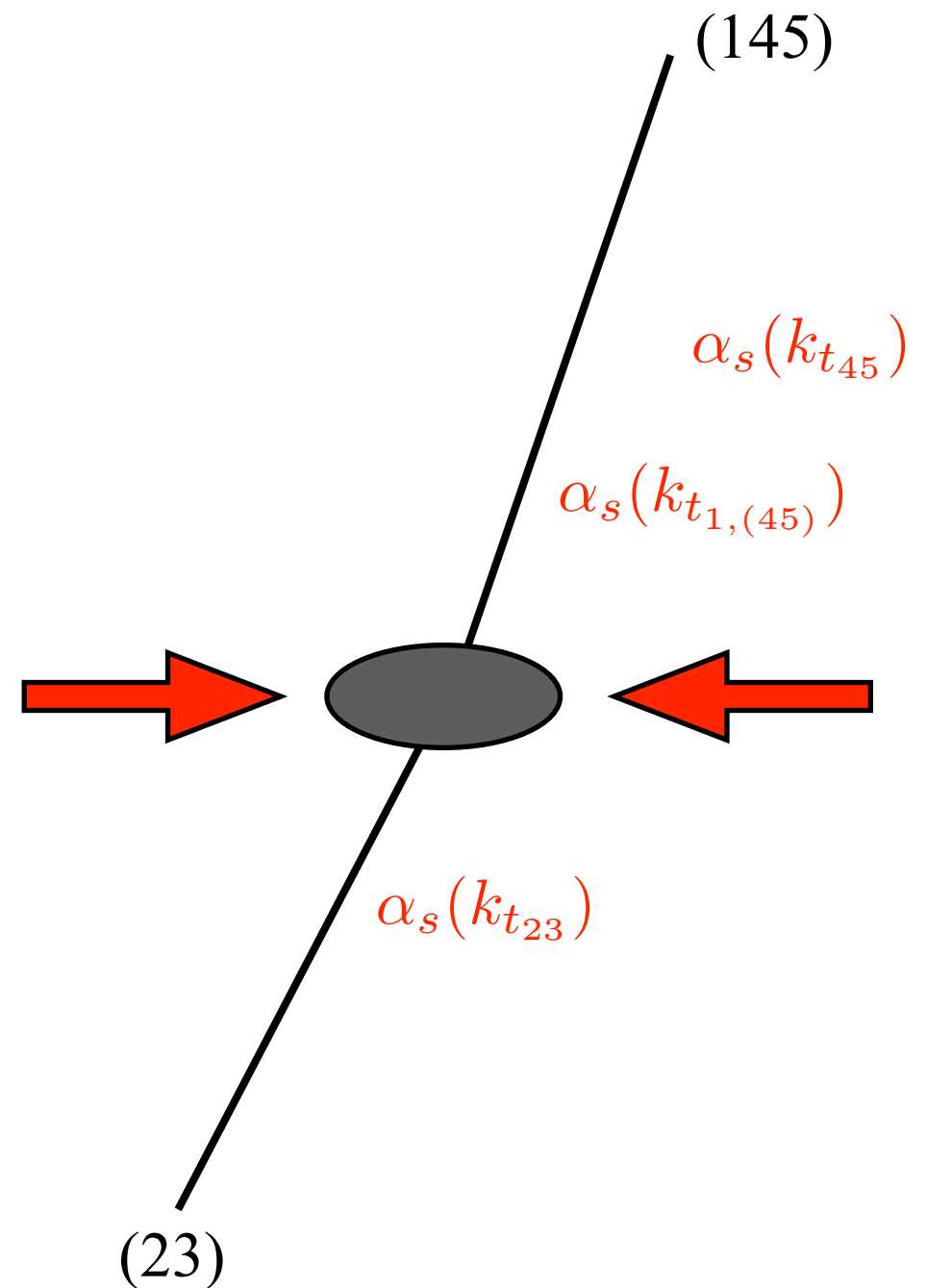
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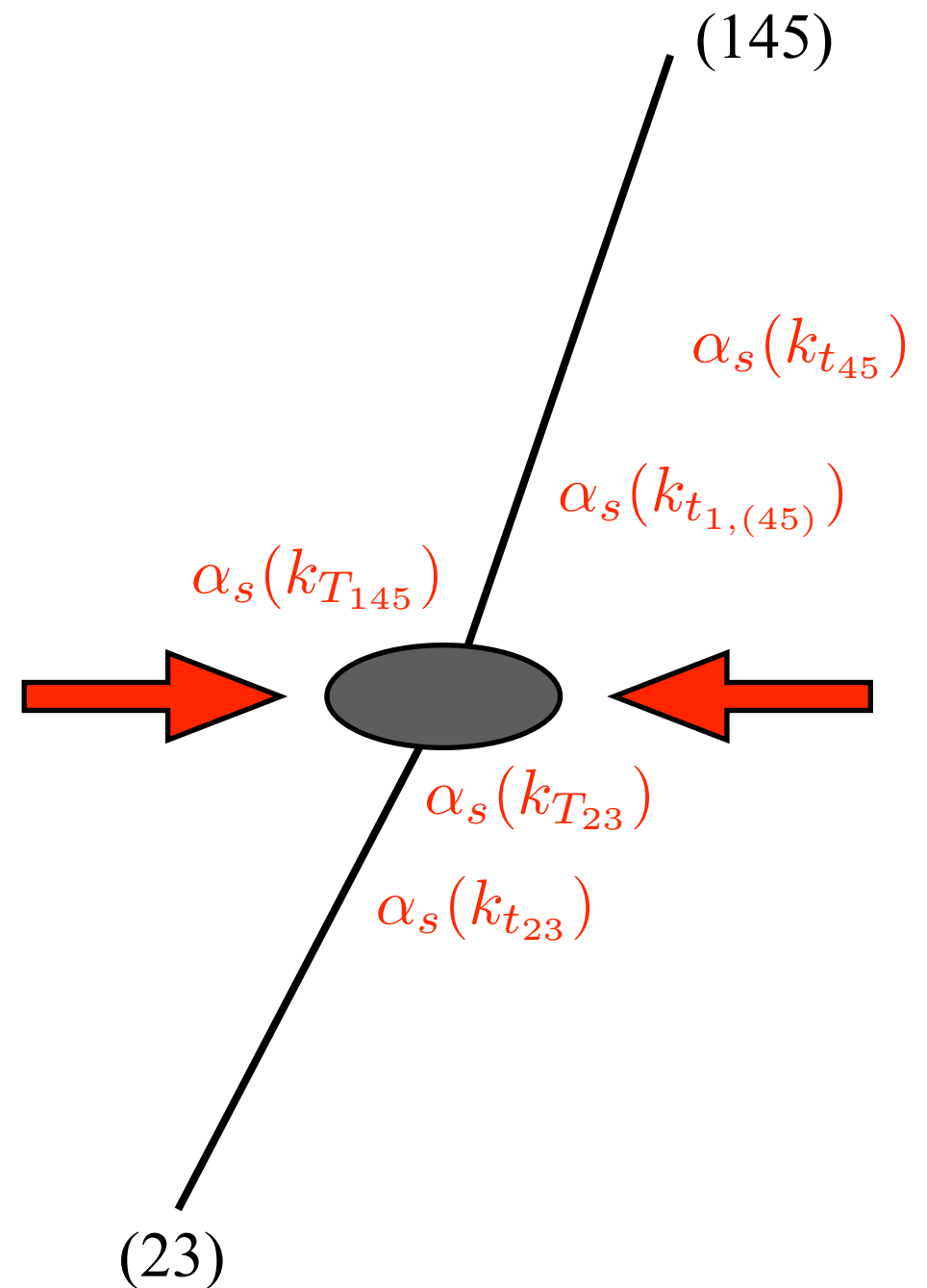
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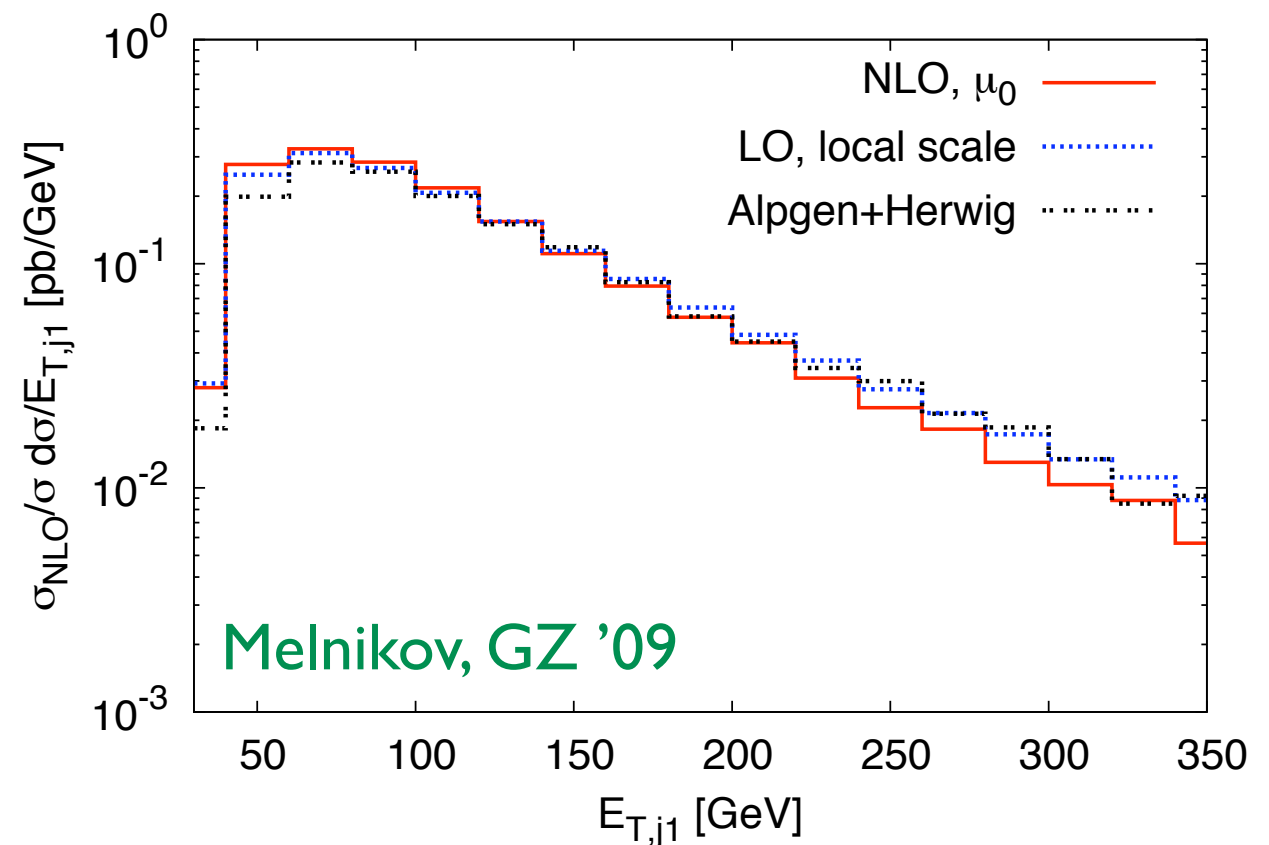
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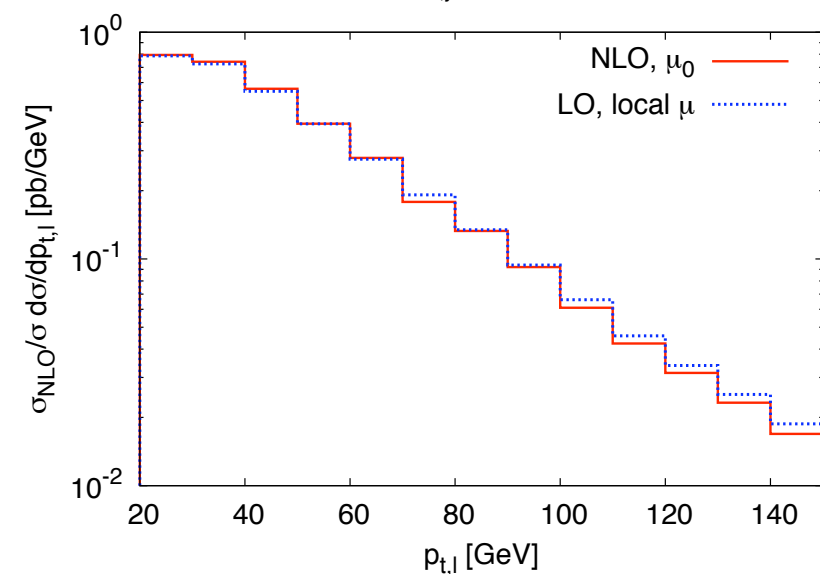
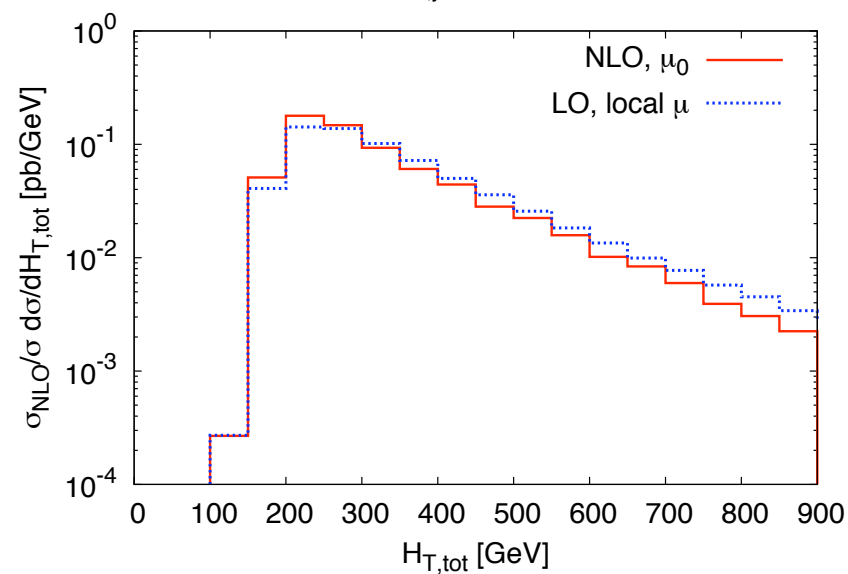
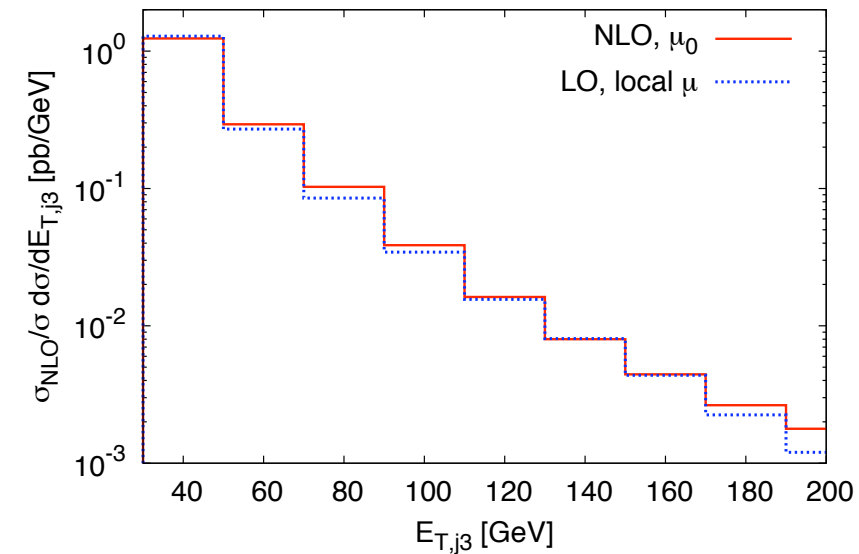
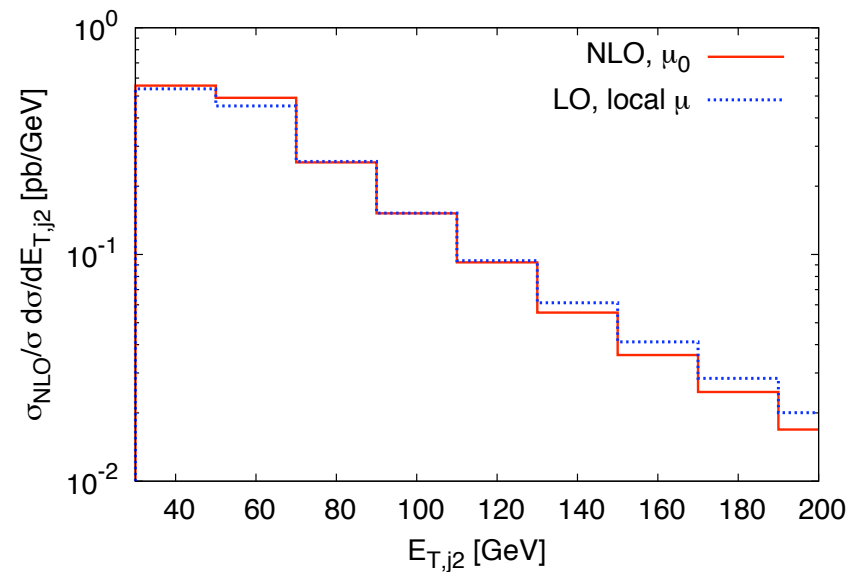
Local (CKKW) scale

Local scale choice (CKKW):

- given a partonic event reconstruct a branching history: cluster partons into jets using k_t -algorithm
 - at each branching the scale in the coupling to set to the relative k_t of the daughter partons
 - local scale = CKKW scale choice, but no Sudakov reweighting, no parton shower
- ➡ local scale choice reproduces the shape of the NLO distribution well
- ➡ the difference between LO with local scale and full Alpgen+Herwig indicative of the importance of the parton shower



Other hadronic distributions

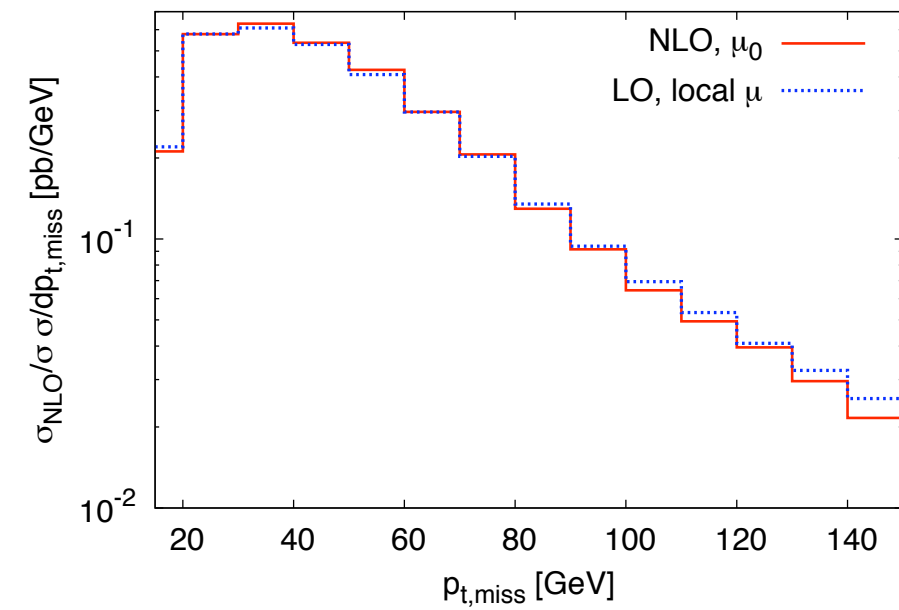
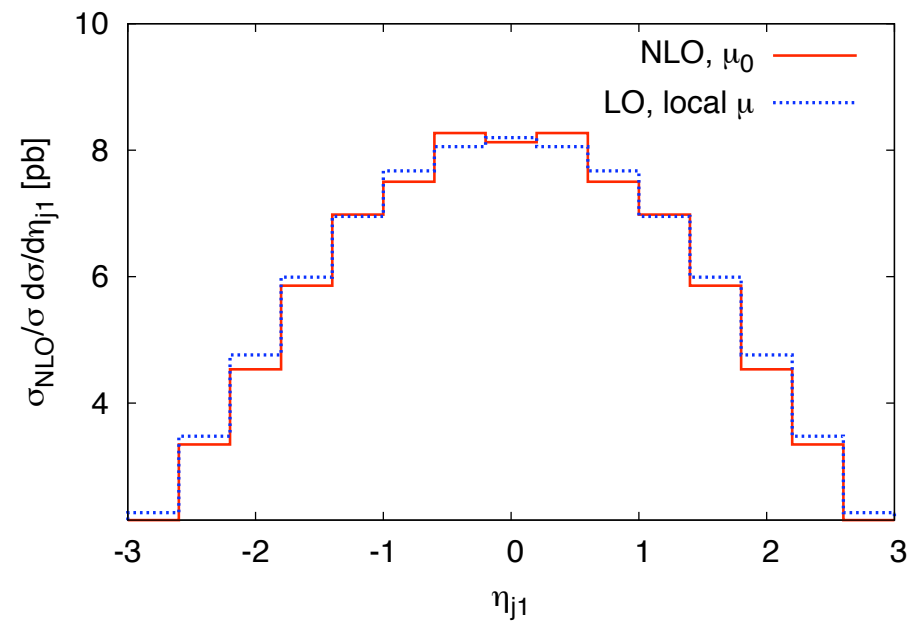


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➡ LO with local scale does a very reasonable job in reproducing shapes

NB:
normalization of LO remains out of control. LO is normalized to NLO in above plots

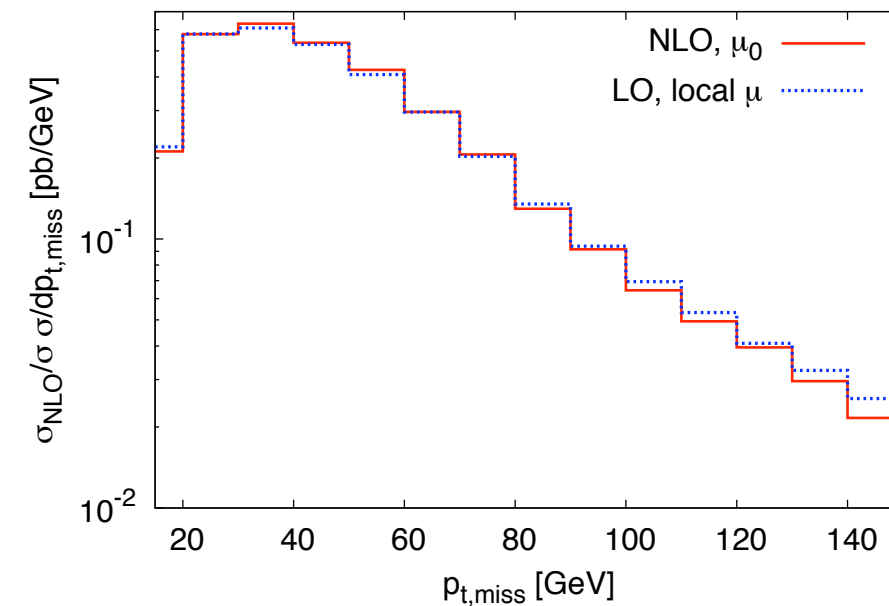
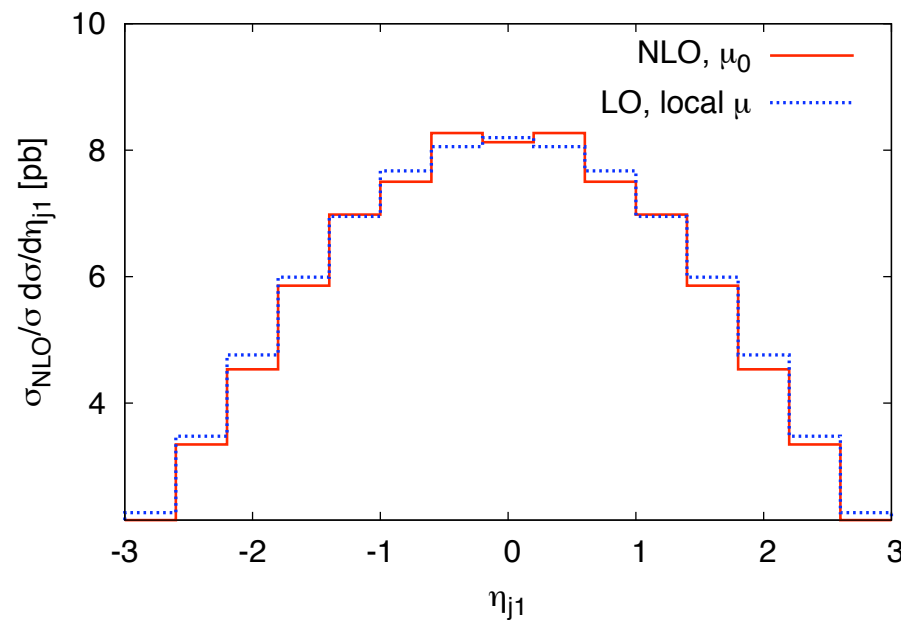
Leptonic distributions



Melnikov GZ '09

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How solid (cut-independent) is this statement ?

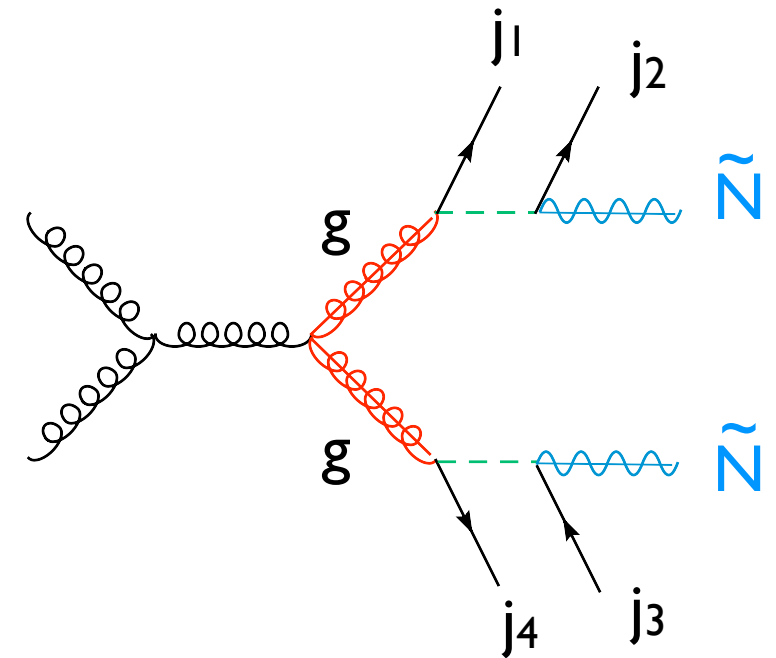
See what happens with different cuts.

Consider two sets of cuts where $W+3\text{jet}$ plays the role of unwanted background

SUSY signature

SUSY with R-parity: e.g. gluino pair production,
each decays into 2 jets and neutralino

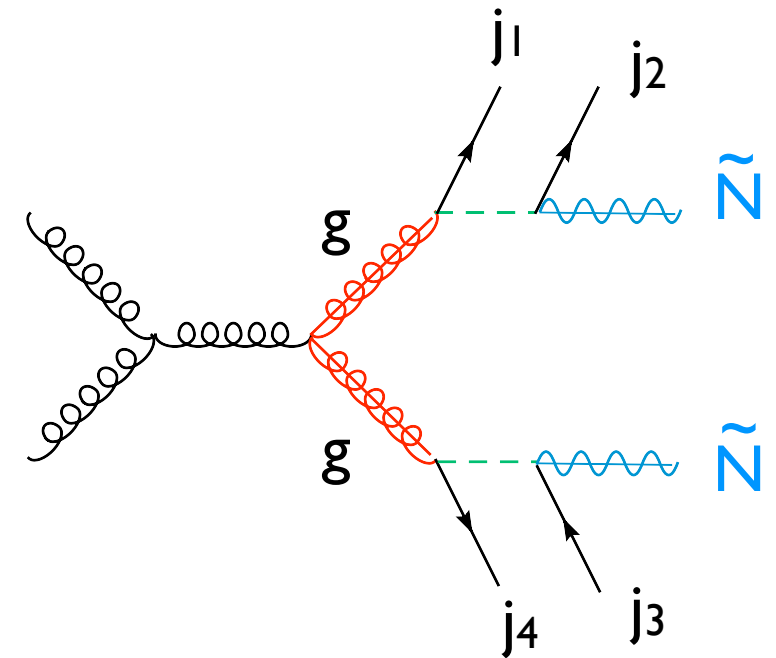
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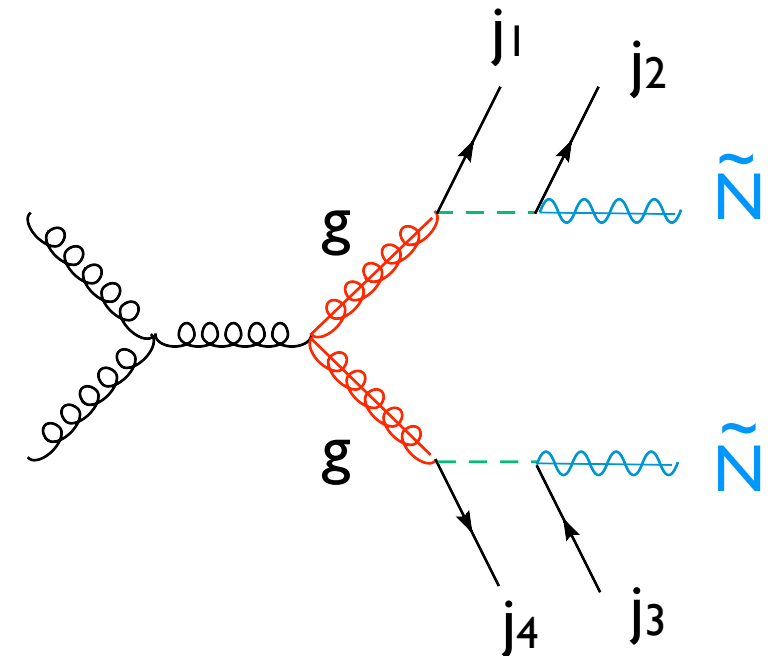


Primary, irreducible background: $Z (\rightarrow \nu\nu) + 4 \text{ jets}$

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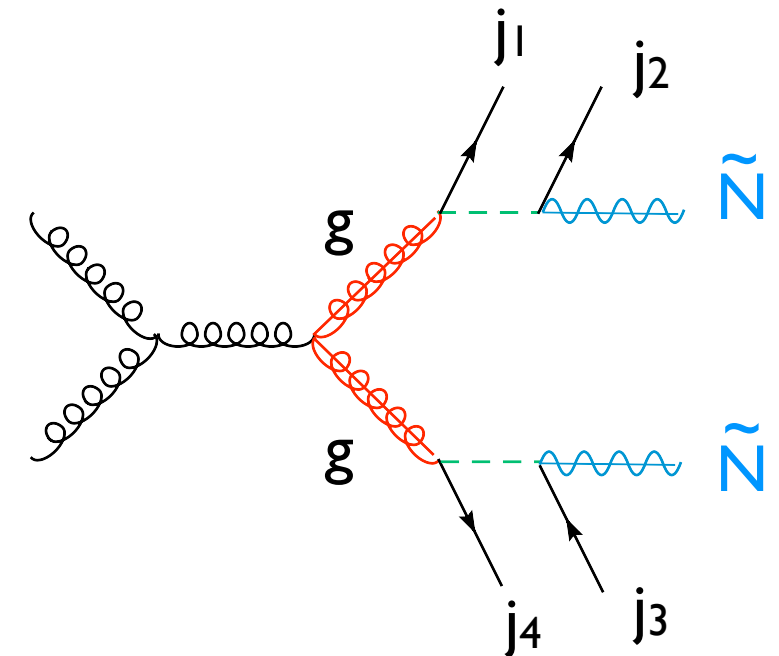
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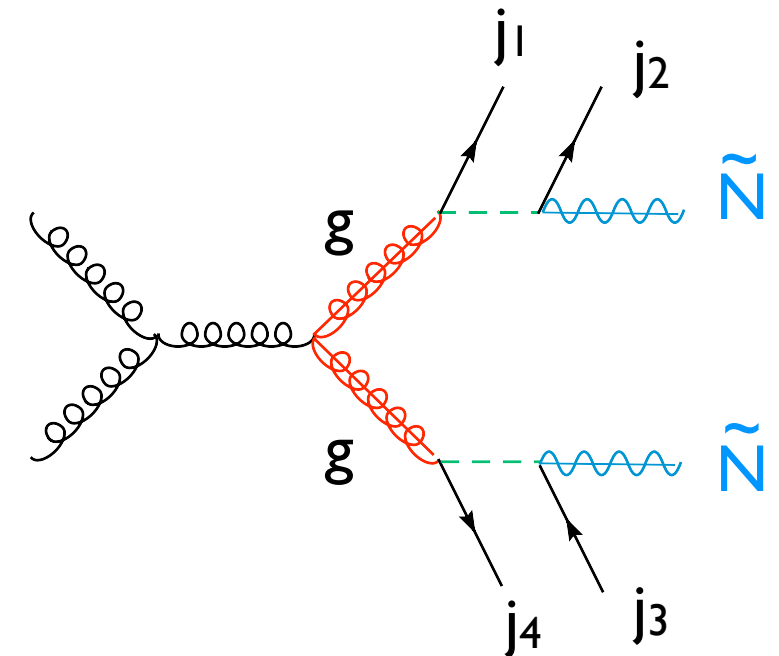
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2) $\sigma(W + 3 \text{ j}) \sim 100 \sigma(Z + 4 \text{ j})$

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\Rightarrow important to consider this source of background as well

Atlas setup

Cuts designed by ATLAS to suppress W+3j background

$$p_{T,j} > 50 \text{ GeV} \quad p_{T,j1} > 100 \text{ GeV} \quad p_{t\bar{t}} < 20 \text{ GeV}$$

$$E_{T,\text{miss}} > \max(100 \text{ GeV}, 0.2 H_T) \quad H_T = \sum_j p_{T,j} + E_{T,\text{miss}}$$

$$S_T > 0.2 \quad |\eta_j| < 3$$

Yamazaki [ATLAS and CMS Col.] 0805.3883

Yamamoto [ATLAS Col.] 0710.3953

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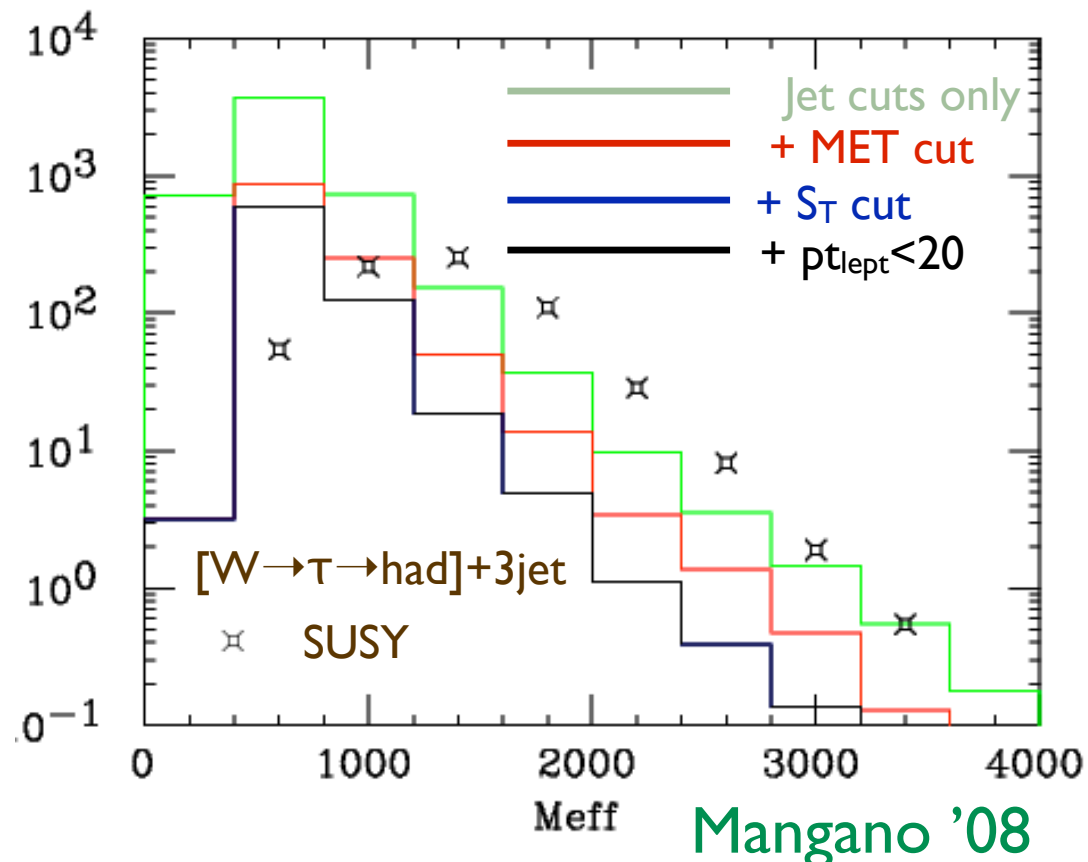
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- each cut suppresses background by factor ~ 3 without modifying the shape
- cut on collinear unsafe sphericity S_T not applied in the following study

SM background from $W+3$ jets

Our calculation includes only the leptonic decay of the W (in e , μ or τ) but not the hadronic subsequent decay of τ . However

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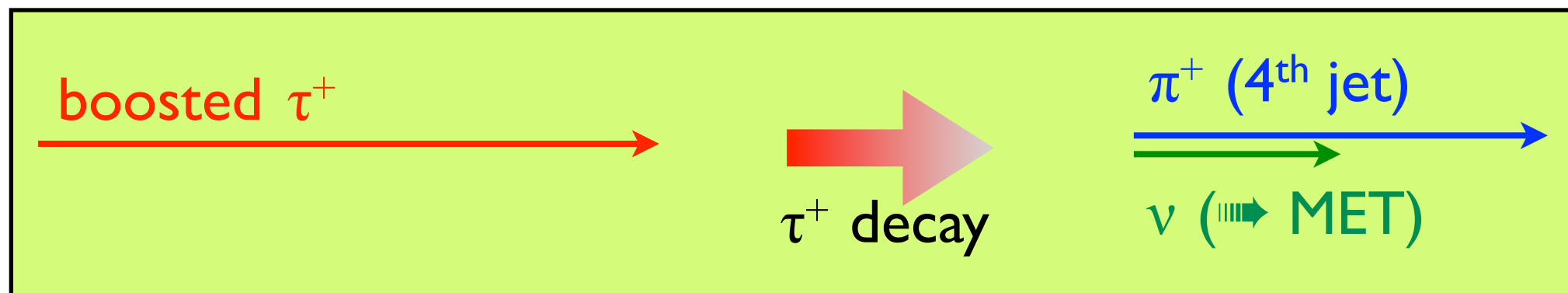
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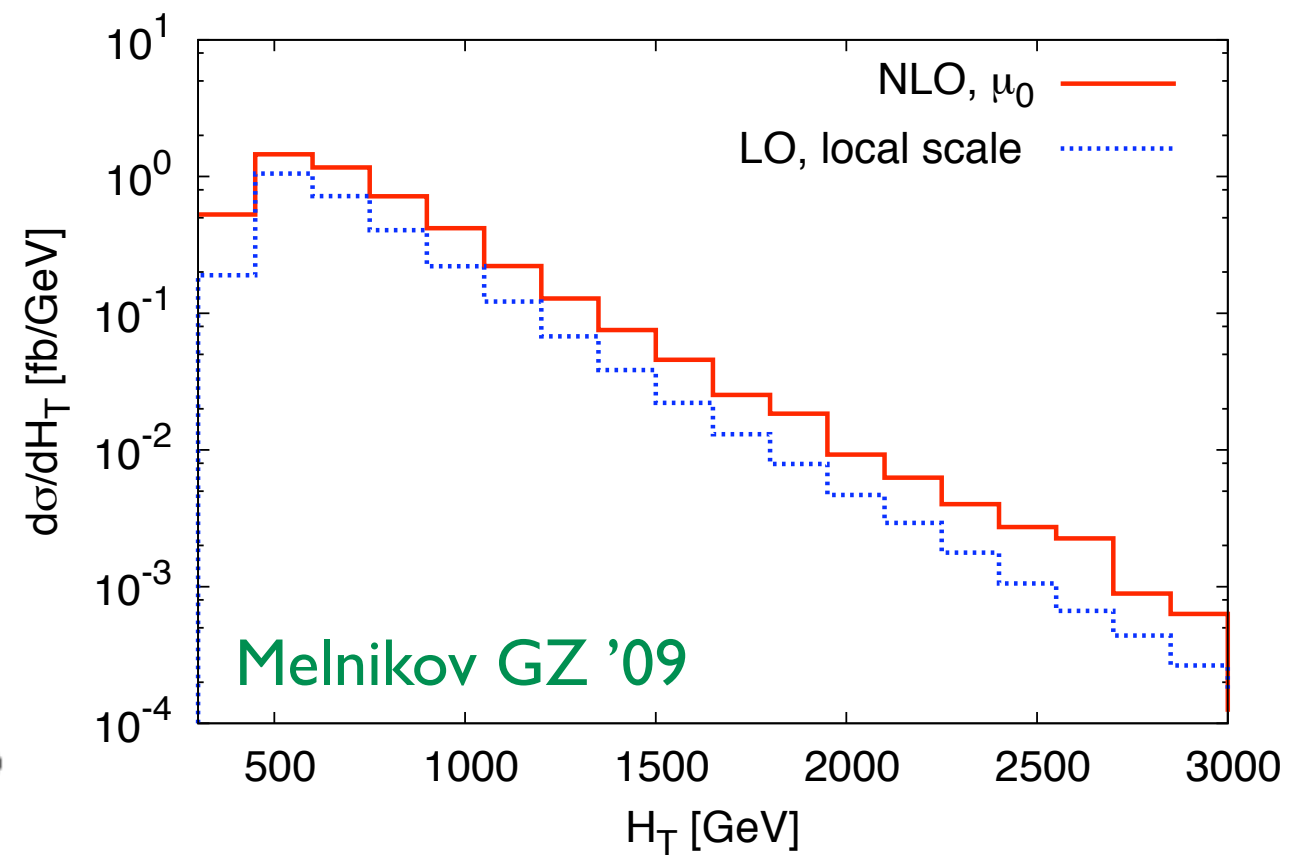
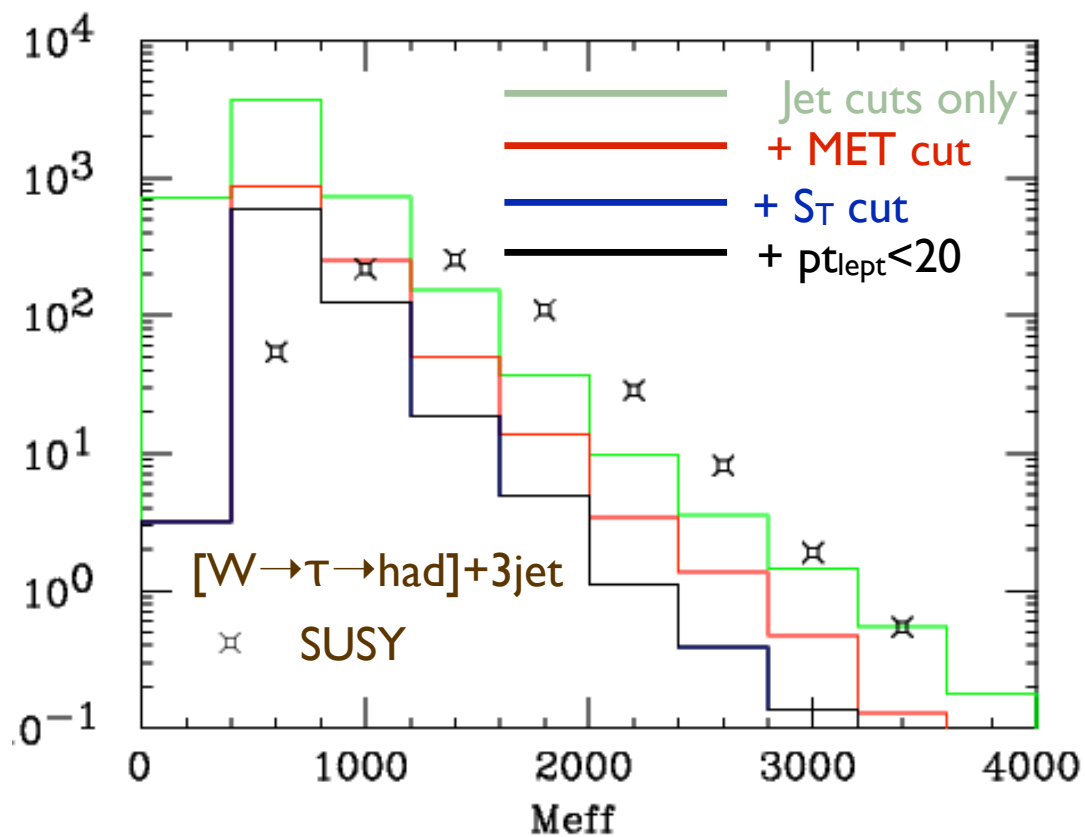
Theoretical robust approximation:

simulate the W decay as a perfect collinear branching with momentum fractions 2/3 (π^+) and 1/3 (ν)

SM background from W+3 jets

Primary observable is H_T (previously called M_{eff}) which ‘measures’ the SUSY scale:

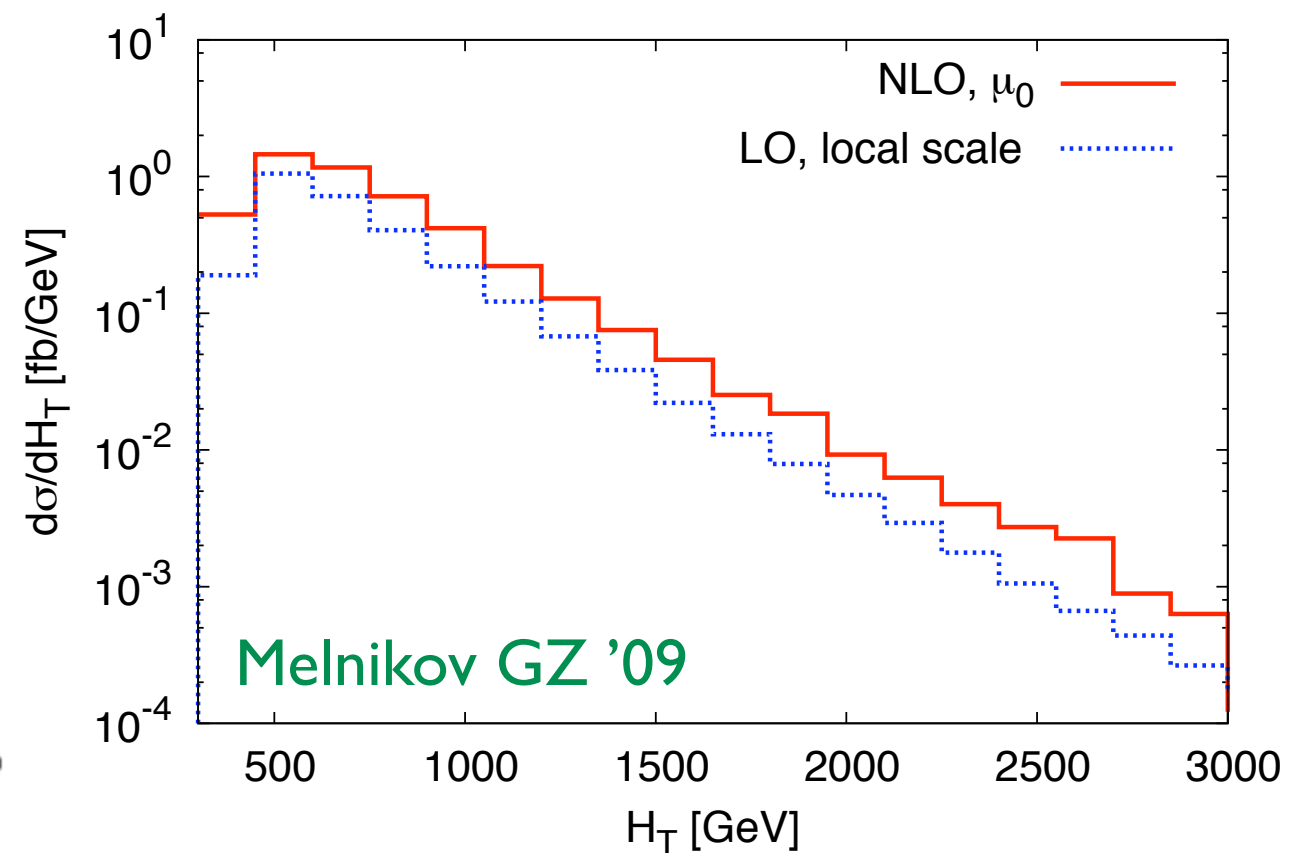
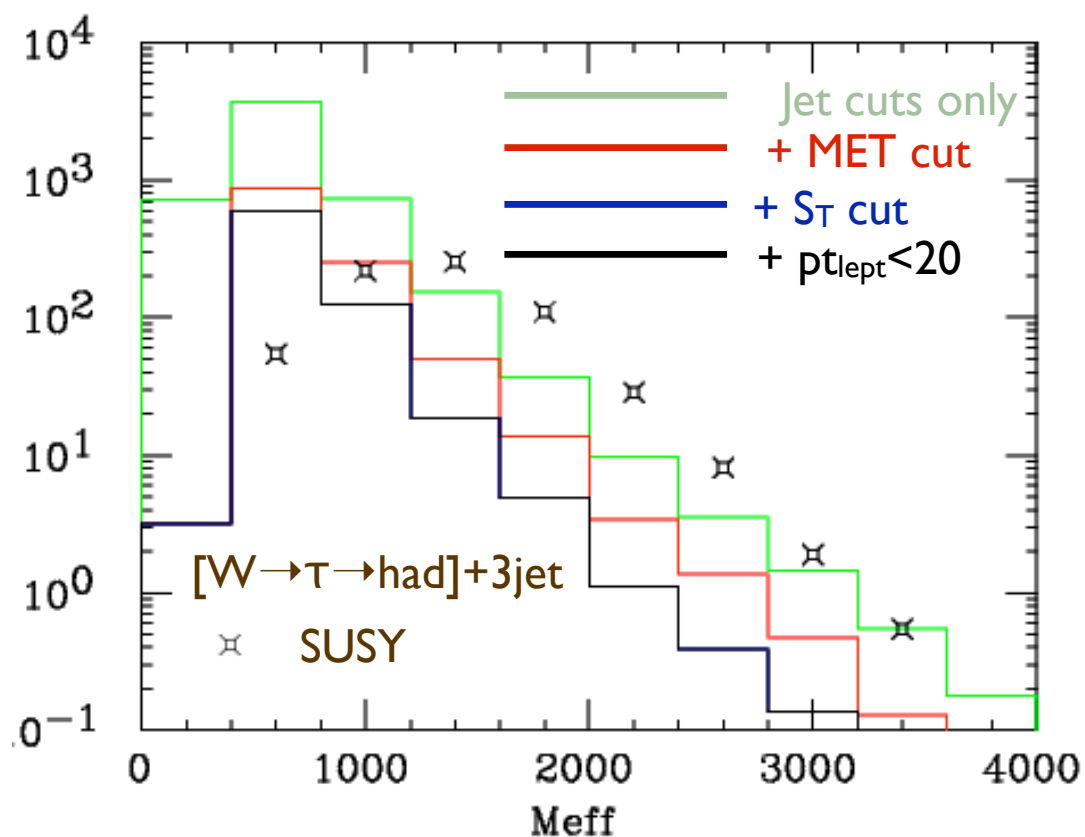
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- ☞ universal enhancement (K-factor ~ 3) of LO without distorting the shape
NB: *same observable* with cuts as shown before had K-factor ~ 1
- ☞ NLO effect similar to that of cuts but *works in opposite direction*

CMS style indirect lepton veto cut

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Indirect lepton veto = no explicit lepton veto, but other cuts force contribution from W +jets to become naturally small

$$p_{T,j} > 30\text{GeV} \quad p_{T,j1} > 180\text{GeV} \quad p_{T,j2} > 110\text{GeV} \quad E_{T,\text{miss}} > 200\text{GeV}$$

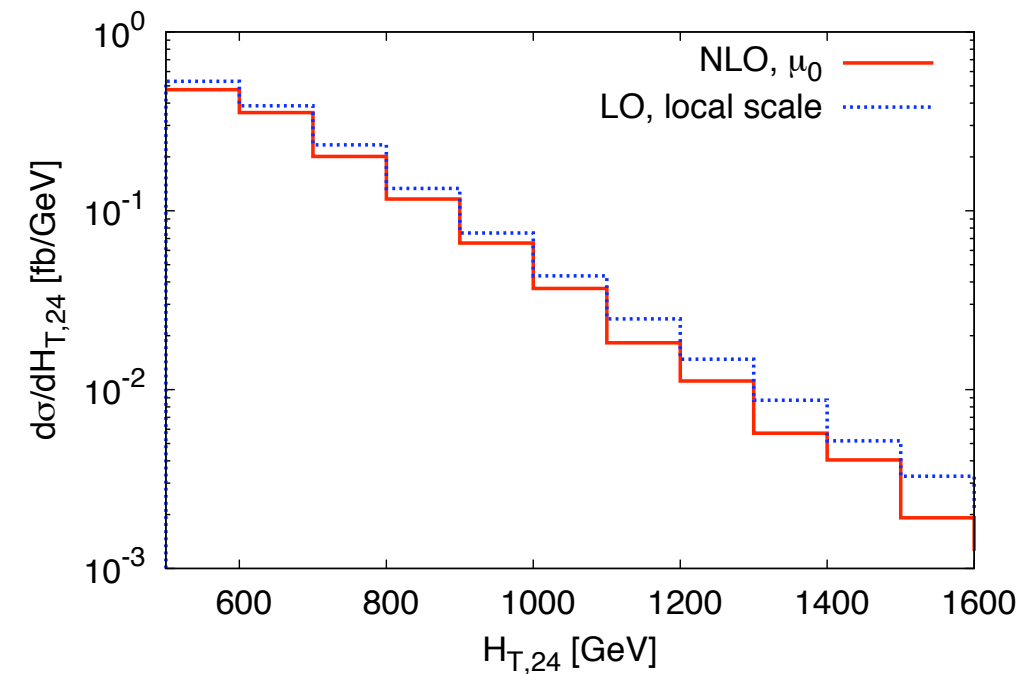
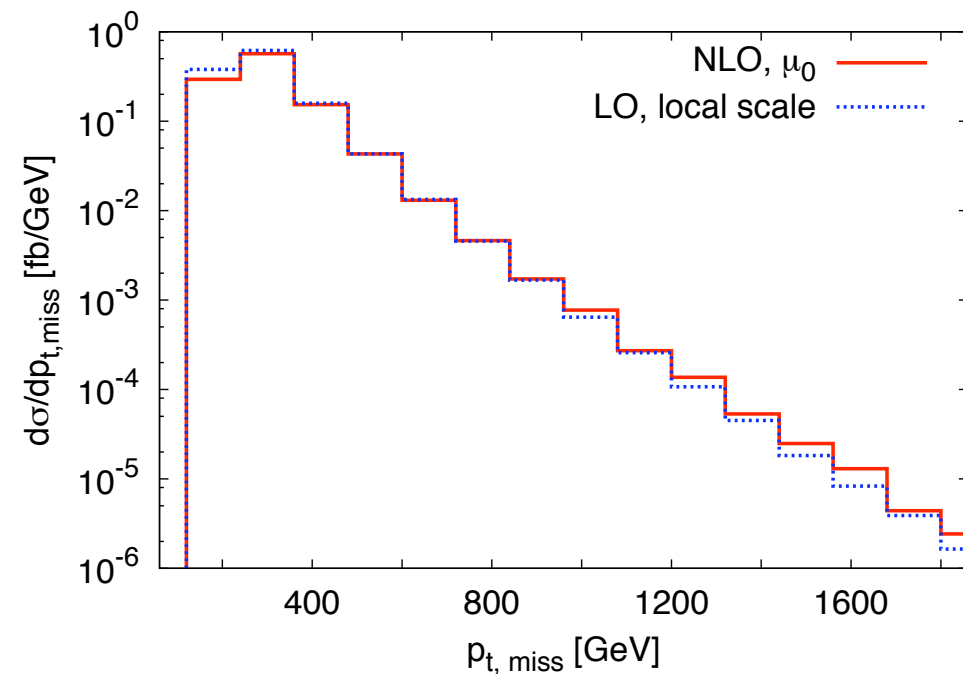
$$|\eta_{\text{lead jet}}| < 1.7 \quad |\eta_{\text{other jets}}| < 3 \quad H_{T,24} = \sum_{j=2}^4 p_{T,j} + E_{T,\text{miss}} > 500\text{GeV}$$

CMS Collaboration Journal Phys. G: Nucl. Part. Phys. 34 (2007) 995

CMS style indirect lepton veto cut

Primary search observables

distribution in transverse missing energy and total effective mass $H_{T,24}$



- NLO correction to cross-section small, K-factor ~ 1
- shapes of LO mostly OK, but moderate shape distortion at high $H_{T,24}$

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- ☑ all this emphasizes the need to extend NLO corrections to other processes ($Z+3j, W+4j \dots$)