

LOW X MEETING  
 ISCHIA ISLAND, ITALY  
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# Heavy Quarkonia: as Seen through the Eyes of Central Exclusive Production at the Tevatron and LHC

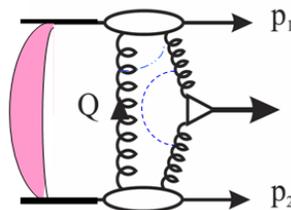


V.A. Khoze (IPPP, Durham, PNPI)



(Based on works with L. Harland-Lang, A.Martin, M.Ryskin and W.J. Stirling)

**main aim:** to demonstrate that CEP can open a new way to study the properties of heavy quarkonia, and, in particular, to serve as a spin-parity analyser

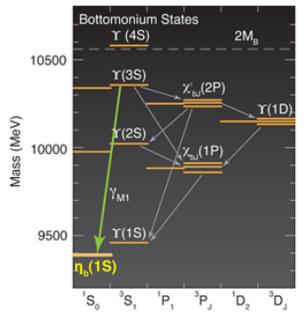


$\chi_{c2}, \chi_{c1}, \chi_{c0}$   
 $\chi_{b2}, \chi_{b1}, \chi_{b0}$



# PLAN

1. Introduction
2. **C**entral **E**xclusive **P**roduction as a heavy meson spin-parity analyser.
3. *Interpretation of the CDF results on charmonium CEP.* (Mike, Jim)
4. *New Run of Durham studies.*
5. *Remarks about the models for soft diffraction.* (Alan, Asher)
6. Conclusion.



# INTRODUCTION

Why an interest to the CEP of  $\chi_c, \chi_b$  ?

- Testing ground for the formalism of CEP used to evaluate the New Physics signals (e.g. 'Diffractive Higgs')
- Open issues in Quarkonium Spectroscopy, such as  $\chi_b$  quantum numbers.  New way to address Quarkonium Physics (numerous new exotic charmonium like states).
- New Encouraging CDF results on CEP of the  $\chi_c$  . (Mike, Jim)

## Heavy Quarkonia

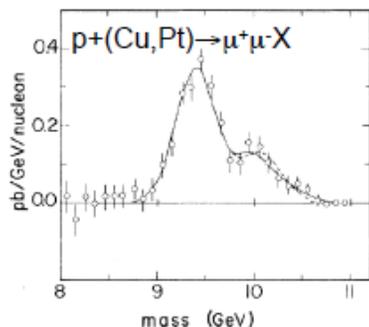
Traditional testing ground for various aspects of QCD

- NRQCD, QCDME, Lattice QCD, QCD sum rules, potential models
- Large NLO..... PT corrections.
- P-states- sensitivity to the derivatives of the wave function, relativistic effects....
- Nature of the new states around 4 GeV; X, Y, Z, other applications of the CEP...

# Bottomonium history started 30 years ago

( PRL 39, 242 (1977) and PRL 39,1240 (1977) )

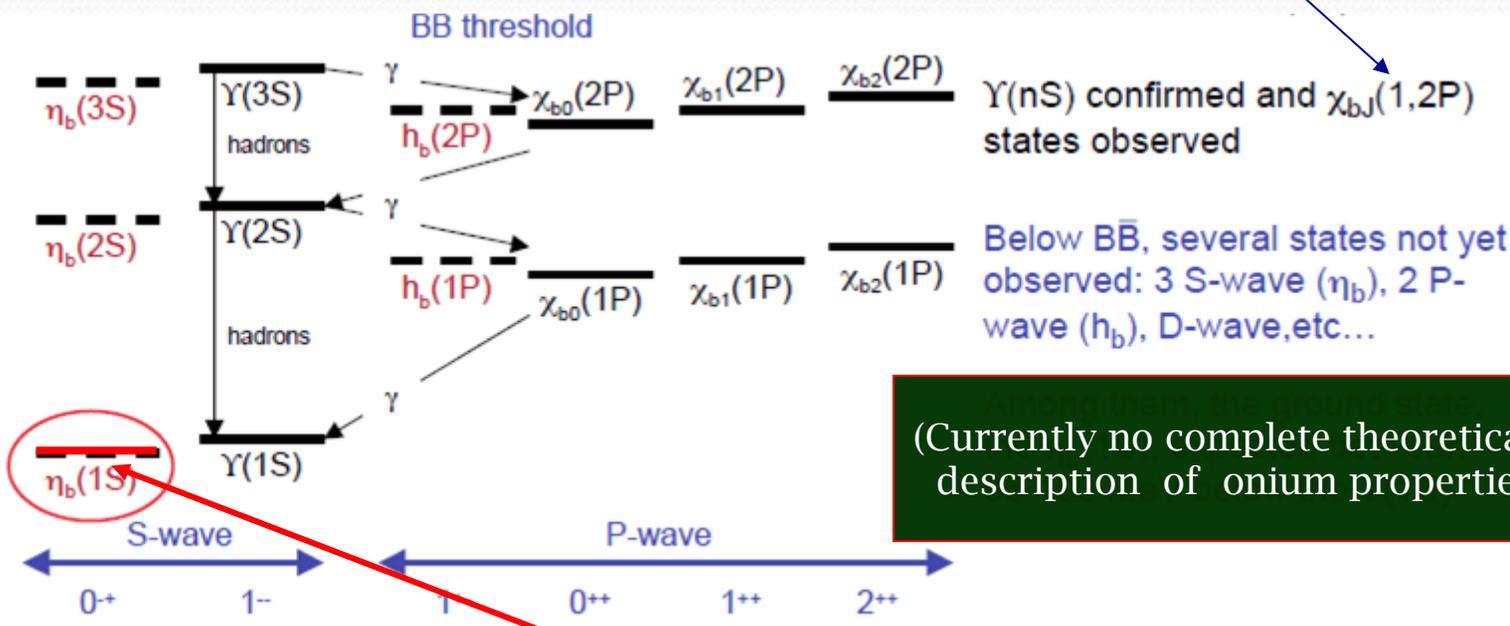
30 years later....



$M(\Upsilon) = 9.40 \pm 0.013$   
 $M(\Upsilon') = 10.00 \pm 0.04$   
 $M(\Upsilon'') = 10.43 \pm 0.12$

FNAL, E288

(spins- still unconfirmed)



(Currently no complete theoretical description of onium properties.)

(BABAR (2008))

(Still puzzles)



The heaviest and most compact quark-antiquark bound state in nature

## The main advantages of the $\chi_c \rightarrow \chi_b$ -CEP



- Quantum number filter/analyser.  
( $0^{++}$  dominance ; C,P-even),
- Clean few-particle final state,
- Favourable background conditions.  
(theoretical estimates,  $\gamma\gamma$ - data).
- New leverage -proton momentum correlations

### Potential (theoretical) problems

- Higher sensitivity to low scales- 'usual suspects'
- Stronger dependence on Enhanced Screening effects



(larger  $s / M_\chi^2$  )

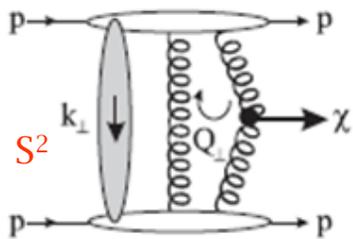
# What we expect within the framework of the Perturbative Durham formalism (KMR-01, KKMR-03, KMRS-04)

O++ -case

$$T = A\pi^2 \int \frac{d^2Q_\perp P(\chi(0^+))}{Q_\perp^2 (\vec{Q}_\perp - \vec{p}_{1\perp})^2 (\vec{Q}_\perp + \vec{p}_{2\perp})^2} f_g(x_1, x'_1, Q_1^2, \mu^2; t_1) f_g(x_2, x'_2, Q_2^2, \mu^2; t_2),$$

$$A^2 = 8\pi\Gamma(\chi \rightarrow gg)/M_\chi^3 \quad *K_{\text{NLO}}$$

$$P(\chi(0^+)) = (\vec{Q}_\perp - \vec{p}_{1\perp}) \cdot (\vec{Q}_\perp + \vec{p}_{2\perp}).$$



● **Strong sensitivity to the polarization structure of the vertex in the bare amplitude.**

■ In the on-shell-gluon approximation **spin -1** is excluded by **Landau-Yang** theorem. Should lead to a strong suppression.

(R. Pasechnik)

■ For forward going protons in the non-relativistic quarkonium approximation, CEP of the **spin-2** meson is strongly suppressed ( $J_z=0$  selection rule)

KMR-01

(A. Alekseev-1958-positronium)

● **Absorption is sizeably distorted by the polarization structure** (affects the b-space distr.)

KMR-02, KKMR-03

●  $\chi_c, \chi_b$  -production is especially sensitive to the effects of enhanced absorption

■ larger available rapidity interval

■ lower scale  $\rightarrow$  larger dipole size  $\rightarrow$  larger absorption

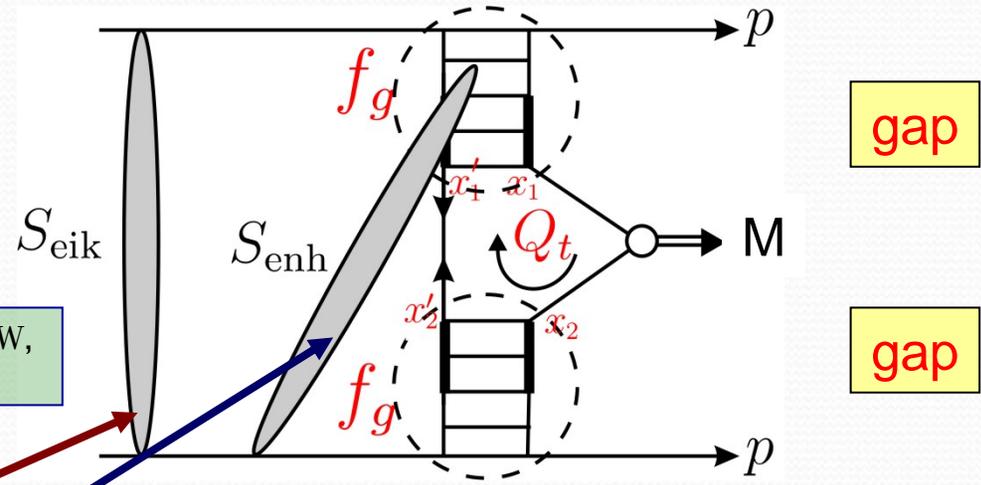
( $S^2\text{enh}$  for  $\chi_c$  at the Tevatron is expected to exceed that for the Higgs at the LHC)

(KMR-08)

“soft” scattering can easily destroy the gaps



$S^2 \rightarrow$  absorption effects -necessitated by unitarity



Everybody's happy (KMR, GLM, FHSW, Petrov et al, BH, GGPS, Luna...MCs)

eikonal rescatt: between protons  
 enhanced rescatt: involving intermediate partons

(Alan,Asher)

(BBKM-06)

Subject of hot discussions :  $S^2_{enh}$



# Interpretation of CDF results

(CDF Collaboration, arXiv:0902.1271, PRL-09)

(Mike, Jim)

Assuming that all events are originated from CEP of  $\chi_c(0^+)$  (limited acceptance)  
(used CHIC MC- Durham based)

CDF  $\frac{d\sigma(\chi_c)}{dy} \Big|_{y=0} = (76 \pm 14) \text{ nb}$

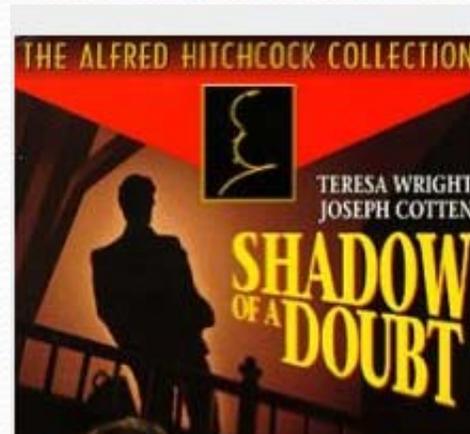
KMRS -2004: 130 nb  $\rightarrow$  80 nb (PDG-2008)



Signal based on:  $\chi_c \rightarrow J/\psi + \gamma$

A certain preference to  $0^+$  in the  $(J/\psi + \gamma)$  mass distribution

Too good to be true ?!



# Devil's Advocate Questions



1. How do we know which particular P- state has been found ?

(P-states are not clearly separated experimentally )

2. Is reconstruction based on CHIC MC still acceptable, what if not  $J^{PC} = 0^{++}$ .

?

3. Are we close to the CEP prescription (role of low mass SD and DD)

Yes, due to the record  
CDF gap coverage  
(KMRS-04)

$\chi_c(0^+)$  dominates CEP, but

$$Br(\chi_c(0^+) \rightarrow J/\Psi + \gamma) = (1.14 \pm 0.11)\%$$

$$Br(\chi_c(1^+) \rightarrow J/\Psi + \gamma) = (34.1 \pm 1.5)\%$$

$$Br(\chi_c(2^+) \rightarrow J/\Psi + \gamma) = (19.4 \pm 0.8)\%$$



- On-mass-shell  $1^+$  production is forbidden due to Landau-Yang theorem, but what about off-mass-shell effects ?

Recently- renewal of interest (R.Pasechnik, A. Szczurek, O.Teryaev-09)

Still numerically small

- Within the non-relativistic framework for forward going protons  $2^+$  is strongly suppressed, but what about non-forward protons and relativistic effects?

Important phenomenon- absorptive corrections are quite sensitive to the meson spin-parity (studied before in the context of scalar/pseudoscalar Higgs-KKMR04)



## New Run of Durham Studies

### Issues addressed:

- New **SUPERCHIC MC** for all  $c\bar{c}$  P -states.
- Absorption effects for CEP of the  $0^+, 1^+, 2^+, 0^-$   $c\bar{c}$  - states revisited
- Proton angular correlations for different  $0^+, 1^+, 2^+, 0^-$   $c\bar{c}$  -states.
- Expectations for the CEP of the  $0^+, 1^+, 2^+, 0^-$   $b\bar{b}$  -states.

### As compared to the previous **K(KMR)S** studies:

- More comprehensive calculation of the absorption effects using the new **KMR-07/08** model for soft diffraction (including the enhanced screening).
- New calculational routine for implementing polarization structure in the b-space.
- New experimental/theoretical results for the parameters of heavy quarkonia, in particular  $\Gamma(\chi \rightarrow gg)$ .

# CHARMONIUM

- The final state (muon) distributions in the  $(J/\psi + \gamma)$  system are sensitive to the meson spin, but after imposing the CDF cuts this dependence is strongly reduced. (L. Harland-Lang, W.J. Stirling).
- Cross-section 'reconstruction' is safe 🤔
- Spin is not discriminated via the  $(J/\psi + \gamma)$  decay products 🤔
- We need to measure better spin-parity analysing final state:  $\pi\pi, KK, p\bar{p}$  or outgoing proton momentum correlations KMRS-04

For normalization purposes- scalar case at the Tevatron

$$\langle S_{eik}^2(0^+) \rangle \approx 0.065$$

$$\langle S_{eff}^2(0^+) \rangle \approx 0.02$$

$$(\Delta=2.3)$$

$$d\sigma(0^+)/dy|_{y=0} \approx 80 \text{ nb}$$

$$d\sigma(0^+)/dy|_{y=0} \approx 30 \text{ nb}$$

- Still within Durham approach-uncertainties, recall, in particular  $(f_g)^4$ - effect

Reasons to believe that enhanced absorption is overestimated- KMR-09

Some of the 'typical' uncertainties cancel in the ratios

$$\diamond \sigma(1^+) / \sigma(0^+) \simeq \langle p_t^2 \rangle / M_\chi^2 * \langle S_{eik}^2(1) \rangle / \langle S_{eik}^2(0) \rangle * R_{NLO}^1 \simeq 0.05$$

■  $\langle S_{eik}^2(J) \rangle$  - the same for all J within ~20% accuracy,

■  $\langle S_{eik}^2(1, 2) \rangle / \langle S_{eik}^2(0) \rangle \simeq 2.5$ ;  $\langle p_t \rangle \simeq 0.5$  GeV.

Production is more peripheral: zero at  $\vec{b} = 0$ , where the absorption is largest  
at larger  $b$  absorption is small anyway.

$$\diamond \sigma(2^+) / \sigma(0^+) \simeq (\langle p_t^2 \rangle / Q_t^2)^2 * \langle S_{eik}^2(2) \rangle / \langle S_{eik}^2(0) \rangle * R_{NLO}^2 \simeq 0.05$$

(  $\langle Q_t \rangle \simeq 1 \text{ GeV}$  )

❖  $\sigma(0^+ \rightarrow J/\psi + \gamma) : \sigma(1^+ \rightarrow J/\psi + \gamma) : \sigma(2^+ \rightarrow J/\psi + \gamma) = 1 : 1.3 : 0.8$

Then  $d\sigma(0^+, 1^+, 2^+) / dy \simeq 95$  nb, as compared to experiment:  $(76 \pm 14)$  nb

After all, 🙌😊🙌 keeping in mind all uncertainties

## Energy Dependence

- Expected to be weak, since the rise of the gluon density at low  $x$  is compensated by stronger enhanced screening.

$$d\sigma(\chi_{c0})/dy \approx 50 \text{ nb at the LHC}$$

- Test of the enhanced absorption (less model dependent):

$$\sigma(\chi_{LHC}) / \sigma(\chi_{Tevatron})$$

various uncertainties cancel (NLO effects, width,...)

- $\sigma(\chi_{c2}) / \sigma(\chi_{c0})$  decreasing with energy ( $\langle Q_t \rangle$  increasing)

## Momentum correlations between outgoing protons

- Separation of different meson states (irrespective of the final state), seen in the CEP of light mesons (WA102 Collab).

Recall the results of Regge theory at low transverse momenta  $(p_{3,4})_t$  (KKMR03)

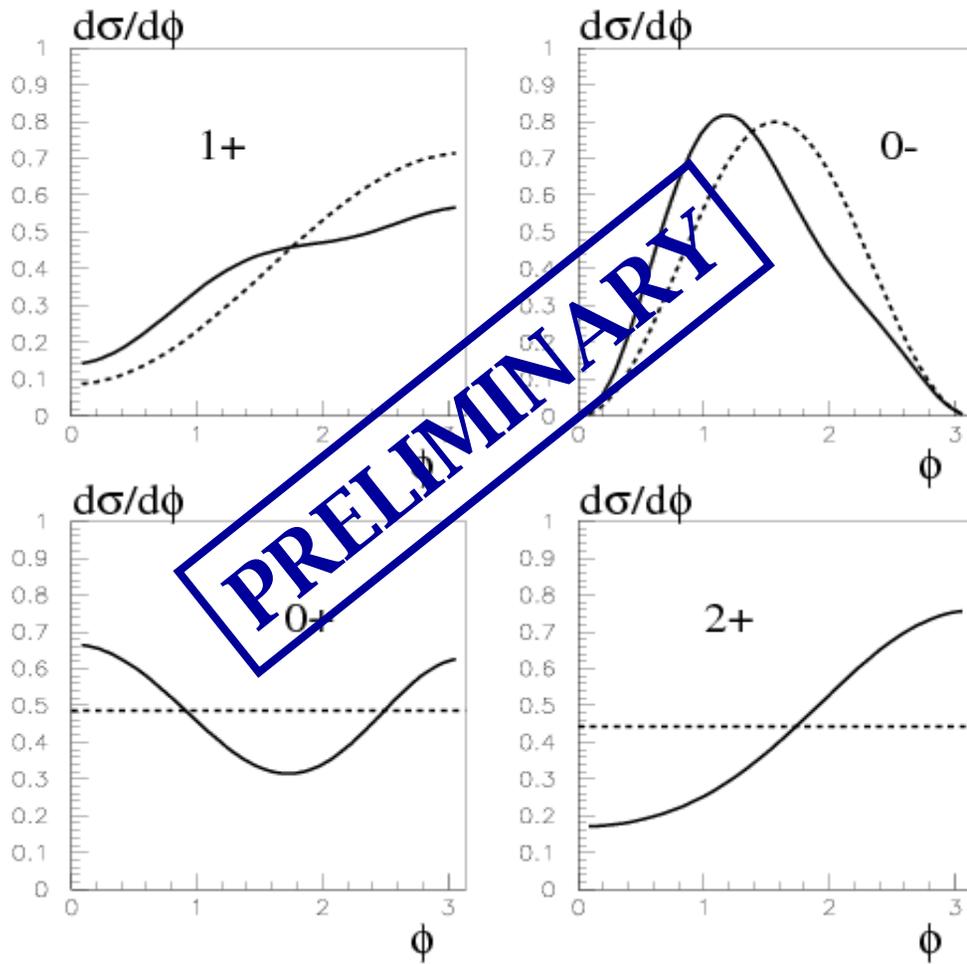
$$d\sigma(0^+) / d\phi \simeq \text{const};$$

$$d\sigma(0^-) / d\phi \simeq \sin^2 \phi;$$

$$d\sigma(1^+) / d\phi \simeq (\vec{p}_3 - \vec{p}_4)_t^2;$$

$$d\sigma(2^+) / d\phi \simeq \text{const};$$

- Some dependence on the choice of unintegrated gluon densities (KMR-02, KKMR-03)
- Serious modification by the absorption effects (amount of suppression strongly depends on the impact parameter  $\vec{b}$ )



# BOTTOMONIUM

Recall , the decay  $\chi_{b0}(1P) \rightarrow Y(1S)\gamma$  has not been seen (yet)

but  $Br(\chi_{b0}(2P) \rightarrow Y(2S)\gamma) = (4.6 \pm 2.1)\%$



According to evaluation by J.T. Lavery et al (2009)

$$\Gamma(\chi_{b0} \rightarrow gg) = 3.7 \text{ MeV} ,$$

$$Br(\eta_b \rightarrow \gamma\gamma) = 3.4 * 10^{-5} \quad (\text{Exp. } Br(\eta_{c0} \rightarrow \gamma\gamma) \approx 2 * 10^{-4} ).$$

$\chi_b$

Higher scale  $\rightarrow$  better PT description

Smaller role of relativistic effects, better knowledge of gluon densities

- $1^+$  -is practically filtered out (strong  $M^2$  -suppression)
- $2^+$  weak-  $\langle Q_t \rangle^2$  dependence
- Enhanced absorption- weaker ( $\sim 2$  times)

## Expectations for $\chi_b, \eta_b$

$$d\sigma / dy |_{y=0} \text{ (in pb)}$$

	$\chi_{b0}$	$\chi_{b1}$	$\chi_{b2}$	$\eta_b$
<b>Tevatron</b>	400	3	10	5
<b>LHC</b>	700	7	20	15

## PROSPECTIVE MEASUREMENTS

- A clear way to resolve the issue of  $\chi_c$  spin-parity identification will be to search for the two-body decays:

$$Br(\chi_{c0} \rightarrow \pi\pi, K^+K^-) \simeq 1.3\% \quad \chi_{c1}, \eta_c \not\rightarrow \pi\pi, KK \quad Br(\chi_{c2} \rightarrow \pi\pi, K^+K^-) \simeq 0.3\%$$

$$Br(\chi_{c0} \rightarrow p\bar{p}) \simeq 2 * 10^{-4} \quad Br(\chi_{c1} \rightarrow p\bar{p}) \simeq 6.6 * 10^{-5} \quad Br(\chi_{c2} \rightarrow p\bar{p}) \simeq 6.7 * 10^{-5}$$

$$Br(\eta_c \rightarrow p\bar{p}) \simeq 0.13\%$$

- Tagged forward protons: spin-parity ID of old and new heavy meson states, detailed tests of absorption effects
- With sufficient statistics of  $\gamma\gamma$  CEP, the measurement of the ratio
 
$$\sigma(\chi_b) / \sigma(\gamma\gamma)$$
 can be quite instructive (the same mass range, various uncertainties cancel).

# UNCERTAINTIES



## Known Unknowns

- N(N)LO- radiative effects (K-factors etc..)  
'...possible inadequacy of PT theory in  $\alpha_s$  ...' R.Barbieri et al-1980
- 'Right' choice of gluon densities, in particular at so low scales as in the  $\chi_c$  case  
(potentiality of a factor of  $\sim 3$  rise for the H-case).
- Complete model for calculation of enhanced absorption.
- $\chi_b$  -experimental widths, decays...

## Unknown Unknowns

- Non- pQCD effects in the meson characteristics.  
Currently no complete description of heavy quarkonium characteristics.  
'Two gluon width does not tell the whole story.'
- Gluons at so low scales, surprises are not excluded at all.



Factor of 5 up or down  
(at best)

S<sup>2</sup>



**Far more** theoretical papers than the expected number of the CED produced Higgs events

‘Well, it is a **possible** supposition  
‘You think so, too?’  
‘I did not say a **probable** one’



# Selection Criteria for the Models of Soft Diffraction

(Aliosha, Asher, Uri, Sergey)

- We have to be **open-eyed** when the soft physics is involved. Theoretical models in the strong coupling regime contain various assumptions and parameters.
- Available data on soft diffraction **at high energies** are still fragmentary, especially concerning the (low mass) diffractive dissociation.



## A viable model should:

- incorporate the inelastic diffraction :SD, DD (for instance 2-3 channel eikonal of **KMR** or **GLM(M)**)
- describe all the existing experimental data on elastic scattering and SD ,DD and **CED** at the Tevatron energies and below (**KMR**; **GLM(M)** )
- be able to explain the existing CDF data on the HERA-Tevatron factorization breaking and on the CED production of the di-jets, di-photons,  $\chi$ ,  $J/\psi$ ,  $Y$ ..., lead. neutrons at HERA
- provide testable **pre-dictions** or at least **post-dictions** for the Tevatron and HERA

So far (only) **Durham** model has passed these tests.

Only a large enough data set would impose the **restriction order** on the theoretical models and to create a confidence in the determination of  $S^2$ .

Program of Early LHC measurements (**KMR**)



LET THE DATA TALK !

# Who's Afraid of the Big, Bad Wolf?

$S^2$  does not affect the signal-to-background ratio- for all irreducible backgrounds (signal evidence is much less affected).

Overlap background  psec (not lifetime of theor. pred<sup>ns</sup>, but **FTD** resol<sup>n</sup>)

Main reduction of the signal (factor of  $\sim 50$ ) comes from the experimental requirements (cuts and efficiencies...) which are currently known mainly for the inclusive environment. Further progress with hard/soft -ware for the **CEP** processes can be expected.

More experimental work needed.

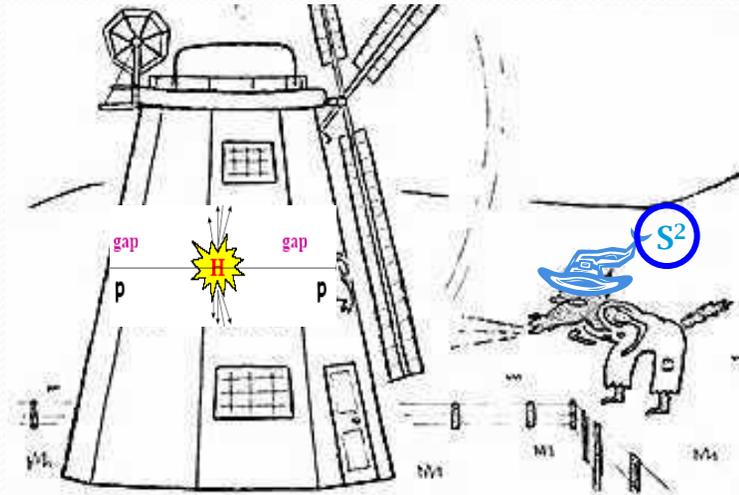


Experimentally we have not seen (at least so far) any evidence in favour of large enhanced absorption (**KKMR, KMR- 2001-2009**).

5

Durham selection of the **UPDF** is quite conservative. Due to the  $(f_g)^4$  behaviour- rise up to a factor of 3 (**Cox et al, KMR**).  
New studies underway.

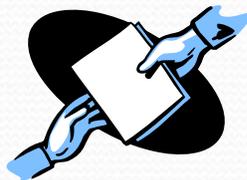
We should be careful with relying on the NLO corrections (e.g. **BBKM-06**). Could be misleading when large parameters are involved. (textbook example: non-relativistic Coulomb corrections)



Up to two orders of magnitude rise in the popular BSM Higgs models.

## CONCLUSION

- CDF data on CEP of the  $\chi_c$  are in a broad agreement with the Durham predictions.
- CEP of heavy mesons - a new way to study quarkonium spectroscopy as well as to address the physics of absorption; can help to establish the nature of newly discovered (exotic) heavy states.
- Promising prospects of studying heavy meson CEP, especially with tagged forward protons.
- Currently active studies are still in progress (both in theory and experiment).





**BACKUP**

## Meson Summary Tables

in the 2008 Review of Particle Physics

<b><math>\chi_{b0}(1P)</math></b> <sup>[d]</sup>	$I^G(J^{PC}) = 0^+(0^{++})$ <i>J</i> needs confirmation		
Mass $m = 9859.44 \pm 0.42 \pm 0.31$ MeV			
<b><math>\chi_{b0}(1P)</math> DECAY MODES</b>	Fraction ( $\Gamma_i/\Gamma$ )	Confidence level	$p$ (MeV/c)
$\gamma \Upsilon(1S)$	<6 %	90%	391

<b><math>\chi_{b1}(1P)</math></b> <sup>[d]</sup>	$I^G(J^{PC}) = 0^+(1^{++})$ <i>J</i> needs confirmation		
$\gamma \Upsilon(1S)$	(35±8) %		423

<b><math>\chi_{b2}(1P)</math></b> <sup>[d]</sup>	$I^G(J^{PC}) = 0^+(2^{++})$ <i>J</i> needs confirmation		
Mass $m = 9912.21 \pm 0.26 \pm 0.31$ MeV			
<b><math>\chi_{b2}(1P)</math> DECAY MODES</b>	Fraction ( $\Gamma_i/\Gamma$ )		$p$ (MeV/c)
$\gamma \Upsilon(1S)$	(22±4) %		442

# What we know from Regge theory (KKMR-2003)

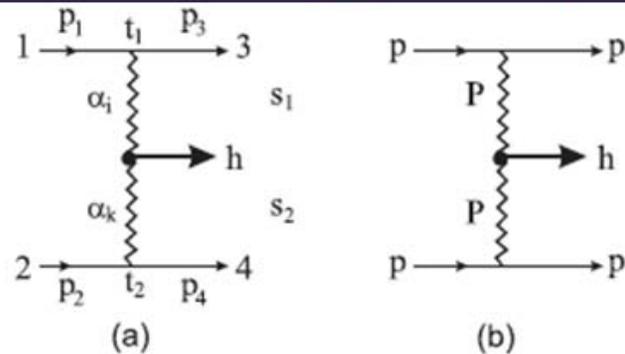


Figure 1: (a) The central production of a state  $h$  by double-Reggeon exchange. (b) The double-Pomeron exchange contribution to  $pp \rightarrow p + h + p$ , which dominates at high energies, where the + signs are used to indicate the presence of Pomeron-induced rapidity gaps.

$$J^P(h) = 0^+$$

Vertex Coupling  $g_{ik}^h = f_{0^+}(p_{3\perp}^2, p_{4\perp}^2, \vec{p}_{3\perp} \cdot \vec{p}_{4\perp})$  depends on dynamics )

$$J^P(h) = 0^-$$

$\sigma \sim |t_1| |t_2| \sin^2 \phi$  observed for  $\eta, \eta'$  by WA102 Group  
(450 GeV,  $pp$  CERN Omega Spectrometer.)

$$J^P(h) = 1^+$$

For small  $p_{it}$

$$g_{PP} = a_{\lambda=0} (p_{3t}^2 - p_{4t}^2) / M^2 (\vec{p}_{3t} \times \vec{p}_{4t}) \vec{e} + f_{\lambda=1} (\vec{K} \times \vec{n}) \vec{e} / M, \quad \vec{K} = (\vec{p}_3 - \vec{p}_4)$$

- Cross section tends to zero at low  $K_t$
- Dominantly produced in the helicity-one state
- Coincide with the NCVC model expectation by F. Close et al (1999)
- Agree with the WA102 data on  $f_1(1420)$  and  $f_1(1285)$

for completeness

$$d\sigma(\eta_c)/dy|_{y=0} = 0.5 \text{ nb at the Tevatron}$$

1.5 nb at the LHC

note: some decay modes are quite sizeable (  $\rho\rho$ ,  $KK\pi$ ,  $\eta\pi\pi$ .. a few % level)

Spin-parity is not confirmed (yet) for some new mesons

- $X(3872)$ :  $1^{++}$  or  $2^{-+}$  ?
- $Y(4140)$  (decay  $J/\psi\phi$ ) (molecule or tetra-quark or..?)

- $X(3940)$ ,  $Y(3940)$  and  $Z(3930)$   $J^{P+}$   
 $\chi_{c1}(2P)$   $\chi_{c2}(2P)$   
or  
 $\eta_c(3S)$  or ?

