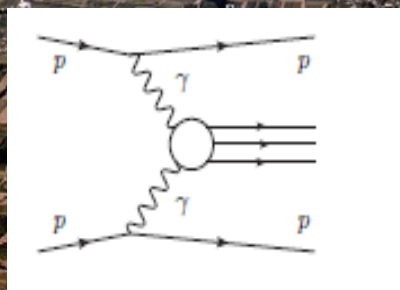
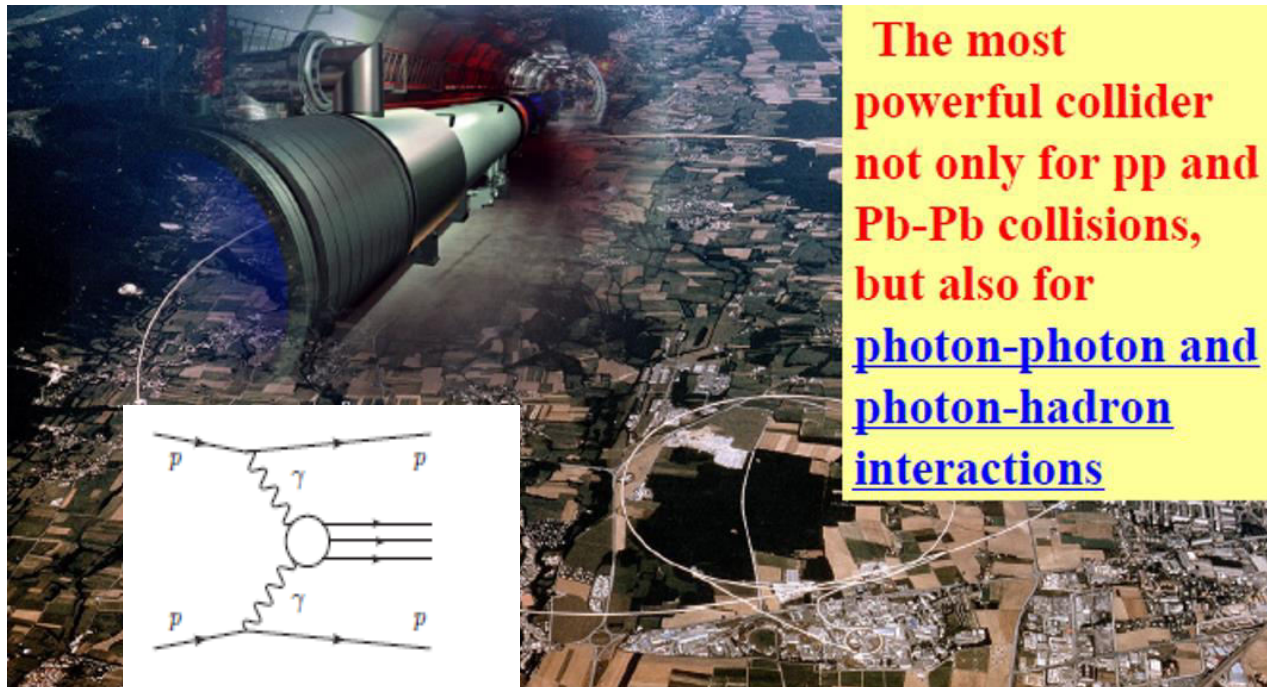




## PHOTON-INDUCED PROCESSES AT HIGH ENERGY HADRON COLLIDERS (selected topics)



Valery Khoze (IPPP, Durham, UK & PNPI, St.Petersburg)





- Introduction and Motivation.
- Photon-Photon collisions . 
- SuperChic MC and Survival Guide
- $\gamma\gamma$  collisions at the LHC- Applications (with an emphasis on BSM physics).
- Other Topical Examples.
- Summary and Outlook. 



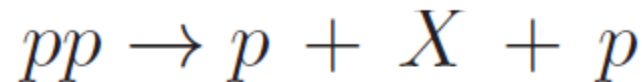
Disclaimer

Emphasis reflects my own personal interests/bias as much as anything.

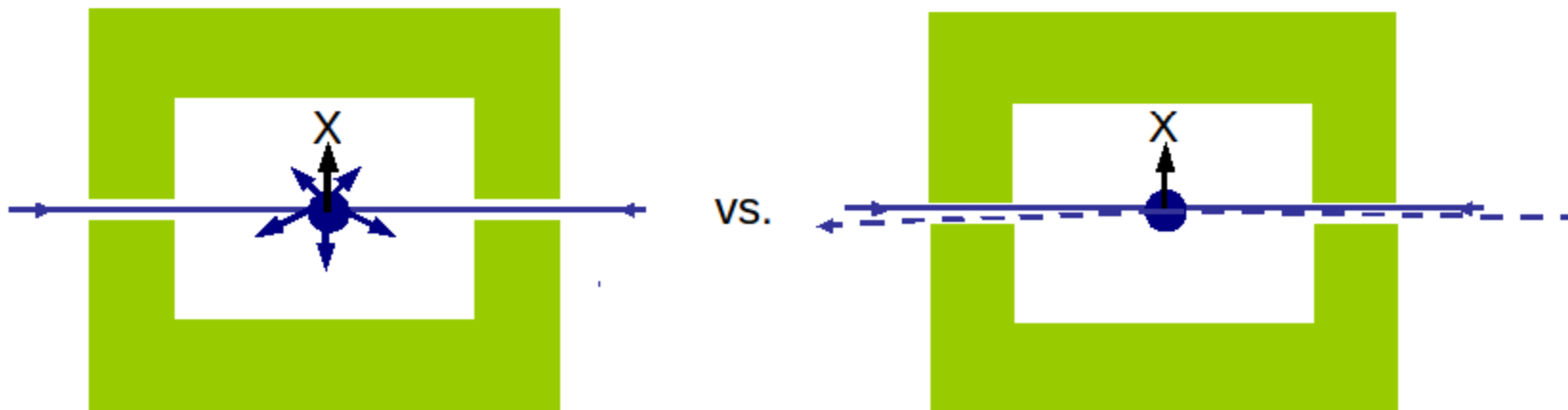


## What is it?

Central Exclusive Production (CEP) is the interaction:



- **CEP** colour singlet exchange between colliding protons, with large rapidity gaps ('+') in the final state. Photons, Pomerons..
- **Exclusive**: hadron lose energy, but remain intact after the collision.
- **Central**: a system of mass  $M_X$  is produced at the collision point and only its decay products are present in the central detector.



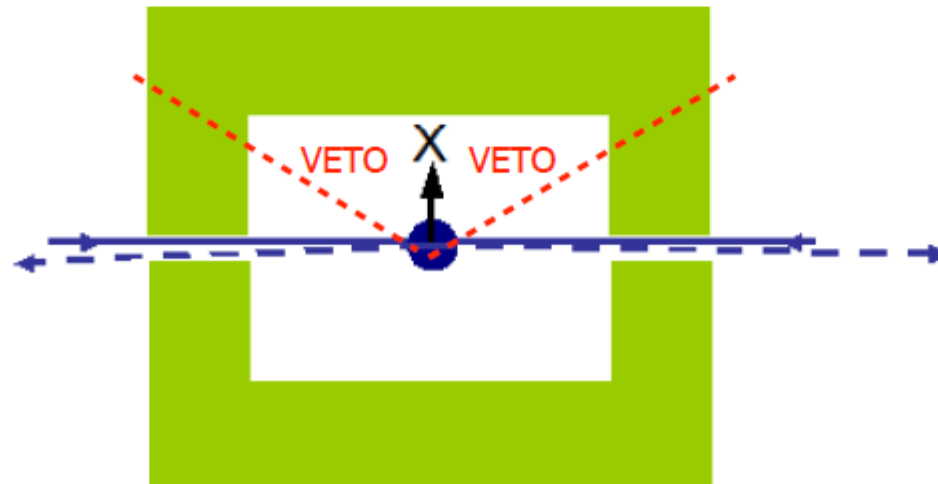
# SELECTING EXCLUSIVE PHOTON-PHOTON EVENTS AT THE LHC



## 1) Gap-based selection: no extra activity in large enough rapidity region.

---

- ▶ No guarantee of pure exclusivity - BG with proton breakup outside veto region. Large enough gap  $\Rightarrow$  BG small and can be subtracted.
- ▶ Pile-up contaminating gap? Either: low pile-up running (dedicated runs/LHCb defocussed beams) or can veto on additional charged tracks only (already used to select charged -  $l^+l^-$ ,  $W^+W^-$  -by ATLAS/CMS/LHCb).

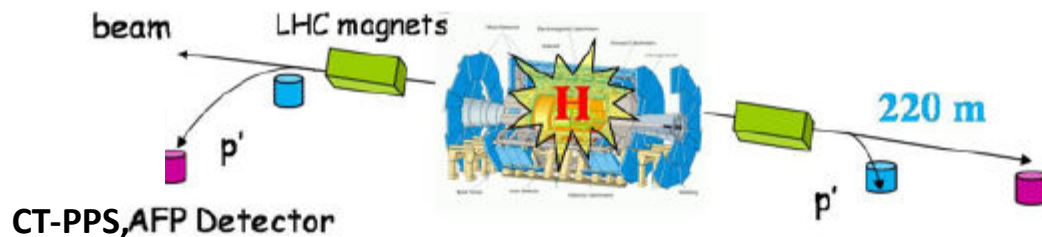


2) Proton tagging:  $pp \rightarrow p + X + p$

- Defining feature of exclusive events: protons intact after collision,
  - If we can measure the outgoing protons, we can identify purely exclusive event samples.

- Basic principle: use LHC magnets to bend outgoing protons out of beam line. After interaction protons are bent out of beam line.
- Ins... (mm) from beam line and  $O(100\text{ m})$  from IP. Reconstruct momenta and measure arrival time of protons.

**The LHC is best mass spectrometer**



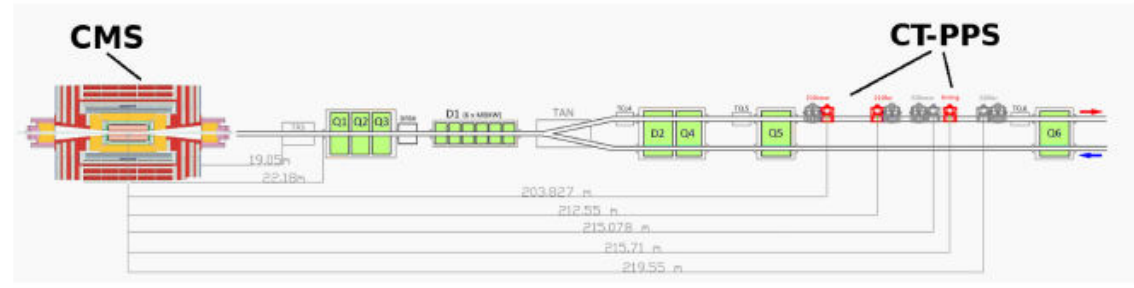
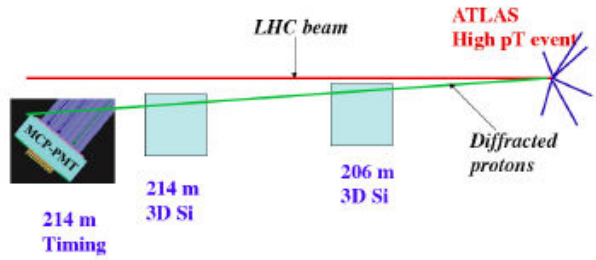
420m under discussion

(Two arms-to come very soon PPS-CMS)

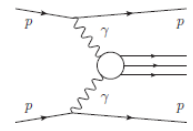
(PPS, AFP first measurements with the one arm proton)



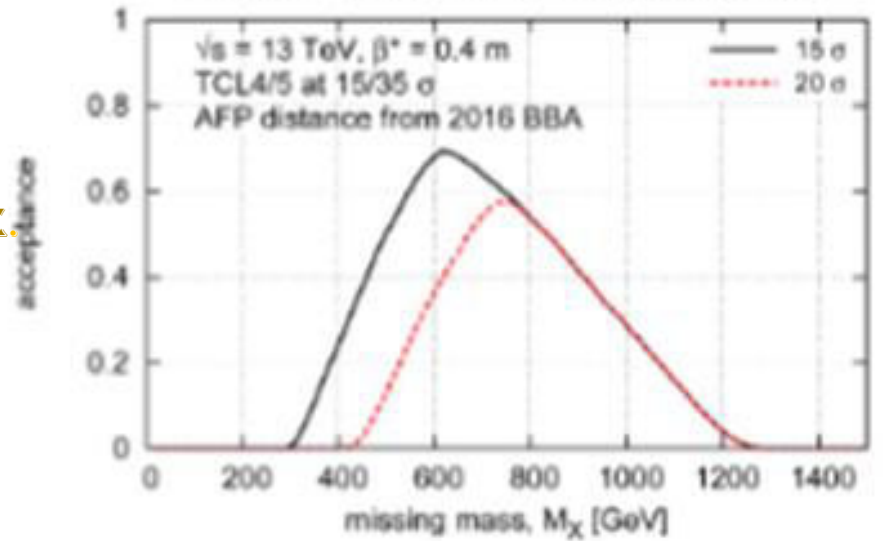
# Proton detectors in CMS-TOTEM/ATLAS: today's configuration



- Tag and measure protons at  $\pm 210$  m both in ATLAS (AFP) and CMS-TOTEM (PPS)
- Detectors: measure proton position (3D pixel or strip Silicon detectors) and time-of-flight (Ultrafast Si, diamond, quartz timing detectors)
- About  $100 \text{ fb}^{-1}$  of data have been accumulated by each experiment ( $\sim 110$  for PPS, 80 for AFP)
- Many analyses in progress



AFP mass acceptance for double tagged events



● THOU SHALT NOT ANNOY MACHINE PEOPLE.



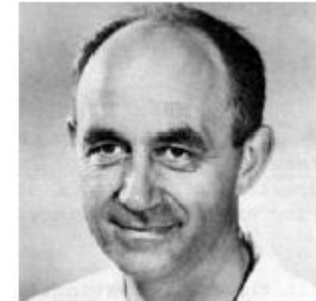
3)

# Ultra Peripheral HI Collisions

Nuovo Cim.,2:143-158,1925

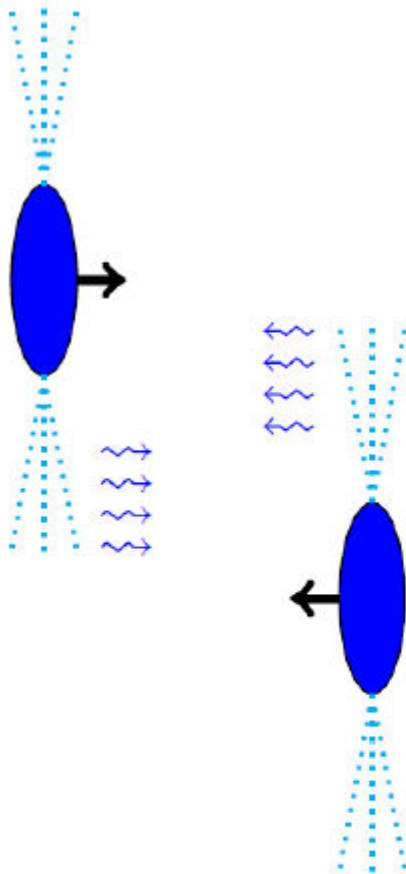
<http://arxiv.org/abs/hep-th/0205086>

*Therefore, we consider that when a charged particle passes near a point, it produces, at that point, a variable electric field. If we decompose this field, via a Fourier transform, into its harmonic components we find that it is equivalent to the electric field at the same point if it were struck by light with an appropriate continuous distribution of frequencies.*



Enrico FERMI

*The electromagnetic field surrounding these protons/ions can be treated as a beam of quasi real photons*



Two ions (or protons) pass by each other with impact parameters  $b > 2R$ . **Hadronic interactions are strongly suppressed**

## pp collisions

### Pros

- harder EPA  $\gamma$  spectrum ( $\omega_{\max} \sim \text{TeV}$ )
- more data available ( $\sim 100 \text{ fb}^{-1}$ )

### Cons

- large pile-up (multiple interactions per bunch crossing)
- problems with triggering on low  $p_T$  objects

## Pb+Pb collisions

### Pros

- AA ( $\gamma\gamma$ ) x-sec  $\propto Z^4$
- gluonic x-sec  $\propto A^2$   $\Rightarrow$  lower QCD bkg.
- low pile-up ( $< 1\%$ )

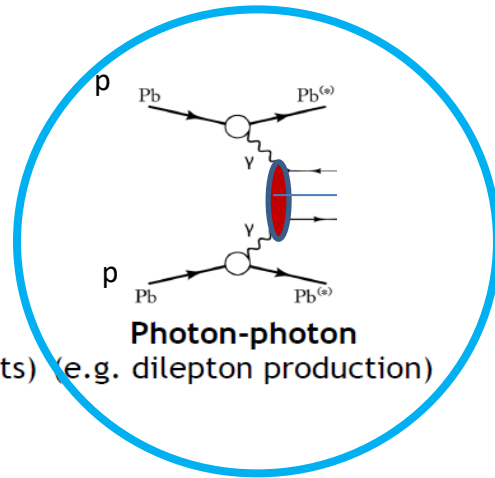
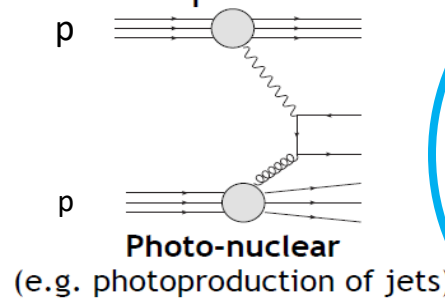
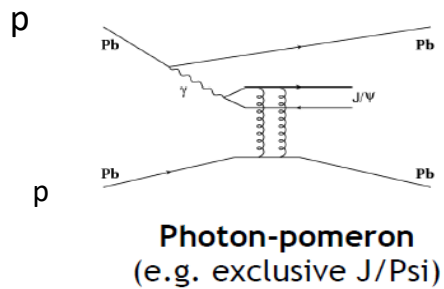
### Cons

- softer EPA  $\gamma$  spectrum ( $\omega_{\max} \sim 0.1 \text{ TeV}$ )
- relatively small data



$A^{1/3}$

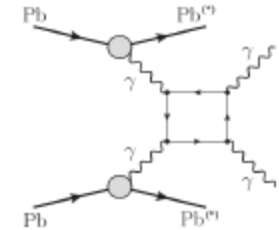
### Various types of interactions possible:



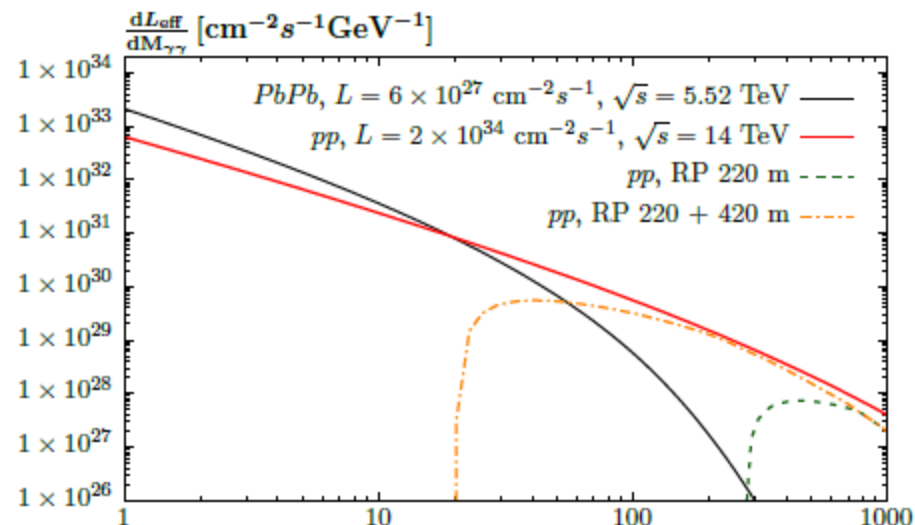
**E** **CEP**   
mphasis



# Heavy Ion Collisions



- **Photon-initiated** CEP equally possible in heavy ion collisions. Indeed in some cases has significant **advantages**:
  - Significant  $\sim Z^4$  enhancement in rate. After accounting for differing luminosities, still  $\sim 2$  relative to  $pp$ , but with no pile-up.
  - QCD-initiated production essentially absent - clear interpretation.
  - Low pile-up - can go to low  $M_{\gamma\gamma}$ .
  - Conversely, steep fall off at high mass -  $pp$  essential here  $\Rightarrow$  **complementary**.


 $M_{\gamma\gamma}$

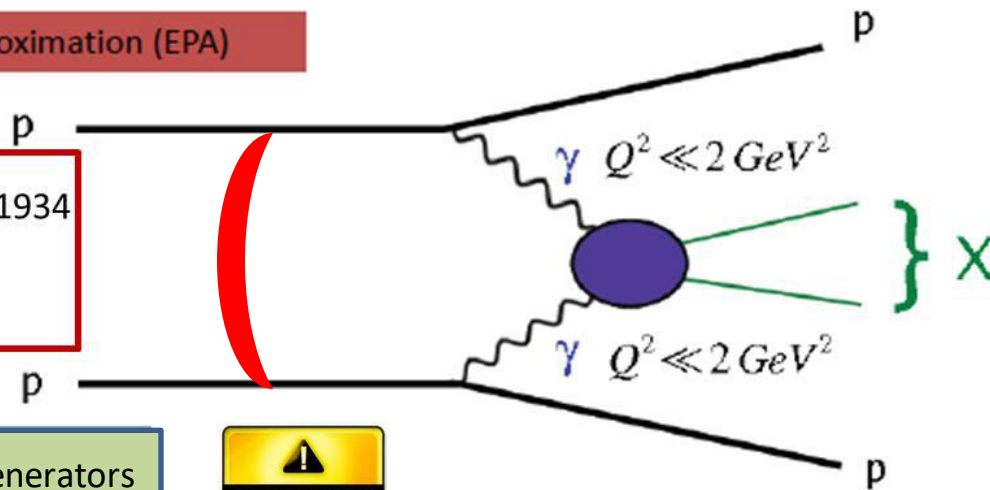
# Modelling Exclusive Photon-Photon collisions

- In exclusive photon-mediated interactions, the colliding protons must both coherently emit a photon, and remain intact after the interaction. How do we model this?
- Answer is well known- the 'equivalent photon approximation' (EPA): cross section described in terms of a flux of quasi-real photons radiated from the proton, and the  $\gamma\gamma \rightarrow X$  subprocess cross section.

Equivalent photon approximation (EPA)

C.F. von Weizsacker, 1934  
E.J. Williams, 1934  
E. Fermi, 1925

introduced to major event generators



# Equivalent photon approximation

- Initial-state  $p \rightarrow p\gamma$  emission can be to v. good approximation factorized from the  $\gamma\gamma \rightarrow X$  process in terms of a flux:

$$n(x_i) = \frac{1}{x_i} \frac{\alpha}{\pi^2} \int \frac{d^2 q_{i\perp}}{q_{i\perp}^2 + x_i^2 m_p^2} \left( \frac{q_{i\perp}^2}{q_{i\perp}^2 + x_i^2 m_p^2} (1 - x_i) F_E(Q_i^2) + \frac{x_i^2}{2} F_M(Q_i^2) \right)$$

- Cross section the given in terms of  $\gamma\gamma$  'luminosity':

$$\frac{d\mathcal{L}_{\gamma\gamma}^{\text{EPA}}}{dM_X^2 dy_X} = \frac{1}{s} n(x_1) n(x_2)$$

THE TWO-PHOTON PARTICLE PRODUCTION MECHANISM.  
PHYSICAL PROBLEMS. APPLICATIONS. EQUIVALENT PHOTON APPROXIMATION

V.M. BUDNEV, I.F. GINZBURG, G.V. MELEDIN and V.G. SERBO  
*USSR Academy of Science, Siberian Division, Institute for Mathematics, Novosibirsk, USSR*

Received 25 April 1974  
Revised version received 5 July 1974

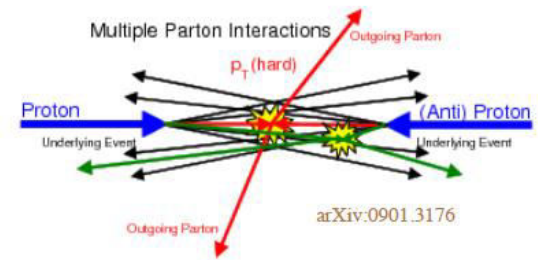
$$\frac{d\sigma^{pp \rightarrow pXp}}{dM_X^2 dy_X} = \langle S_{\text{eik}}^2 \rangle \frac{d\mathcal{L}_{\gamma\gamma}^{\text{EPA}}}{dM_X^2} dy_X \hat{\sigma}(\gamma\gamma \rightarrow X)$$

$$\langle S_{\text{eik}}^2 \rangle = 0.72 \quad : \quad J_P = 0^+$$

$$\langle S_{\text{eik}}^2 \rangle = 0.77 \quad : \quad J_P = 0^-$$

In fact, the situation is more complicated due to the effects caused by the polarization structure of the production amplitude.

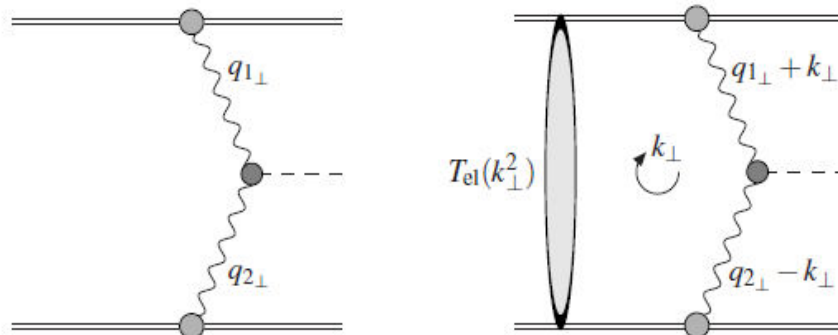




## Soft survival factor

- In any  $pp$  collision event, there will in general be ‘underlying event’ activity, i.e. additional particle production due to  $pp$  interactions secondary to the hard process (a.k.a. ‘multiparticle interactions’, MPI).
- Our  $\gamma\gamma$ -initiated interaction is no different, but we are now requiring final state with no additional particle production ( $X +$  nothing else).

→ Must multiply our cross section by probability of no underlying event activity, known as the soft ‘survival factor’.



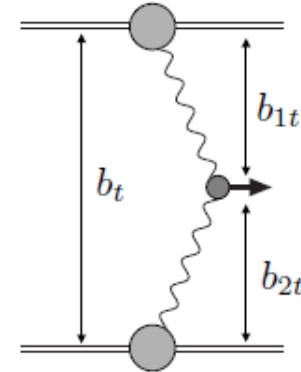


# Soft survival factor

- How do we calculate the survival factor? Work in impact parameter space and apply ‘eikonal’ approach:

$$\langle S^2 \rangle = \frac{\int d^2b_{1t} d^2b_{2t} |T(s, b_{1t}, b_{2t})|^2 \exp(-\Omega(s, b_t))}{\int d^2b_{1t} d^2b_{2t} |T(s, b_{1t}, b_{2t})|^2},$$

$\exp(-\Omega(s, b_t))$ : Poissonian probability of no inelastic scattering at impact parameter  $b_t$ .  
 ↑  
 proton opacity

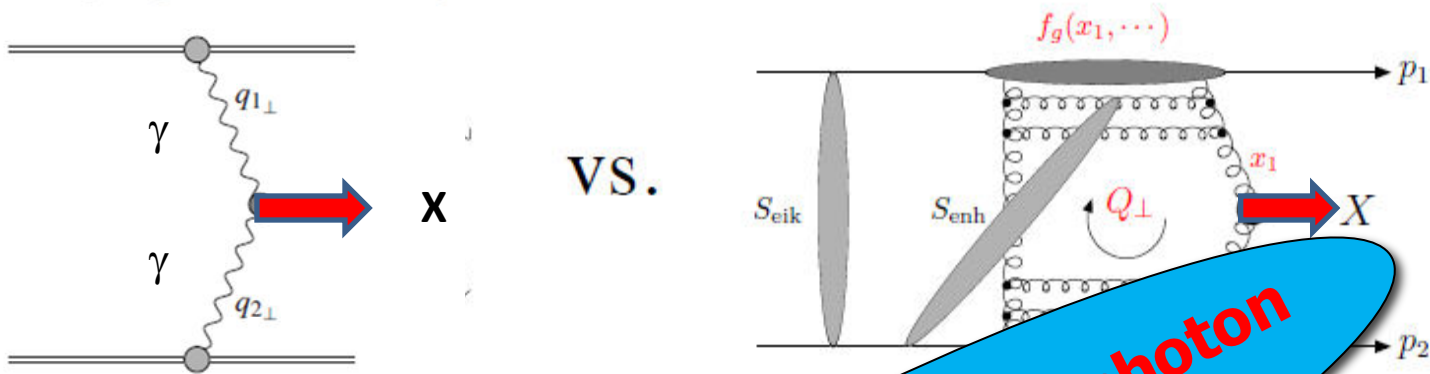


- Underlying event generated by soft QCD. Cannot use pQCD  $\Rightarrow$  take phenomenological approach to this non-pert. observable.

- Have: 
$$\frac{d\sigma^{pp \rightarrow pXp}}{dM_X^2 dy_X} = \langle S^2 \rangle \frac{d\mathcal{L}_{\gamma\gamma}^{\text{EPA}}}{dM_X^2 dy_X} \hat{\sigma}(\gamma\gamma \rightarrow X)$$

KMR, GLM,  
 Ostapchenko,  
 Lund MCs

Well established in the QCD-mediated processes, e.g. CDF-diffractive dijets (2000), CMS/ATLAS- dijets in events with LRG (2015), H1 –diffractive dijet photoproduction (2011).



**The LHC is the high energy photon collider**

- Naively expect strong interactions to be suppressed by  $\alpha_s$ .
- However QCD enhancement is important. Weakness: exclusive event requires no extra particles in final state. Requires introduction of Sudakov form factor:



'Large' Pomeron size in the production of the small size objects.

$$T_g(Q_{\perp}^2, \mu^2) = \exp\left(-\int_{Q_{\perp}^2}^{\mu^2} \frac{dk_{\perp}^2}{k_{\perp}^2} \frac{\alpha_s(k_{\perp}^2)}{2\pi} \int_0^{1-\Delta} \left[ zP_{gg}(z) + \sum_q P_{qg}(z) \right] dz\right)$$

- Increasing  $M_X \Rightarrow$  larger phase space for extra gluon emission stronger suppression in exclusive QCD cross section. Gluons like to radiate!

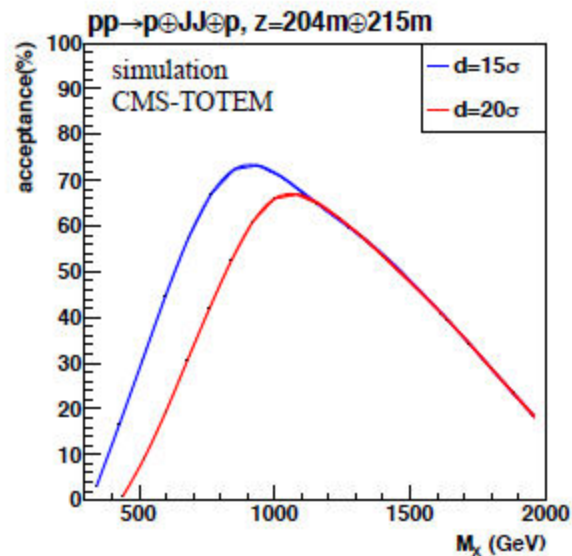
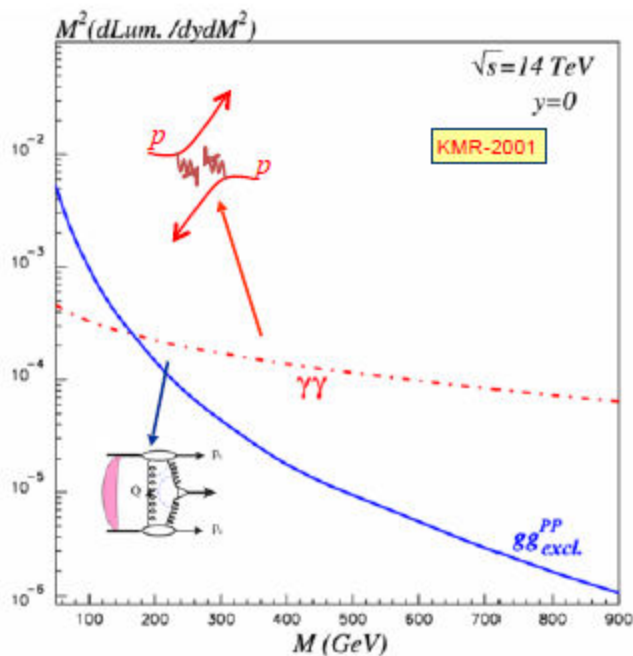
+ absorptive/rescattering effects- survival factor  $S_{\text{soft}}^2$

# $gg$ vs. $\gamma\gamma$

- Situation summarised in ‘effective’ exclusive  $gg$  and  $\gamma\gamma$  luminosities. This Sudakov suppression in QCD cross section leads to enhancement in  $\gamma\gamma$  already\* for  $M_X \gtrsim 200$  GeV - well before CT-PPS/AFP mass acceptance region.



→ Can study  $\gamma\gamma$  collisions at the LHC with unprecedented  $s_{\gamma\gamma}$ .

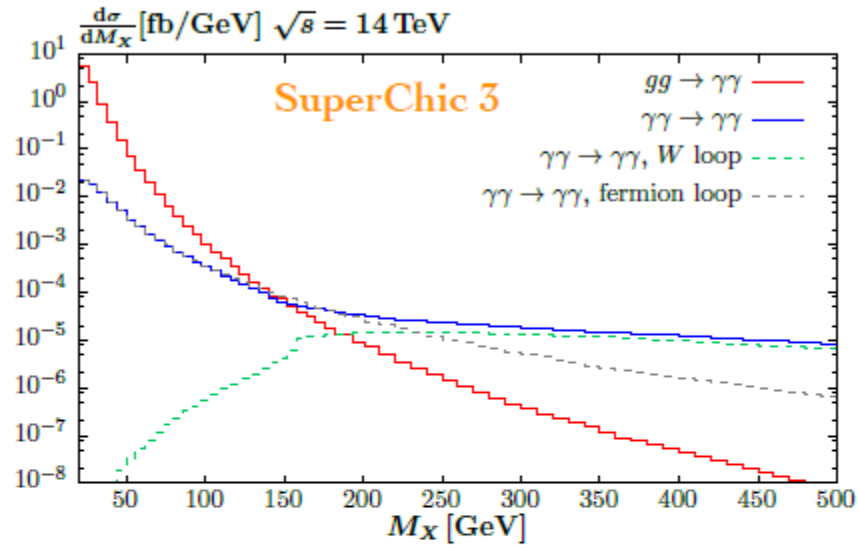


\*Caveat - this is enhancement in initial state only.

16 Of course depends on coupling to produced state.

★ As mass of central system  $M_X$  increases, QCD-initiated production cross section suppressed by no radiation probability  $\Rightarrow$  BG often low<sup>\*</sup>.

- Example of  $\gamma\gamma$  production:



- CEP: unique possibility to observe photon-initiated production of states with EM coupling in clean/well understood environment.
  - However typically considering high mass region (RPs) and relatively low cross sections (EM couplings). **Statistics limited.**
- Increased statistics from HL-LHC running offer **clear advantage** here, in particular in terms of pushing to higher mass.

5 <sup>\*</sup>Precise level depends on particular process.



# CEP(PHOTON)-TOOLS

If a craftsman wants to do good work, he must first sharpen his tools.

**K'ung Fu-tzu**  
**Confucius**

551 - 479 B.C.E.  
Links, Quotes, Bibliography, Sayings, Notes



FPMC; Starlight; LPAIR, SuperChic-2,3 (photon-induced processes)

ExHuME; SuperChic-2,3;FPMC; CEP@PYTHIA; CepGen; Dime; ExDiff;GenEx (CEP-QCD), Lund-Leif&Radek

**SuperChic-3** incorporates **pp**,  
**soft survival** effects.



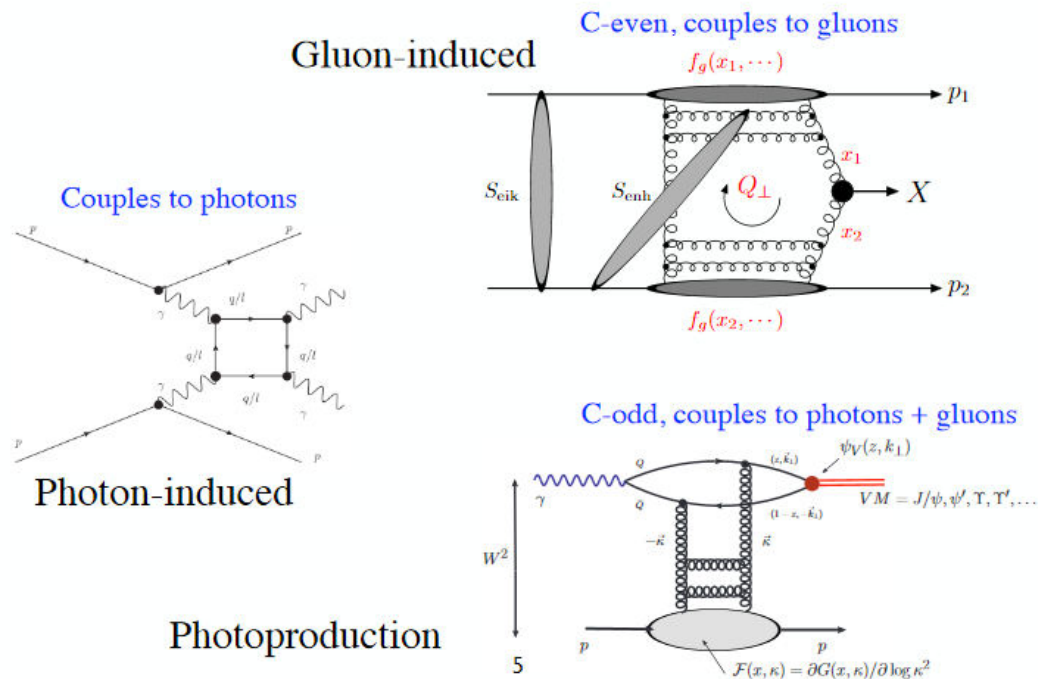
collisions and a complete treatment of



# Photon-photon collisions in Superchic

## Production mechanisms

Exclusive final state can be produced via three different mechanisms, depending on kinematics and quantum numbers of state:



# Tools for future studies - SuperChic MC

- A MC event generator for CEP processes. **Common platform** for:
  - QCD-initiated CEP.
  - Photoproduction.
  - Photon-initiated CEP.
- Previously generated  $pp$  collisions only, but recently updated to include  $pA$  and  $AA$  collisions.

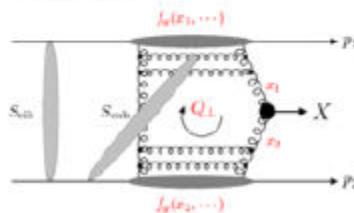
arXiv:1810.06567

superchic is hosted by Hepforge, IPPP Durham

## SuperChic 3 - A Monte Carlo for Central Exclusive Production

- Home
- Code
- References
- Contact

SuperChic is a Fortran based Monte Carlo event generator for central exclusive production in proton and heavy ion collisions. A range of Standard Model final states are implemented, in most cases with spin correlations where relevant, and a fully differential treatment of the soft survival factor is given. Arbitrary user-defined histograms and cuts may be made, as well as unweighted events in the HEPEVT, HEPMC and LHE formats. For further information see the [user manual](#).



15 Oct 2018

## Exclusive LHC physics with heavy ions: SuperChic 3

L.A. Harland-Lang<sup>1</sup>, V.A. Khoze<sup>2,3</sup>, M.G. Ryskin<sup>2</sup>

<sup>1</sup>Rudolf Peierls Centre for Theoretical Physics, University of Oxford, Clarendon Laboratory, Parks Road, Oxford OX1 3PU, United Kingdom.

<sup>2</sup>Institute for Particle Physics Phenomenology, University of Durham, Durham, DH1 3LE

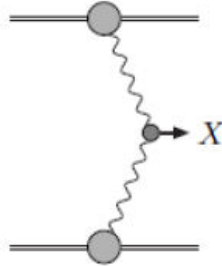
<sup>3</sup>Petersburg Nuclear Physics Institute, NRC Kurchatov Institute, Gatchina, St. Petersburg, 188300, Russia

### Abstract

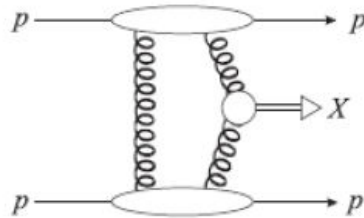
We present results of the updated SuperChic 3 Monte Carlo event generator for central exclusive production. This extends the previous treatment of proton-proton collisions to

<https://superchic.hepforge.org>

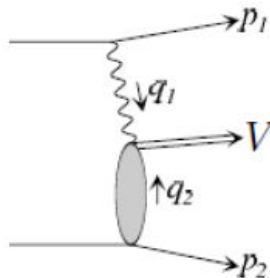
# SuperChic MC - processes generated



★  $W^+W^-$ ,  $l^+l^-$ , LbyL, SM Higgs, ALPs, monopoles, monopolum.



★ SM Higgs, dijets, trijets, light meson pairs, heavy quarkonia (single and double),  $\gamma\gamma$ .

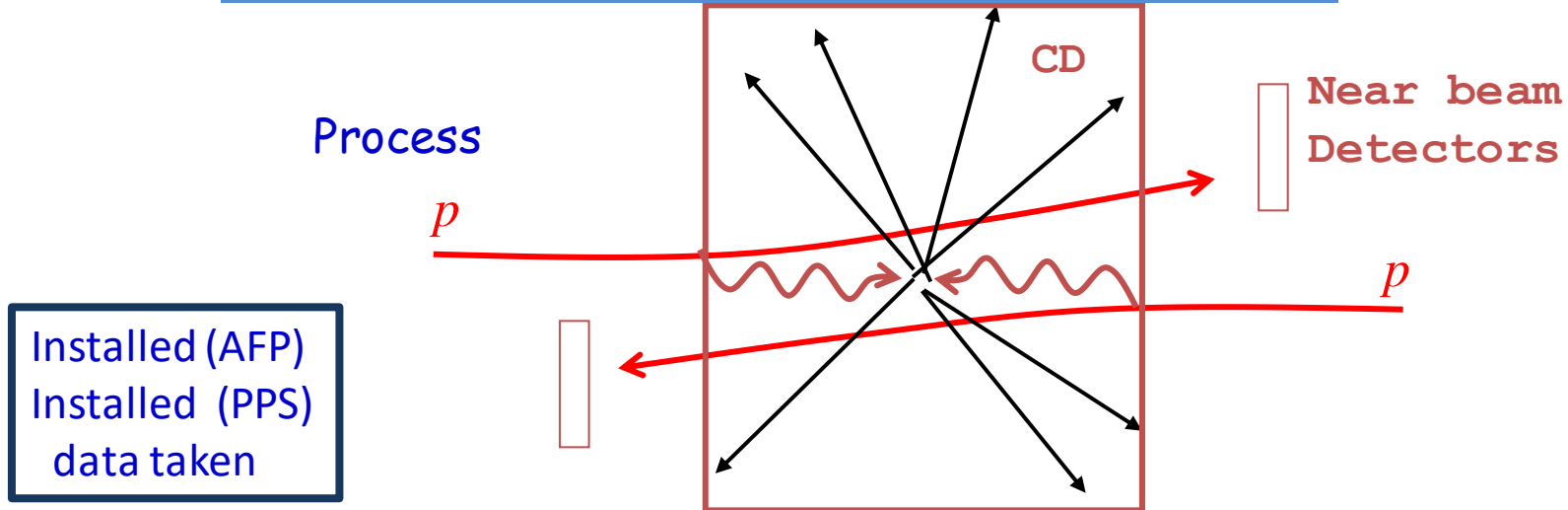


★ Light to heavy vector meson production.

- In all cases in  $pp$ ,  $pA$  and  $AA$  collisions.



# $\gamma\gamma$ collisions- applications



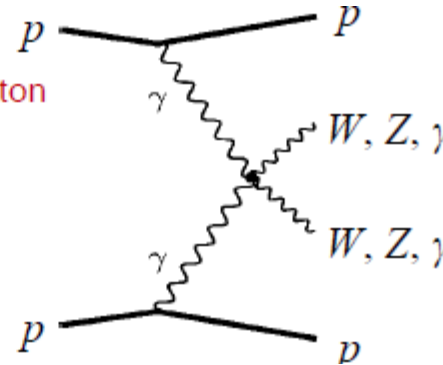
## Extensive Program

- $\gamma\gamma \rightarrow \mu\mu, ee$  QED processes
- $\gamma\gamma \rightarrow$  QCD (jets..)
- $\gamma\gamma \rightarrow WW$  anomalous couplings
- $\gamma\gamma \rightarrow$  squark, top... pairs
- $\gamma\gamma \rightarrow$  Charginos (natural SUSY)
- New BSM objects



# Anomalous Gauge Quartic Couplings

- Low Cross sections: ~few fb
  - AFP has a Missing-Mass resolution (from the proton measurements) of 2-4 %
- Match with invariant central object mass is efficient: ( $Z \rightarrow ee, \gamma\gamma$ )
  - powerful rejection of non-exclusive backgrounds
- Much interest in this from theory side
  - e.g. “LHC Forward Physics” CERN-PH-LPCC-2015-001)

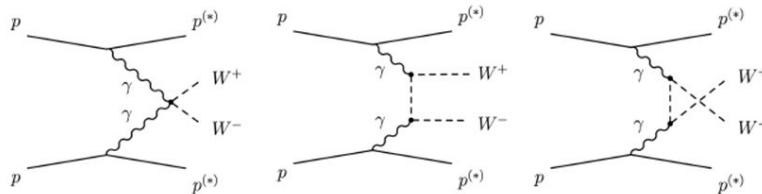


“Probing anomalous quartic gauge couplings using proton tagging at the Large Hadron Collider”, M. Saimpert, E. Chapon, S. Fichet, G. von Gersdorff, O. Kepka, B. Lenzi, C. Royon; 23/05/2014

hep-ph/09082020, Louvain Group

• Exclusive  $W^+W^-$  production: no contribution from  $q\bar{q} \rightarrow W^+W^- \Rightarrow$  sensitive to  $\gamma\gamma \rightarrow W^+W^-$  process alone.

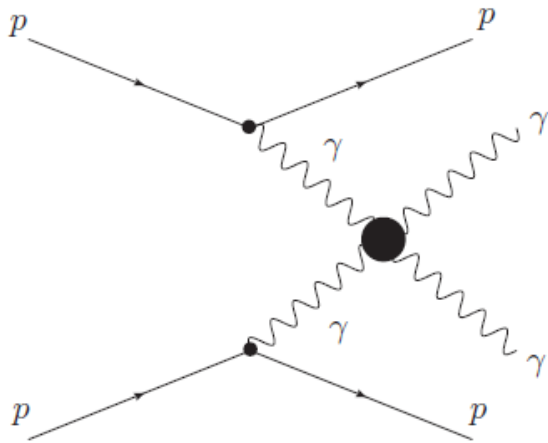
→ Directly sensitive to any deviations from the SM gauge couplings. Predicted in various BSM scenarios. Composite Higgs, warped extra dimensions....



• Limits have been set at LEP, and in inclusive final-states at the Tevatron and LHC. How does the exclusive case compare?

Currently very encouraging ATLAS & CMS data

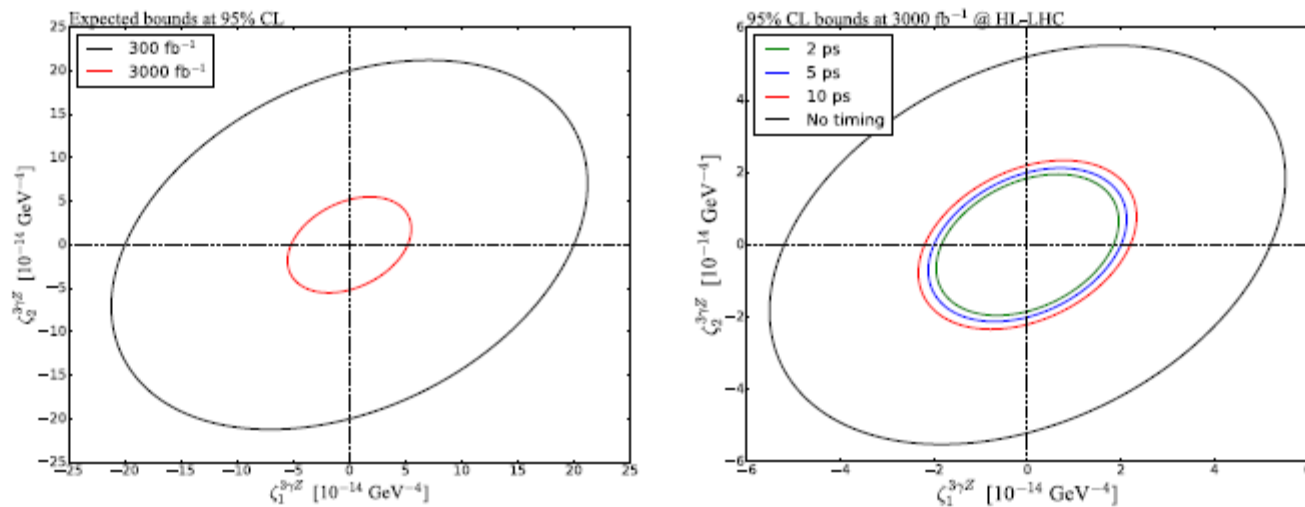
# Search for extra dimensions in the universe using $\gamma$ induced processes



- Additional channels:  $WW$ ,  $ZZ$ ,  $\gamma Z$ , dilepton production ( $\gamma\gamma$  described in more detail as an example)

- Search for production of two photons and two intact protons in the final state:  $pp \rightarrow p\gamma\gamma p$
- Number of events predicted to be increased by extra-dimensions, composite Higgs models, dark matter particles
- Anomalous couplings can appear via loops of new particles coupling to photons or via resonances decaying into two photons
- In addition, SM production of  $WW$  or dileptons will be still statistically limited after run III

# Anomalous couplings@ HL-LHC



Left: Comparison between  $300 \text{ fb}^{-1}$  and  $3000 \text{ fb}^{-1}$ . Right: (Zoomed-in; change of scale in X-Y axis) comparison between the use of timing  $\delta t = 2, 5, 10 \text{ ps}$ .  $Z\gamma \rightarrow \text{hadrons} + \gamma$  benefits the most from the use of timing.

- Expected **improvements** from HL-LHC impressive  $\sim$  an **order of magnitude** ( $\sim 5$  orders of magnitude better than current best inclusive limits).



# Recent Developments

• prospects for compressed SUSY searches in photon-initiated CEP.

- Signal for this study:

$$\gamma\gamma \rightarrow \tilde{l}^+\tilde{l}^- \rightarrow l^+l^-\tilde{\chi}_0\tilde{\chi}_0$$

implemented in SuperChic, as well as more challenging case:

$$\gamma\gamma \rightarrow \tilde{\chi}^+\tilde{\chi}^- \quad (\text{hadron/leptonic decays})$$

- In addition, contribution from proton dissociation included in effective way, interfaced to SuperChic.
- Future work: include (more) complete treatment of dissociation. Stay tuned!

IPPP/18/103

LHC Searches for Dark Matter in Compressed Mass Scenarios:  
Challenges in the Forward Proton Mode

L.A. HARLAND-LANG<sup>1\*</sup>, V.A. KHOZE<sup>2,3†</sup>, M.G. RYSKIN<sup>2‡</sup>  
AND M. TASHINSKY<sup>4§</sup>

<sup>1</sup>Rudolf Peierls Centre, Beecroft Building, Parks Road, Oxford, OX1 3PU, UK

<sup>2</sup>IPPP, Department of Physics, University of Durham, Durham, DH1 3LE, UK

<sup>3</sup>Petersburg Nuclear Physics Institute, NRC "Kurchatov Institute", Gatchina, St. Petersburg, 188300, Russia

<sup>4</sup>Institute of Physics, Czech Academy of Sciences, CS-18221 Prague 8, Czech Republic



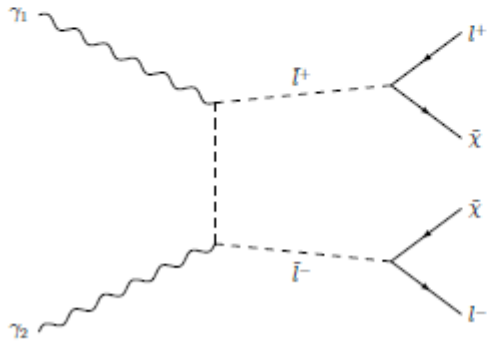
$$pp \rightarrow p + \gamma\gamma + p,$$

$$\gamma\gamma \rightarrow X^+X^-,$$

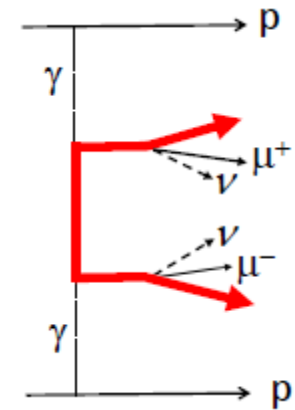
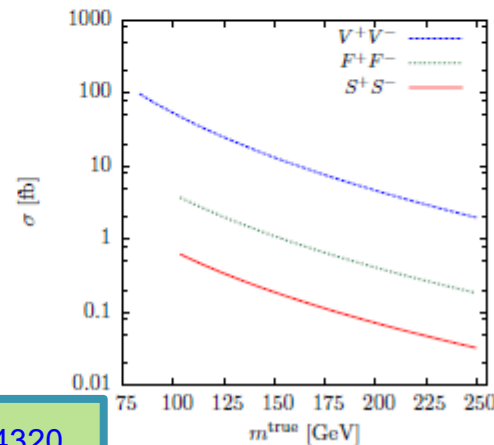
## Diphoton X-Pair Production

where  $X = W$ -boson, lepton, slepton, chargino...

- If particle decays semi-invisibly, then additional information from tagged proton momenta can be used to measure masses and discriminate BG.



[HKSS, arXiv:1110.4320](https://arxiv.org/abs/1110.4320)



- Consider exclusive production of chargino pair  $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ , decaying via

$$\tilde{\chi}_1^+ (\tilde{\chi}_1^-) \rightarrow l^+ (l^-) + \nu (\bar{\nu}) + \tilde{\chi}_1^0,$$

electroweakinos

where the  $\tilde{\chi}_1^0$  is an LSP neutralino.

- For cases that  $\Delta M = M(\tilde{\chi}_1^0) - M(\tilde{\chi}_1^\pm)$  is relatively small, can be difficult to observe inclusively. (compressed mass BSM scenarios)

• to report current status of our ongoing long-term studies on prospects of searches at the LHC for **ELECTROWEAKINO** pair production via photon fusion with forward proton detectors (AFP, CT-PPS)

• exemplified within the framework of the compressed mass **MSSM**

First discussed: **KMR**, J.Phys. G44 (2017) no.5, 055002

**HKRT JHEP 1904 (2019) 010**

Some recent studies:

L.Beresford, Jesse Liu, ArXiv:1811.0645

S.I. Godunov et al, ArXiv: 1906.08568



**SUSY** — solution to various shortcomings of SM (as an example only)

If (it looks like) squarks and gluinos are too heavy, sleptons, charginos, neutralinos- the main target.



**(null search result so far)**

**MSSM** : charginos  $\tilde{\chi}_{1,2}^{\pm}$  four neutralinos  $\tilde{\chi}_{1,2,3,4}^0$

$\tilde{\chi}_1^0$ , natural candidate for cold Dark Matter –**LSP**

[arXiv:1710.02406](https://arxiv.org/abs/1710.02406)

(and quite a few other papers)

natural SUSY:  
existence of light nearly  
mass-degenerate Higgsinos/charged  
Mass  $\sim 100-200 \text{ GeV}$   
mass splittings

Most challenge  
scenario  
between

Motivated by naturalness,  
cosmological observations and (g-2)  
phenomenology.

## Naturalness and light Higgsinos: why ILC is the right machine for SUSY discovery

**Howard Baer**  
University of Oklahoma, Norman, OK 73019, USA  
E-mail: baer@ou.edu

**Mikael Berggren, Suvi-Leena Lehtinen**  
DESY, Hamburg, Germany  
E-mail: berggren@desy.de, lehtinen@desy.de

**Yuta Koide**  
University of Tokyo, Tokyo, Japan E-mail: koide@hep.t.u-tokyo.ac.jp

**What about the LHC during MY lifetime ?**

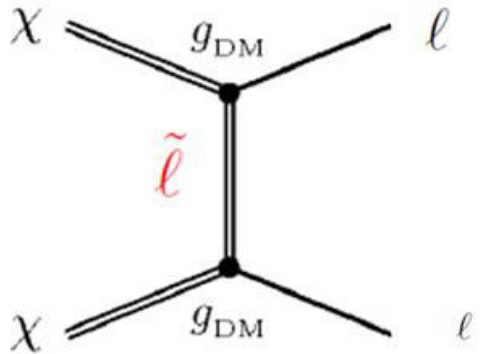
Supersymmetry, a theoretically and experimentally well-motivated theory, predicts the existence of four light, nearly mass-degenerate Higgsinos with mass  $\sim 100 - 200 \text{ GeV}$  (not too far above  $m_Z$ ). The small mass splittings amongst the higgsinos, typically 4-20 GeV, results in very little visible energy arising from decays of the heavier higgsinos. Given that other SUSY particles are considerably heavy, this makes detection challenging at hadron colliders. On the other hand, the clean environment of an electron-positron collider with  $\sqrt{s} > 2m_{\text{higgsino}}$  would enable a decisive search of these required higgsinos, and thus either the discovery or exclusion of natural SUSY. We present a detailed simulation study of precision measurements of higgsino masses and production cross sections at  $\sqrt{s} = 500 \text{ GeV}$  of the proposed International Linear Collider currently under consideration for construction in Japan.

$$e^+e^- \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^- \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0 q\bar{q}' e\nu_e (\mu\nu_\mu).$$

# Co-annihilation

(1702.00750, model-1a)

## Dark matter annihilation

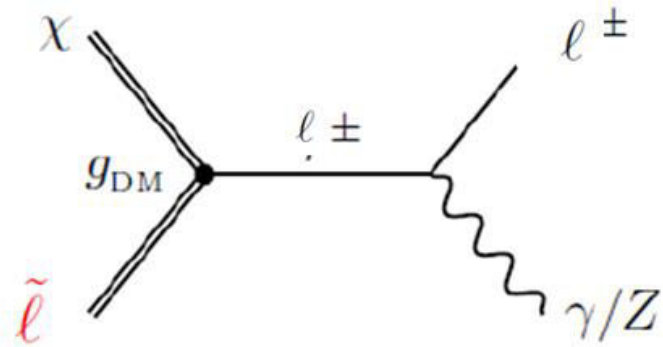


to bring DM abundance down to the observed value

Initially DM in thermal equilibrium with SM, later it freezes out

- Overproduces dark matter (Unless large couplings)
- We need a mechanism to reduce the DM relic density

### Co-annihilation:



Freeze-out temperature  $T_F \sim m_{DM}/25$

Boltzmann factor  $\exp\left(-\frac{\Delta M}{T}\right)$   $\longrightarrow$

$$\Delta M \lesssim m_{DM}/25$$

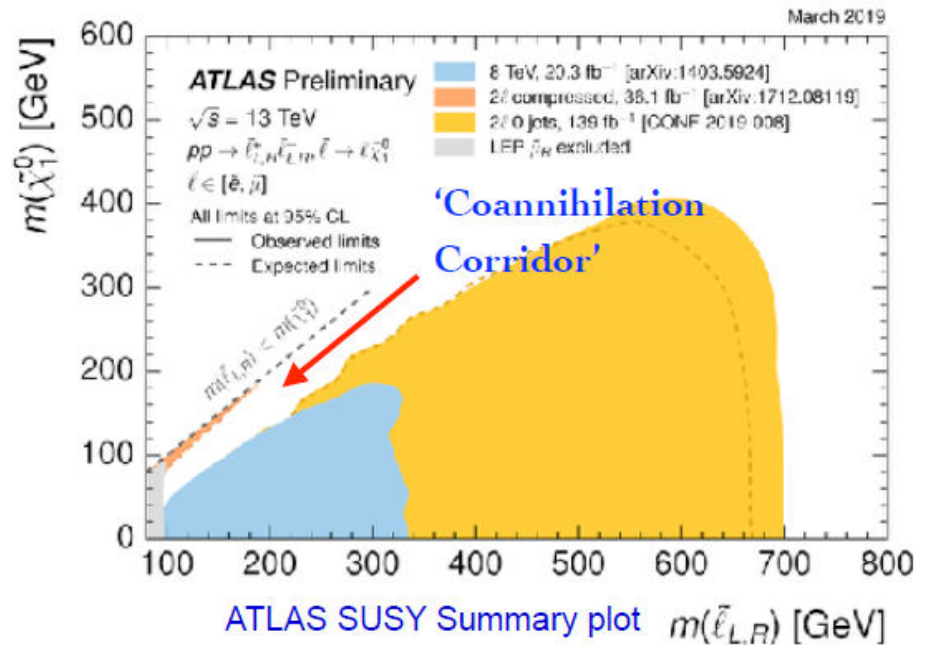
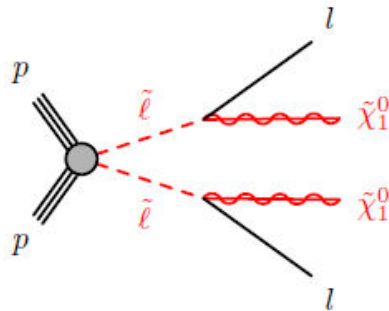
We need **mass splitting of 4% of  $m_{DM}$**



# SUSY at the LHC

- **Pre-LHC**: EW-scale SUSY theoretically well motivated BSM scenario: hierarchy problem, coupling unification, natural DM candidate...
- **Post-LHC** folklore: no EW-scale SUSY to be seen! ↗ Lightest SUSY particle = 'LSP'
- Only half true: most significant limits based on 'classic' large missing  $E_{\perp}$  signal, requiring **largish** SUSY particle **mass splittings**.

• For e.g. small slepton-neutralino mass differences, LEP constraints still dominant!

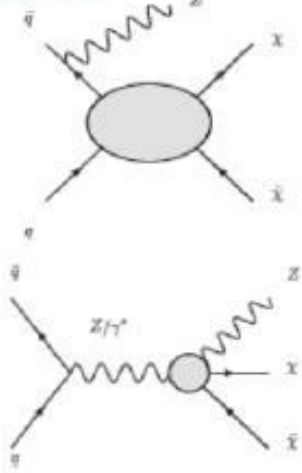




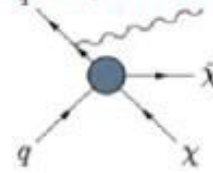
# Mono-Mania (at the LHC)

DM Searches @ LHC: O. Buchmüller

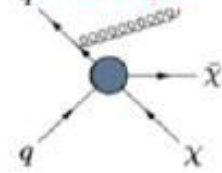
Mono-Z



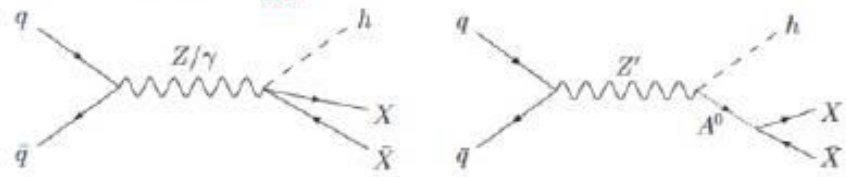
Mono-photon



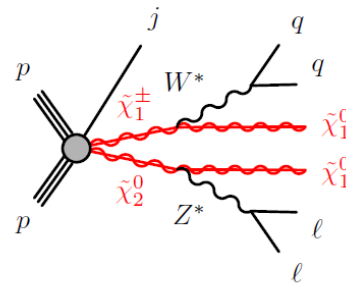
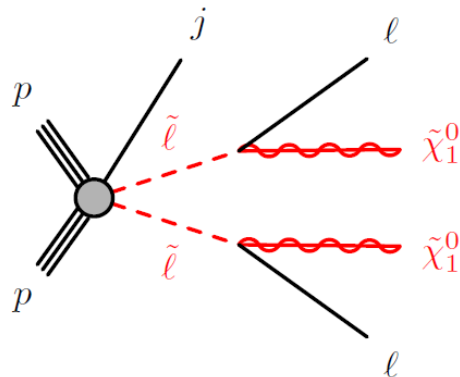
Mono-jet



Mono-Higgs



## Searches for Electroweakinos at the LHC



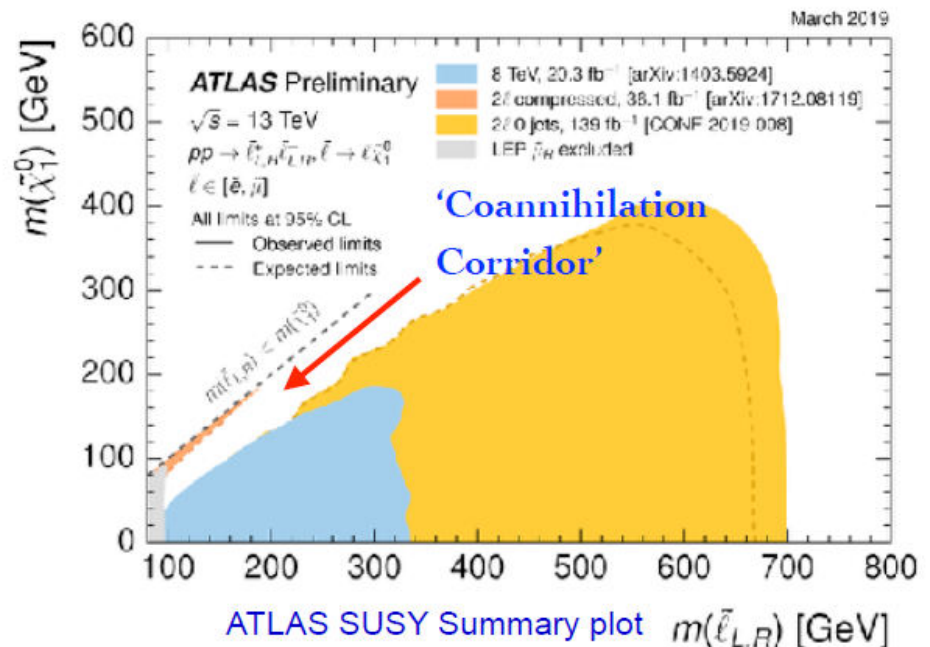
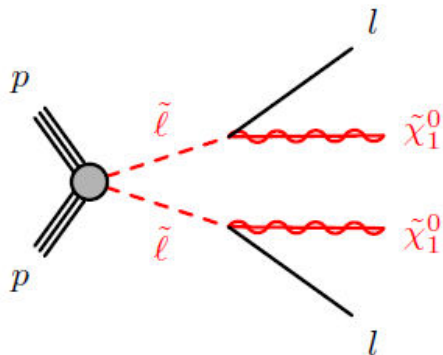
Model dependence

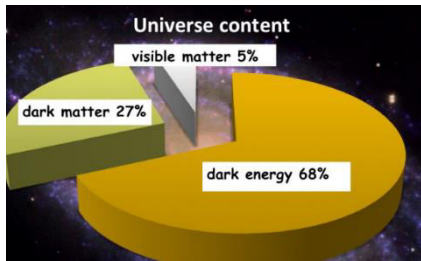


# SUSY at the LHC

- Such '**compressed SUSY**' scenarios not just dreamt up to avoid limits.
- **Theoretically motivated** by naturalness,  $(g - 2)$  phenomenology, and cosmological considerations (coannihilation  $\Rightarrow$  correct DM abundance).
- Inclusive cross sections not small (up to  $\sim 100$ s of fb)  $\Rightarrow$  huge number of events may be produced at LHC, but **lost in BG**.

- How can **Photon-initiated** production help?





(SUSY, DM -null LHC search results so far)

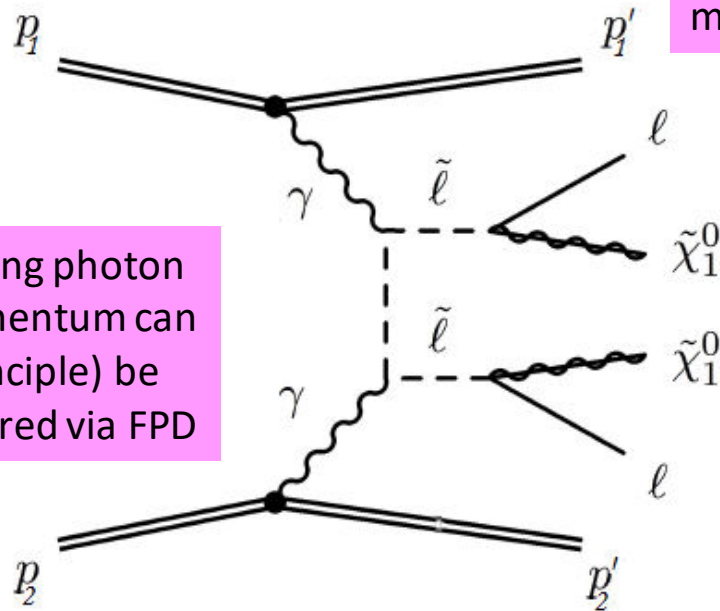


## K'ung Fu-tzu Confucius

551 - 479 B.C.E.  
Links, Quotes, Bibliography, Sayings, Notes

**Better to  
light a  
candle  
than  
curse the  
darkness.**

# Advantage of exclusivity & compressed mass



Incoming photon 4-momentum can (in principle) be measured via FPD

Outgoing proton 4-momentum measured in FPD

Lepton 4-momentum measured in Central detector

4-momentum of system of 2 DM particles could be constrained from photon & lepton 4-momenta

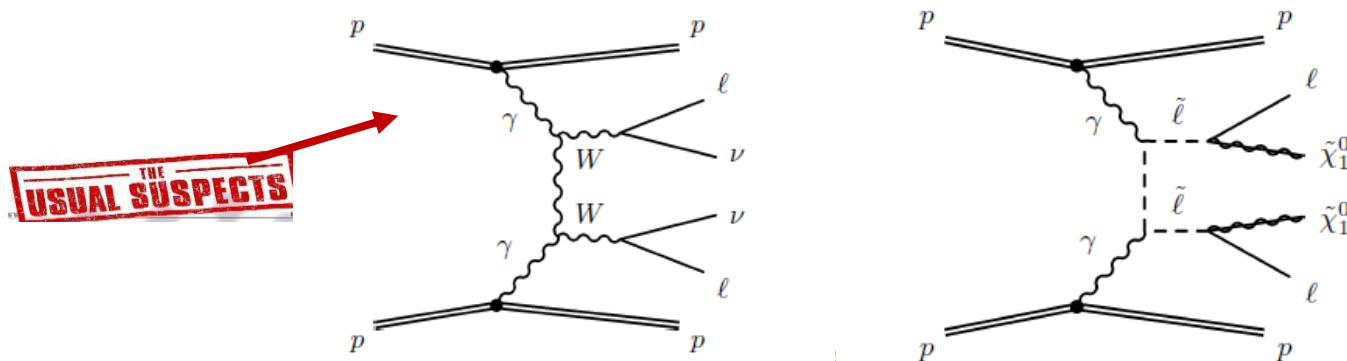
$$\xi_i = 1 - \frac{E_{p'_i}}{E_{p_i}}, i=1,2$$

measured precisely in FPD

FPD measures precisely mass  $M_{\tilde{l}} - M_{\tilde{\chi}_1^0}$  system. If mass splitting low  $\rightarrow$  FPD can give quite a precise hint about  $2m_{DM}$

# CEP and SUSY

- Possibility of  $\sim 100$  GeV mass slepton/chargino production at LHC begin swamped by huge inclusive BGs.
- **Exclusive photon-initiated** production a natural mechanism:
  - ★ **Well understood**, model-independent signal cross section.
  - ★ Irreducible WW BG can be controlled. No need for large missing  $E_{\perp}$ .
  - ★ **Proton tagging**: can reconstruct mass of central system from protons alone ('missing mass'). Crucial handle for BGs.
- But, how feasible is this in **high pile-up** environment?





# Classes of Background

$$(M_{\tilde{l}}, M_{\tilde{\chi}_1^0}) = (120, 110) \text{ GeV}$$



$$(M_{\tilde{l}}, M_{\tilde{\chi}_1^0}) = (300, 280) \text{ GeV}$$

- Take **slepton pairs** for concreteness. Signal selection:

★ Low  $\Delta M_{\tilde{l}\tilde{\chi}_0} \Rightarrow$  two relatively low  $p_{\perp}$  leptons, with low  $m_{ll}$ , in central detector.

★ Two proton hits in AFP/PPS ( $\sim 220\text{m}$ ) acceptance.

$$5 < p_{T,l_1,l_2} < 40 \text{ GeV}$$

$$2 < m_{ll_2} < 40 \text{ GeV}$$

$$|\eta_{l_1,l_2}| < 2.5 \text{ (4.0)}$$

$$0.02 < \xi_{1,2} < 0.15$$

- What are **backgrounds**?

★ Irreducible CEP of **W pairs**.

★ Reducible **semi-exclusive** production ( $l^+l^- \dots$ ) with proton from dissociation system giving hit in forward proton detector (FPD).

★ Reducible **pile-up** background: coincidence of non-diffractive event with hits from independent diffractive events.

- Realistic analysis must consider all three.

$$M_{\tilde{l}} = 120-300 \text{ GeV}, \quad \Delta M = M_{\tilde{l}} - M_{\tilde{\chi}_1^0} = 10-20 \text{ GeV}$$

## Major backgrounds

- $\gamma\gamma \rightarrow W^+W^- \rightarrow l^+\nu + l^-\bar{\nu}$
- **Low mass  $\gamma\gamma \rightarrow l^+l^-$  production**  
Semi-exclusive process with proton from (SD,DD) dissociation detected in the FPD.
- Semi-exclusive QCD-initiated BGs due to low-pt (mainly c-quark) jets, with SD and DD followed by proton hits in the FPD.
- Coincidence of inelastic lepton pair production with two independent SD/DD events from the PU interactions that mimics the signal.  
(danger for other New Physics searches with  $\sigma \leq 1 \text{ fb}$ )
- $\tau\tau$ , dimeson, vector resonances etc...



# Results

$$(M_{\tilde{l}}, M_{\tilde{\chi}_1^0}) = (120, 110) \text{ GeV} \longrightarrow (M_{\tilde{l}}, M_{\tilde{\chi}_1^0}) = (300, 280) \text{ GeV}$$

$$|\eta| < 2.5$$

Event yields / $\mathcal{L} = 300 \text{ fb}^{-1}$	$\langle \mu \rangle_{PU}$		
	0	10	50
Excl. sleptons	0.6—3.9	0.5—3.3	0.3—1.9
Excl. $l^+l^-$	1.4	1.2	0.7
Excl. $K^+K^-$	$\sim 0$	$\sim 0$	$\sim 0$
Excl. $W^+W^-$	0.7	0.6	0.3
Excl. $c\bar{c}$	$\sim 0$	$\sim 0$	$\sim 0$
Excl. $gg$	$\sim 0$	$\sim 0$	$\sim 0$
Incl. ND jets	$\sim 0(\sim 0)$	0.1(0.1)	1.8(2.4)

$$|\eta| < 4.0$$

Event yields / $\mathcal{L} = 300 \text{ fb}^{-1}$	$\langle \mu \rangle_{PU}$		
	0	10	50
Excl. sleptons	0.7—4.3	0.6—3.6	0.3—2.1
Excl. $l^+l^-$	1.1	0.9	0.5
Excl. $K^+K^-$	$\sim 0$	$\sim 0$	$\sim 0$
Excl. $W^+W^-$	0.6	0.5	0.3
Excl. $c\bar{c}$	$\sim 0$	$\sim 0$	$\sim 0$
Excl. $gg$	$\sim 0$	$\sim 0$	$\sim 0$
Incl. ND jets	$\sim 0(\sim 0)$	0.03(0.05)	0.6(0.7)

- Final signal yield- **handful** of events.
- Irreducible WW BG under control. Most significant BG from **pile-up**, with dilepton production + **dissociation** a close runner-up.

# Future Improvements

- What improvements might we expect in the future?
  - ★ Cut on distance between **secondary** and **primary** vertex: reduce BG from decays of heavier particles (dominant part of inclusive BG).
  - ★ Improved **ToF resolution** in **FPDs** (ToF rejection increases linearly with decreasing resolution).
  - ★ Radiation hard **ZDCs** with timing to suppress proton dissociation BG.
  - ★ Add **timing info** to **central detector** - considered for HL-LHC upgrades at forward rapidity, and envisaged by CMS centrally.

# CONCLUSIONS

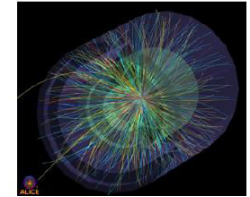


- Have discussed possibility to search for compressed SUSY scenarios via exclusive photon-initiated production at LHC.
- Highly attractive proposal, as very hard to probe via inclusive channels.
- However, important to consider all sources of backgrounds in pile-up heavy nominal LHC environment.
- Possible to bring the backgrounds under control, at the price of a limited significance  $S \sim 2$ ,  $B \sim 2$  events for  $300 \text{ fb}^{-1}$ .
- But not the end of the story- only a first study, and many potential avenues for improvement to explore.
- Ongoing work: more complete treatment of proton dissociation in SuperChic. Stay tuned!





# Processes of interest



## Heavy Ion Collisions

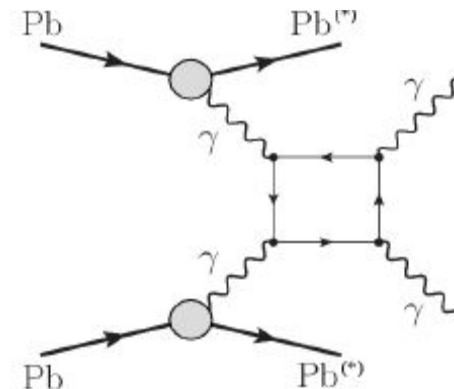
- ★ Light-by-light SM signal, but also e.g. Born-Infeld extensions.
  - ★ Axion-like particle production.
  - ★ Magnetic monopoles.
  - ★ Other possibilities? Gravitons, radions, unparticles, SUSY?
- A principle drawback of heavy ions for these studies is the low luminosity  $\Rightarrow$  benefit from increased datasets can be significant.

SM

- [D. d'Enterria et al. PRL 111 (2013) 0804
- [A. Szczurek et al. PRC 93 (2016) 4, 044]

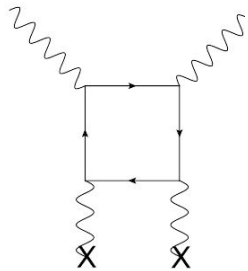
Production mode	BSM particle/interaction	Remarks
Ultraperipheral	Axion-like particles	$\gamma\gamma \rightarrow a$ , $m_a \approx 0.5-100$ GeV
	Radion	$\gamma\gamma \rightarrow \phi$ , $m_\phi \approx 0.5-100$ GeV
	Born-Infeld QED	via $\gamma\gamma \rightarrow \gamma\gamma$ anomalies
	Non-commutative interactions	via $\gamma\gamma \rightarrow \gamma\gamma$ anomalies
Schwinger process	Magnetic monopole	Only viable in HI collisions
Hard scattering	Dark photon	$m_{A'} \lesssim 1$ GeV, advanced particle ID
	Long-lived particles (heavy $\nu$ )	$m_{LLP} \lesssim 10$ GeV, improved vertexing
Thermal QCD	Sexaquarks	DM candidate

Table 1: Examples of new-physics particles and interactions accessible in searches with HI collisions at the LHC, listed by production mechanism. Indicative competitive mass ranges and/or the associated measurement advantages compared to the  $pp$  running mode are given.



# Long and chequered history

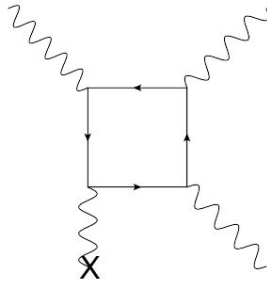
## (nonlinear effects of QED)



Delbrück **1933**

Scattering of gamma-rays by a Coulomb field of heavy nuclei.  
First observed-1953 for 1.33 MeV on lead nuclei.  
Most accurate high-energy results- Novosibirsk,VEPP-4M 1998.

Delbrück scattering



First claims of observation- DESY, PRD 8(1973) 3813.  
Criticised by V.A.Khoze et al, ZhETF Pis.Red.19 (1974) 47.  
First observation- Novosibirsk, VEPP-4M 2002.

Photon splitting in atomic Coulomb field

first direct observation of  
 $\gamma\gamma \rightarrow \gamma\gamma$  scattering



(ArXiv:1702.01625)

(CMS-Arxiv:1902.0139)



# New processes

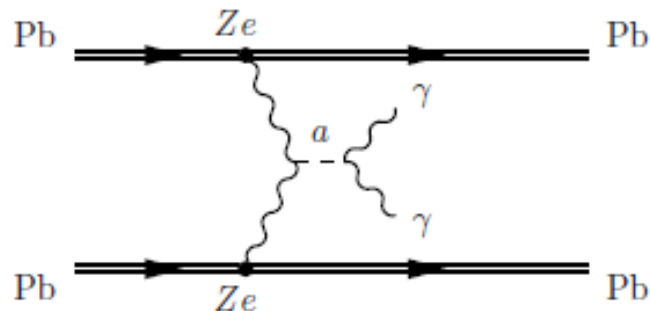
$$\frac{d\sigma^{pp \rightarrow pXp}}{dM_X^2 dy_X} \sim \frac{d\mathcal{L}_{\gamma\gamma}^{\text{EPA}}}{dM_X^2 dy_X} \hat{\sigma}(\gamma\gamma \rightarrow X)$$

- SuperChic has the capability to simulate any arbitrary process given the  $\gamma\gamma \rightarrow X$  amplitudes.

→ Simple to implement new processes within framework.  
Suggestions/collaboration welcome!

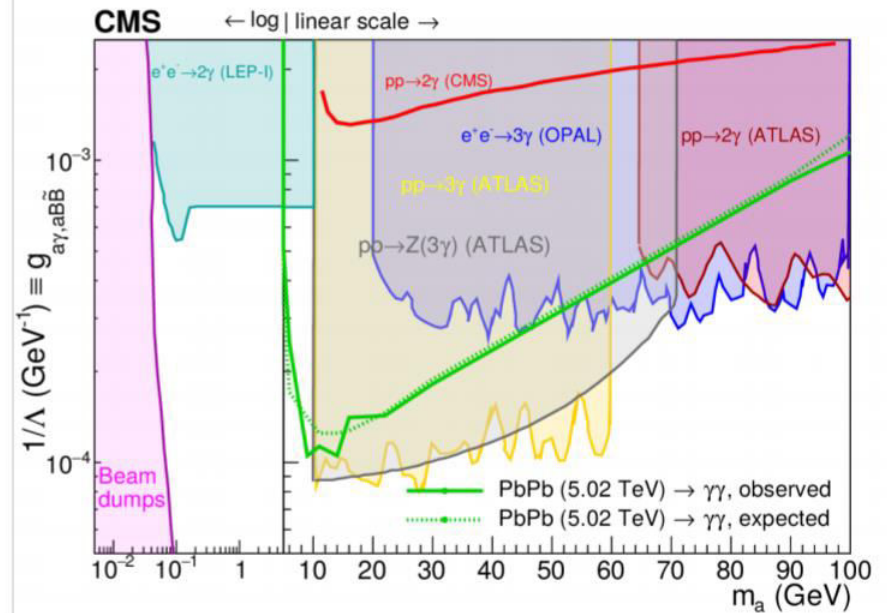
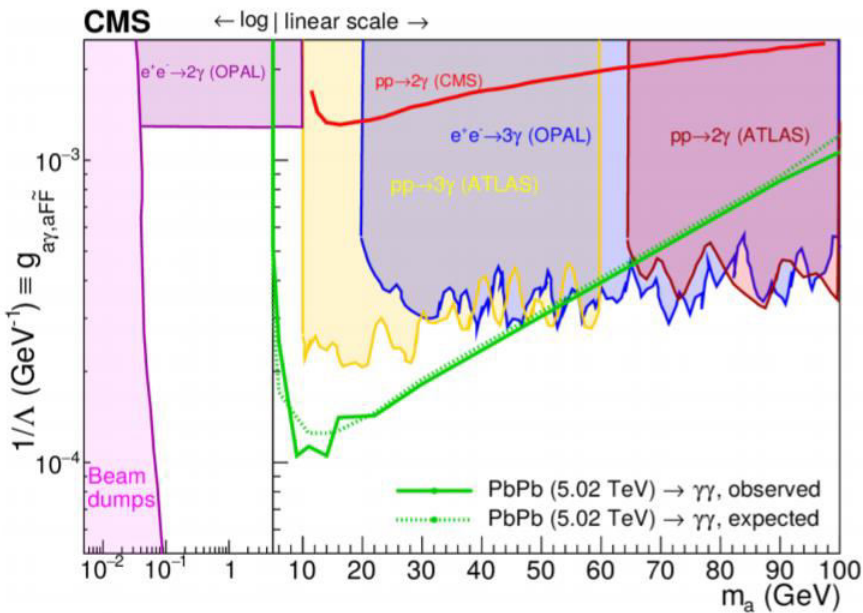
axion-like particles

ALPs



- Limits in cross-section  $\rightarrow$  limits in  $g_{a\gamma}$  vs.  $m_a$  plane
- ALPS coupling only to photons (left plot)
  - Best exclusion limits so far over the  $m_a = 5-50$  GeV
- ALPS coupling to photons or hypercharge (right plot)
  - New constraints in the  $m_a = 5-10$  GeV.

arXiv: 1810.04602



PAPER

## Can invisible objects be 'seen' via forward proton detectors at the LHC?

V A Khoze<sup>1,2</sup>, A D Martin<sup>1,3</sup> and M G Ryskin<sup>1,2</sup>

Published 7 April 2017 • © 2017 IOP Publishing Ltd

[Journal of Physics G: Nuclear and Particle Physics](#), Volume 44, [Number 5](#)



$$pp \rightarrow p + \text{invisible} + p,$$

An attractive idea, but huge backgrounds caused by soft proton dissociation, photon bremsstrahlung and PU (at high lumi)



$$p \rightarrow p + \gamma, N^* \rightarrow p + \gamma \text{ and } N^* \rightarrow p + \pi^0.$$

$$p \rightarrow p\pi^+\pi^-$$

Measurements at low lumi ( $\mu \sim 1$ ) with 'veto' detectors (like ZDC and FSC/ADA/ADC)

LHCb, ALICE



# New Result - The Odderon

- Recent result from CMS-TOTEM: evidence for QCD Odderon.



## Experiments at the Large Hadron Collider uncover possible evidence of elusive 'odderon' that physicists have sought after for decades

- Physicists have been looking for subatomic quasiparticle, 'odderon' since 1970s
- It involves collisions in which an odd number of gluons are exchanged
- While it hasn't been seen in earlier experiments, technology is now more precise

By CHEYENNE MACDONALD FOR DAILYMAL.COM

PUBLISHED: 00:46, 2 February 2018 | UPDATED: 00:46, 2 February 2018

## New Scientist

HOME NEWS TECHNOLOGY SPACE PHYSICS HEALTH ENVIRONMENT TRAVEL LIFE SCIENCE EVENTS AND

SUBSCRIBE LOGIN

THIS WEEK 7 February 2018

## LHC finds hints of strange quasiparticle called the odderon

WE MAY have glimpsed an odd member of the particle family tree.

EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH



TOTEM-2017-002  
16 December 2017



CERN-EP-2017-335  
19 December 2017

First determination of the  $\rho$  parameter at  $\sqrt{s} = 13$  TeV – probing the existence of a colourless three-gluon bound state

The TOTEM Collaboration

## International Business Times

IBT UK World Business Finance Technology Science Sport Entertainment Opinion Video Pictures Vouchers Newsletter Sign up

### Science

## What is Odderon? LHC uncovers possible subatomic quasiparticle which existed only in theory

- Scientists have been searching for the elusive odderon since the 1970's.



By India Ashok

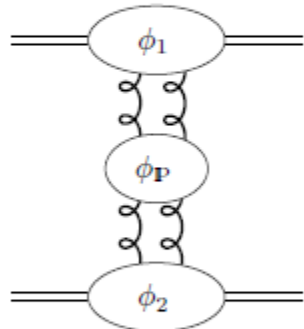
February 2, 2018 11:54 GMT



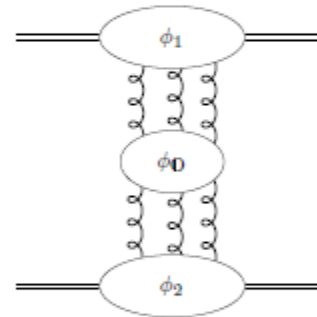
# What is the Odderon?

- Most easily pictured perturbatively. For high energy scattering ( $s/|t| \gg 1$ ) with  $\alpha_s \ll 1$ , dominant elastic amplitude given by resummation of  $\sim \alpha_s \ln s$ .
- What colour-singlet gluon configurations can we exchange?

$gg$  : C-even, dominant.  
QCD Pomeron.



$ggg$  : can be C-odd.  
QCD Odderon.



- Odderon: prediction of QCD\*. Detailed calculation  $\rightarrow$  does not die away with energy. Although moving to non-perturbative regime will complicate matters, expect on general grounds for Odderon to contribute.
- General QCD expectation - C-odd (i.e.  $\sigma_{pp} \neq \sigma_{p\bar{p}}$ ) exchange at high energies. Yet for a long time no clear experimental evidence. Where is it?

# SEARCHES for ODDERON



## Odderon signals

- **pp scatt** Odderon exch. is a small correction to even-signature term  $(g_{pO})^2$

- **photoproduction of C even mesons**  $\pi^0, f_2, \eta \dots$

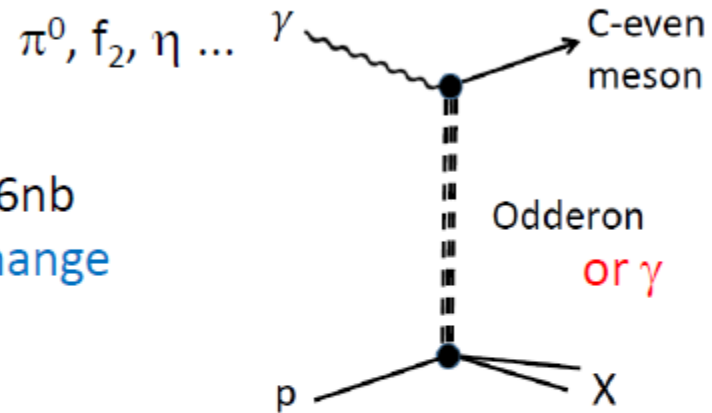
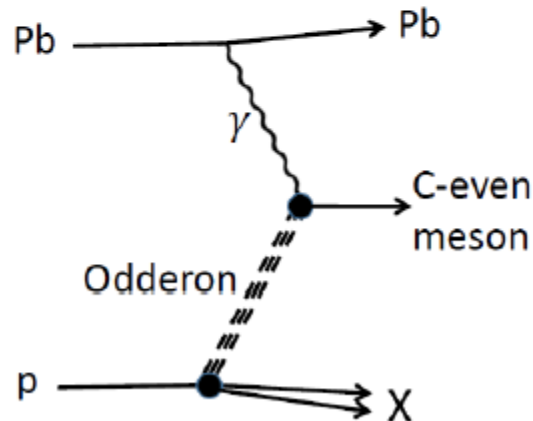
No evidence in HERA data

upper limits  $\sigma(\pi^0)=39\text{nb}$ ,  $\sigma(f_2)=16\text{nb}$

Need to suppress back<sup>gd</sup> due to  $\gamma$  exchange

- **ultraperipheral production in p-Pb collisions**

$Z^2$  in photon flux



## Odderon signal in p-Pb collisions?

$d\sigma/dy_M _{y_M=0}$	Expected upper limits [ $\mu\text{b}$ ]
$\pi^0$	7.4
$\eta$	3.4
$f_2(1270)$	3.0

# Summary & Outlook

- No immediate plans for a future  $\gamma\gamma$  collider, but the LHC is already a photon-photon collider!
- The  $\gamma\gamma$  initial state naturally leads to exclusive events, with intact outgoing protons.
- Theory well understood, and use as highly competitive and clean probe of EW sector and BSM physics already demonstrated at LHC. Much further data with tagged protons to come.
- Such studies equally possible (with higher  $s_{\gamma\gamma}$ ) at FCC.
- - **SuperChic** - a MC event generator for CEP processes.
  - Unified platform for QCD-induced, photoproduction and photon-photon collisions.
  - Fully differential treatment of survival factor.
- Photon-mediated processes provide a very efficient tool to probe various aspects of the SM and BSM physics.







*BACKUP*

# Motivation: photon-induced CEP

- **Photon-initiated** CEP of particular interest:

- ★ Very well understood initial state, via **equivalent photon approximation**:

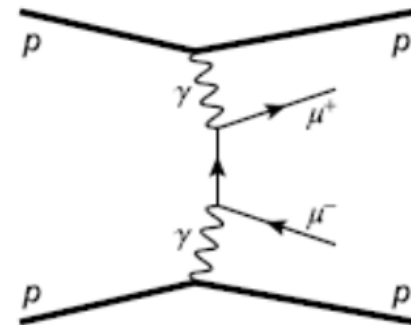
(Well known) EM Form Factors

$$n(x_i) = \frac{1}{x_i} \frac{\alpha}{\pi^2} \int \frac{d^2 q_{i\perp}}{q_{i\perp}^2 + x_i^2 m_p^2} \left( \frac{q_{i\perp}^2}{q_{i\perp}^2 + x_i^2 m_p^2} (1 - x_i) F_E(Q_i^2) + \frac{x_i^2}{2} F_M(Q_i^2) \right)$$

Photon flux

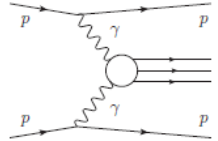
- ★ Low photon  $Q^2 \Rightarrow$  large proton-proton impact parameter  
impact of QCD ('survival factor') small and under control.


→ LHC as a  $\gamma\gamma$  collider! **Clean probe** of **BSM**  
with EW couplings.




# FP420 AND RESURRECTION OF 'DIFFRACTIVE HIGGS' (15 YEARS ON)

- searching for lower mass new objects in CEP requires far away ( $\sim 400\text{m}$ ) FDs





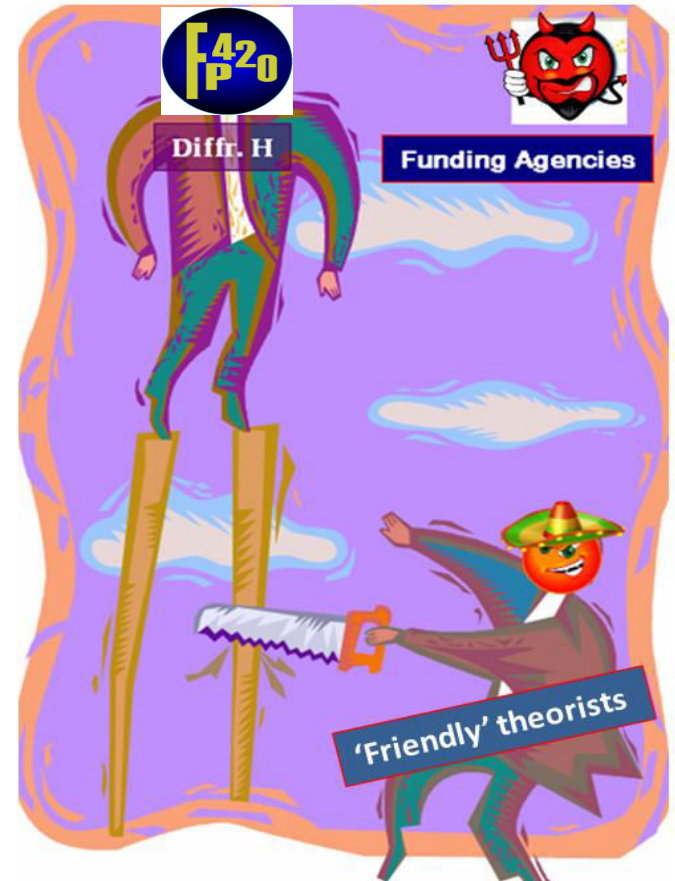
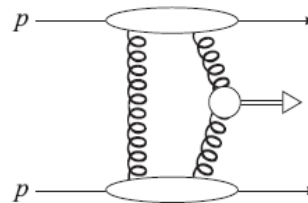
The FP420 R&D project (2004-2009)



- FP420 was a joint R&D collaboration between CMS and ATLAS to develop a forward proton detector system to tag outgoing protons.
- Key questions:
  - Can suitable forward detectors be placed close to the LHC beam
  - What is the physics potential of these detectors?
  - Will they cover an interesting region of Higgs mass?
- Final report is available at [JINST 4:T10001,2009 \[arXiv:0806.0302\]](https://arxiv.org/abs/0806.0302)

❖ QCD-initiated production: potential for e.g. exclusive jet and F studies analysed (though there are more).

- ★ **Jets:**  $gg$  colour-singlet initial state range of unique QCD studies.
- ★ **Higgs:** completely unseen mode, Higgs properties (CP, couplings) via independent method.



96 GeV-'light Higgs'

