

## Heavy Quarkonia: as Seen through the Eyes of Central Exclusive Production at the Tevatron and LHC



Durham  
University

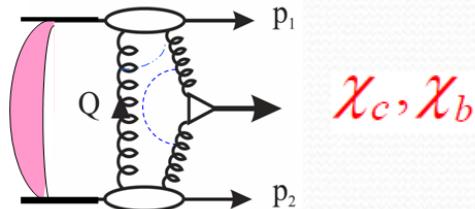
V.A. Khoze (IPPP, Durham)



June 26<sup>th</sup> 09

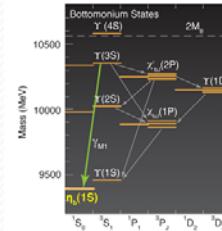
(Based on collaboration with L. Harland-Lang, A..Martin, M.Ryskin and W.J. Stirling)

**main aim:** to demonstrate that CEP can open a new way to study the properties of heavy quarkonia , and, in particular, to serve as a spin-parity analyser



## PLAN

1. Introduction
2. Central Exclusive Production as a heavy meson spin-parity analyser.
3. What is known from the general rules (Regge theory).
4. Expectations within the pQCD Durham approach.
5. Interpretation of the CDF results on charmonium *CEP*.
6. New Run of Durham studies (*still preliminary*).
7. Prospective measurements.
8. A few remarks about the models for soft diffraction.
9. Conclusion.



## INTRODUCTION

Why an interest to the CEP of  $\chi_c, \chi_b$  ?

- Testing ground for the formalism of CEP used to evaluate the New Physics signals (e.g. 'Diffractive Higgs')
- Open issues in Quarkonium Spectroscopy, such as  $\chi_b$  quantum numbers. New way to address Quarkonium Physics.
- New Encouraging CDF results on CEP of the  $\chi_c$  . . . . (Mike)

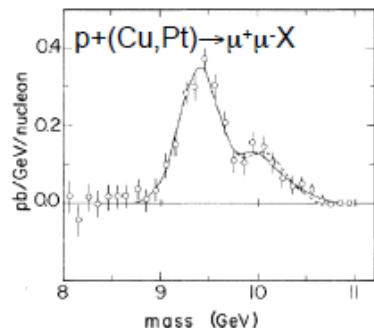
### Heavy Quarkonia

Traditional testing ground for various aspects of QCD

- NRQCD, QCDME, Lattice QCD, QCD sum rules, potential models,...
- Large NLO..... PT corrections.
- P-states- sensitivity to the derivatives of the wave function, relativistic effects....

Bottomonium history started 30 years ago

( PRL 39, 242 (1977) and PRL 39,1240 (1977) )



$$M(Y) = 9.40 \pm 0.013$$

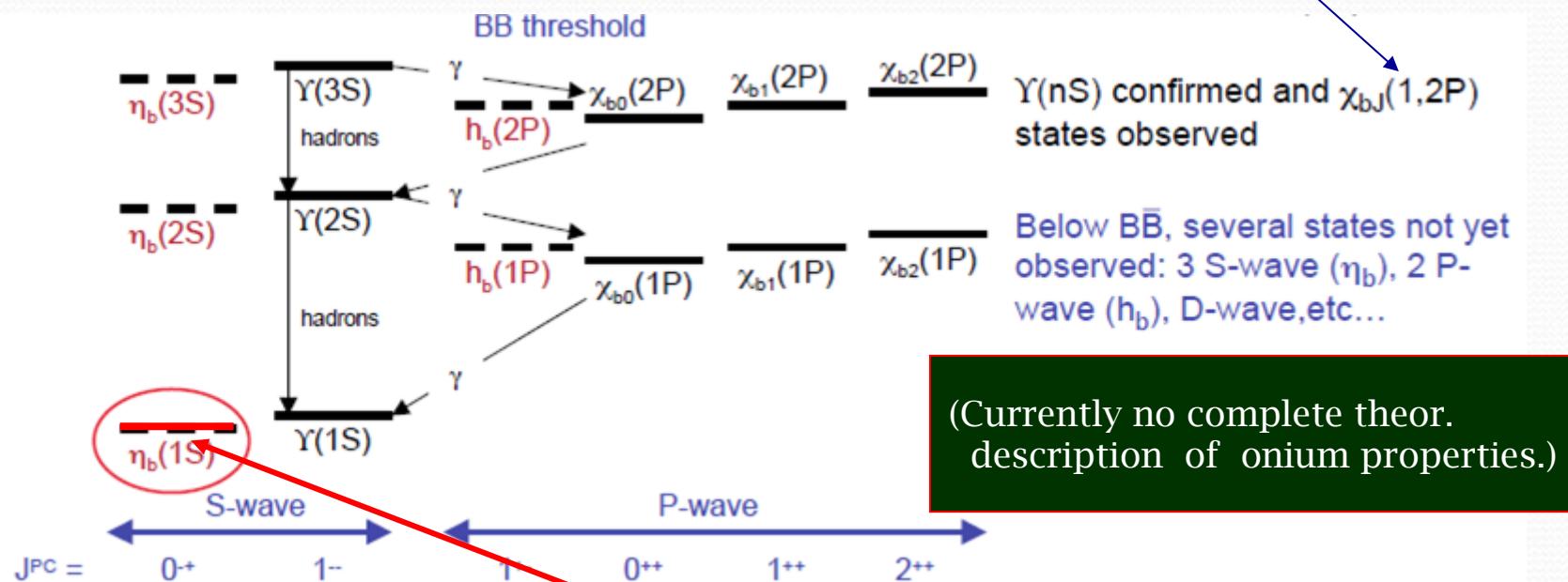
$$M(Y') = 10.00 \pm 0.04$$

$$M(Y') = 10.43 \pm 0.12$$

FNAL, E288

30 years later....

(spins- still unconfirmed)



(Currently no complete theor.  
description of onium properties.)

(BABAR (2008))

(Still some puzzles)

# The heaviest and most compact quark-antiquark bound state in nature



PDG

Summary Tables

## Meson Summary Tables

in the 2008 Review of Particle Physics

 **$\chi_{b0}(1P)$**  [d] $I^G(J^{PC}) = 0^+(0^{++})$  $\cancel{I \text{ needs confirmation}}$ Mass  $m = 9859.44 \pm 0.42 \pm 0.31$  MeV **$\chi_{b0}(1P)$  DECAY MODES**Fraction ( $\Gamma_i/\Gamma$ )

Confidence level

 $p$  $\gamma \ U(1S)$  $\cancel{<6 \%}$ 

90%

391

 **$\chi_{b1}(1P)$**  [d] $I^G(J^{PC}) = 0^+(1^{++})$  $\cancel{I \text{ needs confirmation}}$  $\gamma \ U(1S)$  $(35 \pm 8) \%$ 

423

 **$\chi_{b2}(1P)$**  [d] $I^G(J^{PC}) = 0^+(2^{++})$  $\cancel{I \text{ needs confirmation}}$ Mass  $m = 9912.21 \pm 0.26 \pm 0.31$  MeV **$\chi_{b2}(1P)$  DECAY MODES**Fraction ( $\Gamma_i/\Gamma$ ) $p$  (MeV/c) $\gamma \ U(1S)$  $(22 \pm 4) \%$ 

442



- Quantum number filter/analyser.  
( 0<sup>++</sup> dominance ; C,P-even),
- Clean few-particle final state,
- Favourable background conditions.  
(theoretical estimates,  $\gamma\gamma$ - data).
- New leverage -proton momentum correlations

### Potential (theoretical) problems



- Higher sensitivity to low scales- ‘usual suspects’
- Stronger dependence on Enhanced Screening effects  
(larger  $s/M_\chi^2$  )

## What we know from Regge theory (KKMR-2003)

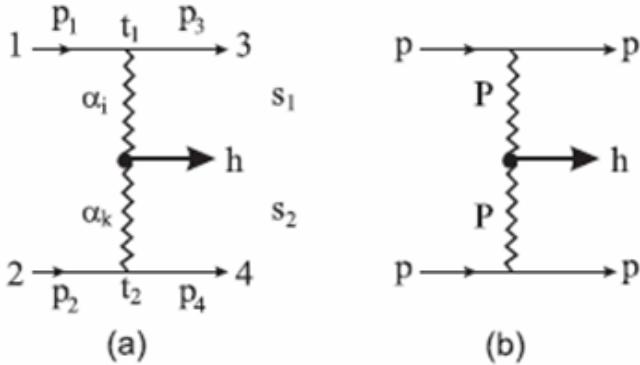


Figure 1: (a) The central production of a state  $h$  by double-Reggeon exchange. (b) The double-Pomeron exchange contribution to  $pp \rightarrow p + h + p$ , which dominates at high energies, where the + signs are used to indicate the presence of Pomeron-induced rapidity gaps.

$$J^P(h) = 0^+$$

Vertex Coupling  $g_{ik}^h = f_{0+}(p_{3\perp}^2, p_{4\perp}^2, \vec{p}_{3\perp} \cdot \vec{p}_{4\perp})$  depends on dynamics

$$J^P(h) = 0^-$$

$\sigma \sim |t_1||t_2| \sin^2 \phi$  observed for  $\eta, \eta'$  by WA102 Group  
(450 GeV, pp CERN Omega Spectrometer,)

$$J^P(h) = 1^+$$

For small  $p_{it}$

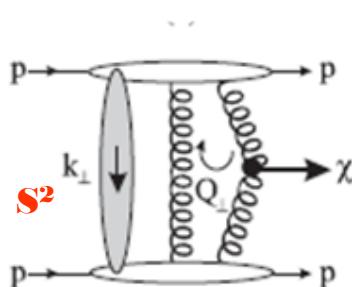
$$g_{PP} = a_{\lambda=0} (p_{3t}^2 - p_{4t}^2) / M^2 (\vec{p}_{3t} \times \vec{p}_{4t}) \vec{e} + f_{\lambda=1} (\vec{K} \times \vec{n}) \vec{e} / M, \quad \vec{K} = (\vec{p}_3 - \vec{p}_4)$$

- Cross section tends to zero at low  $K_t$  (also Petrov et al-04, Szczerba et al)
- Dominantly produced in the helicity-one state
- Coincide with the NCVC model expectation by F. Close et al (1999)
- Agree with the WA102 data on  $f_1(1420)$  and  $f_1(1285)$

# What we expect within the framework of the Perturbative Durham formalism

(KMR-01, KKMR-03, KMRS-04)

## O++ -case



$$T = A\pi^2 \int \frac{d^2 Q_\perp P(\chi(0^+))}{Q_1^2 (\vec{Q}_\perp - \vec{p}_{1\perp})^2 (\vec{Q}_\perp + \vec{p}_{2\perp})^2} f_g(x_1, x'_1, Q_1^2, \mu^2; t_1) f_g(x_2, x'_2, Q_2^2, \mu^2; t_2),$$

$$A^2 = 8\pi\Gamma(\chi \rightarrow gg)/M_\chi^3 \quad *K_{NLO}$$

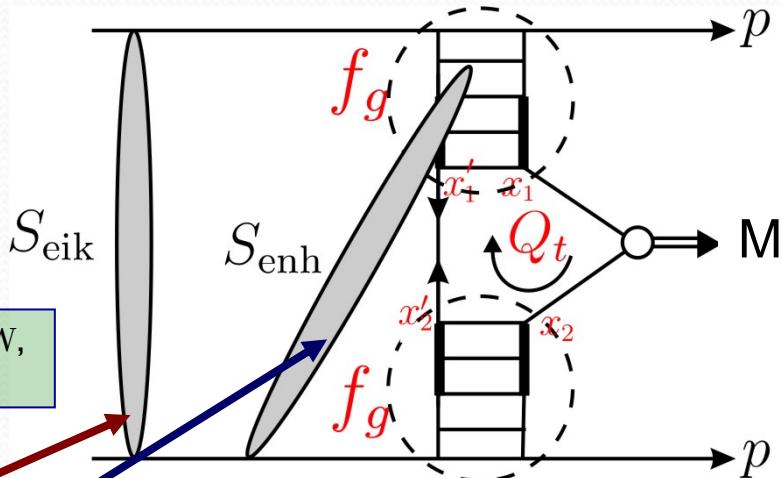
$$P(\chi(0^+)) = (\vec{Q}_\perp - \vec{p}_{1\perp}) \cdot (\vec{Q}_\perp + \vec{p}_{2\perp}).$$

- Strong sensitivity to the polarization structure of the vertex in the bare amplitude.
- In the on-shell-gluon approximation spin-1 is excluded by Landau-Yang theorem. Should lead to a strong suppression.
- For forward going protons in the non-relativistic quarkonium approximation, CEP of the spin-2 meson is strongly suppressed ( $J_z=0$  selection rule) KMR-01
- Absorption is sizeably distorted by the polarization structure (affects the b-space distr.) KMR-02, KKMR-03
- $\chi_c, \chi_b$ -production is especially sensitive to the effects of enhanced absorption
  - larger available rapidity interval
  - lower scale → larger dipole size → larger absorption  
( $S^2_{\text{enh}}$  for  $\chi_c$  at the Tevatron is expected to exceed that for the Higgs at the LHC)



“soft” scattering can easily destroy the gaps

$S^2 \rightarrow$  absorption effects -necessitated by unitarity



Everybody's happy (KMR, GLM, FHSW, Petrov et al, BH, GGPS, Luna...MCs)

eikonal rescatt: between protons  
enhanced rescatt: involving intermediate partons

soft-hard factoriz<sup>n</sup>  
conserved broken

Subject of hot discussions :  $S^2_{\text{enh}}$



# Interpretation of CDF results

(CDF Collaboration, arXiv:0902.1271 [hep-ex], PRL in press)

(Mike)

Assuming that all events are originated from CEP of  $\chi_c(0^+)$  ( limited acceptance)  
(used CHIC MC- Durham based )

$$\text{CDF} \quad \frac{d\sigma(\chi_c)}{dy} \Big|_{y=0} = (76 \pm 14) \text{ nb}$$

KMRS -2004: 130 nb → 90 nb (PDG-2008)

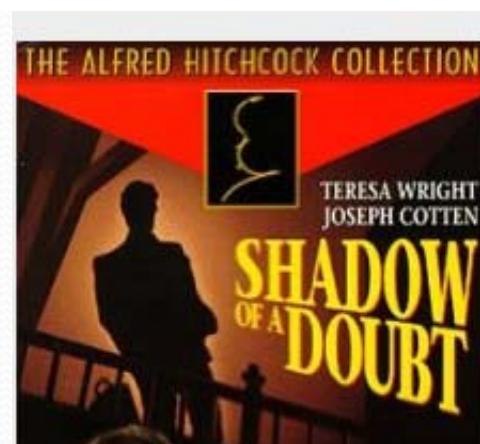


Signal based on:  $\chi_c \rightarrow J/\psi + \gamma$



A certain preference to  $0^+$  in the  $(J/\psi + \gamma)$  mass distribution

Too good to be true ?!



## Devil's Advocate Questions



1. How do we know which particular P-state has been found ?  
(P-states are not clearly separated experimentally )
2. Is reconstruction based on CHIC MC still acceptable, what if not  $J^{PC} = 0^{++}$  ?
3. Are we close to the CEP prescription (role of low mass SD and DD)  
 $\chi_c(0^+)$  dominates CEP, but

Yes, due to the record  
CDF gap coverage (7.4)  
(KMRS-04)

$$Br(\chi_c(0^+) \rightarrow J/\psi + \gamma) = (1.28 \pm 0.11)\%$$



$$Br(\chi_c(1^+) \rightarrow J/\psi + \gamma) = (36.0 \pm 1.9)\%$$

$$Br(\chi_c(2^+) \rightarrow J/\psi + \gamma) = (20.0 \pm 1.0)\%$$

- On-mass-shell  $1^+$  production is forbidden due to Landau-Yang theorem,  
but what about off-mass-shell effects ?  
Recently- renewal of interest ([R.Pasechnik](#), [A.Szczerba](#), [O.Teryaev-09](#))  
Still numerically small
- Within the non-relativistic framework for forward going protons  
 $2^+$  is strongly suppressed, but what about non-forward protons and relativistic effects ?
- Important phenomenon- absorptive corrections are quite sensitive to the meson spin-parity  
( studied before in the context of scalar/pseudoscalar Higgs-[KKMR04](#) )

### Issues addressed:

- New CHIC MC for all  $c\bar{c}$  P-states.
- Absorption effects for CEP of the  $0^+, 1^+, 2^+, 0^-$   $c\bar{c}$  - states revisited
- Proton angular correlations for different  $0^+, 1^+, 2^+, 0^-$   $c\bar{c}$  -states.
- Expectations for the CER of the  $0^+, 1^+, 2^+, 0^-$   $b\bar{b}$  -states.

As compared to the previous K(KMR)S studies:

- More comprehensive calculation of the absorption effects using the new KMR-07/08 model for soft diffraction (including the enhanced screening).
- New calculational routine for implementing polarization structure in the b-space.
- New experimental/theoretical results for the parameters of heavy quarkonia, in particular  $\Gamma(\chi \rightarrow gg)$ .

**PRELIMINARY**

## CHARMONIUM

- The final state ( muon) distributions in the  $(J/\psi + \gamma)$  system are sensitive to the meson spin, but after imposing the CDF cuts this dependence is strongly reduced. (L. Harland-Lang, W.J. Stirling).
- Cross-section ‘reconstruction’ is safe 😊
- Spin is not discriminated via the  $(J/\psi + \gamma)$  decay products 😱
- We need to measure better spin-parity analysing final state:  $\pi\pi, KK, p\bar{p}$  or outgoing proton momentum correlations KMR-04

For normalization purposes- scalar case at the Tevatron

$$\langle S^2_{eik}(0^+) \rangle \approx 0.06$$

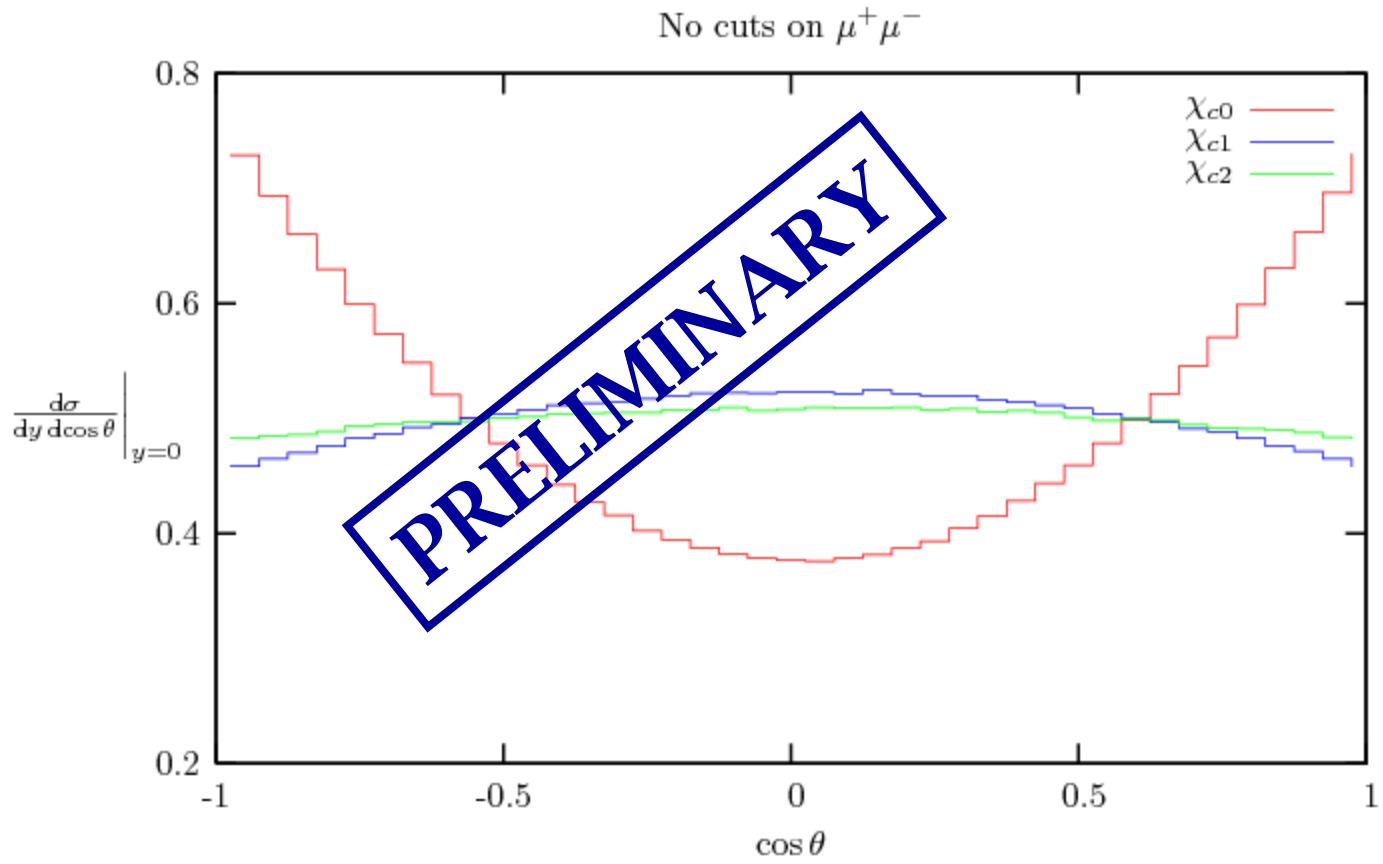
$$\langle S^2_{eff}(0^+) \rangle \approx 0.02$$

( $\Delta=2.3$ )

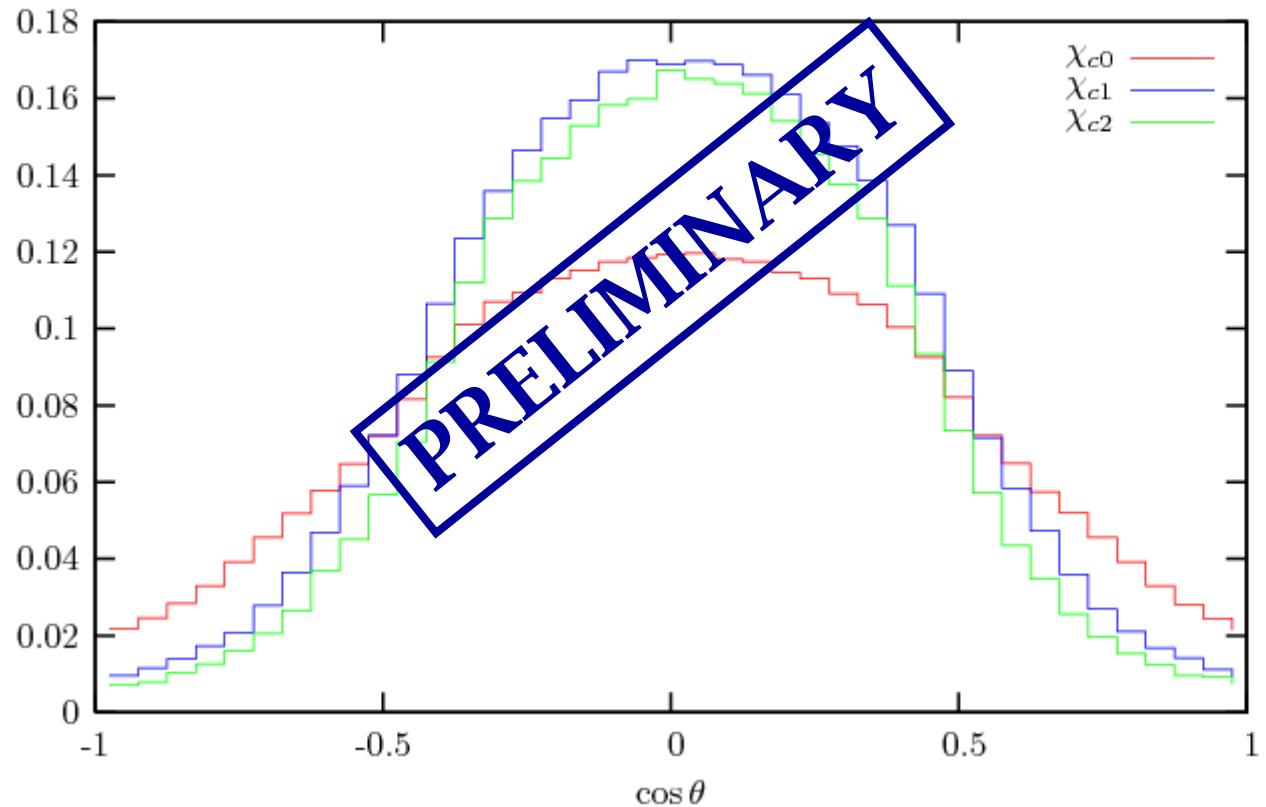
$$d\sigma(0^+)/dy|_{y=0} \approx 90 \text{ nb}$$

$$d\sigma(0^+)/dy|_{y=0} \approx 35 \text{ nb}$$

- Still within Durham approach-uncertainties, recall, in particular  $(f_g)^4$ - effect  
Reasons to believe that enhanced absorption is overestimated- KMR-09



$|\eta| < 0.6$  and  $p_{\perp} > 1.4$  GeV cuts on  $\mu^+\mu^-$



Some of the ‘typical’ uncertainties cancel in the ratios

- ❖  $\sigma(1^+)/\sigma(0^+) \approx \langle p_t^2 \rangle / M_\chi * \langle S_{eik}^2(1) \rangle / \langle S_{eik}^2(0) \rangle * R_{NLO}^1 \approx 0.05$

- $\langle S_{enh}^2(J) \rangle$  - the same for all J within  $\sim 20\%$  accuracy,

- $\langle S_{eik}^2(1,2) \rangle / \langle S_{eik}^2(0) \rangle \approx 2.5; \langle p_t \rangle \approx 0.5 \text{ GeV}.$

Production is more peripheral: zero at  $\vec{b} = 0$ , where the absorption is largest at larger  $b$  absorption is small anyway.

- ❖  $\sigma(2^+)/\sigma(0^+) \approx (\langle p_t^2 \rangle / Q_t^2)^2 * \langle S_{eik}^2(2) \rangle / \langle S_{eik}^2(0) \rangle * R_{NLO}^2 \approx 0.05$   
( $\langle Q_t \rangle = 1 \text{ GeV}$ )

- ❖  $\sigma(0^+ \rightarrow J/\psi + \gamma) : \sigma(1^+ \rightarrow J/\psi + \gamma) : \sigma(2^+ \rightarrow J/\psi + \gamma) = 1 : 1.4 : 0.77$

Then  $d\sigma(0^+, 1^+, 2^+)/dy \approx 110 \text{ nb}$ , as compared to experiment:  $(76 \pm 14) \text{ nb}$

After all ,



keeping in mind all uncertainties

## Energy Dependence

- Expected to be weak, since the rise of the gluon density at low  $x$  is compensated by stronger enhanced screening.

$$d\sigma(\chi_{c0})/dy \approx 50 \text{ nb at the LHC}$$

- Test of the enhanced absorption (less model dependent):

$$\sigma(\chi_{LHC})/\sigma(\chi_{Tevatron})$$

various uncertainties cancel (NLO effects, width,...)

- $\sigma(\chi_{c2})/\sigma(\chi_{c0})$  decreasing with energy ( $\langle Q_t \rangle$  increasing)

## Momentum correlations between outgoing protons

- Separation of different meson states (irrespectively of the final state), seen in the CEP of light mesons (WA102 Collab).

Recall the results of Regge theory at low transverse momenta  $(p_{3,4})_t$

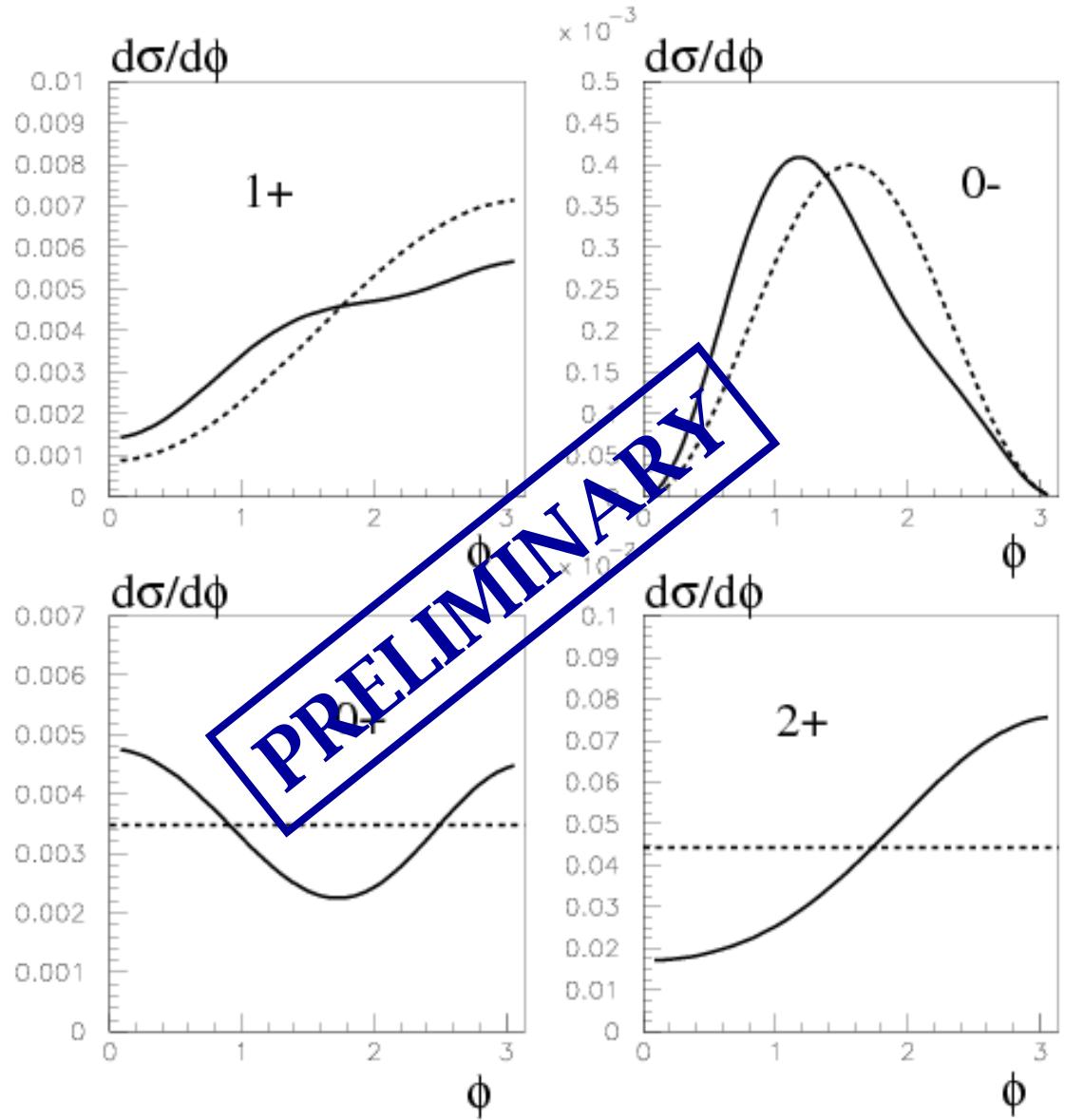
$$d\sigma(0^+)/d\phi \simeq \text{const};$$

$$d\sigma(0^-)/d\phi \simeq \sin^2 \phi;$$

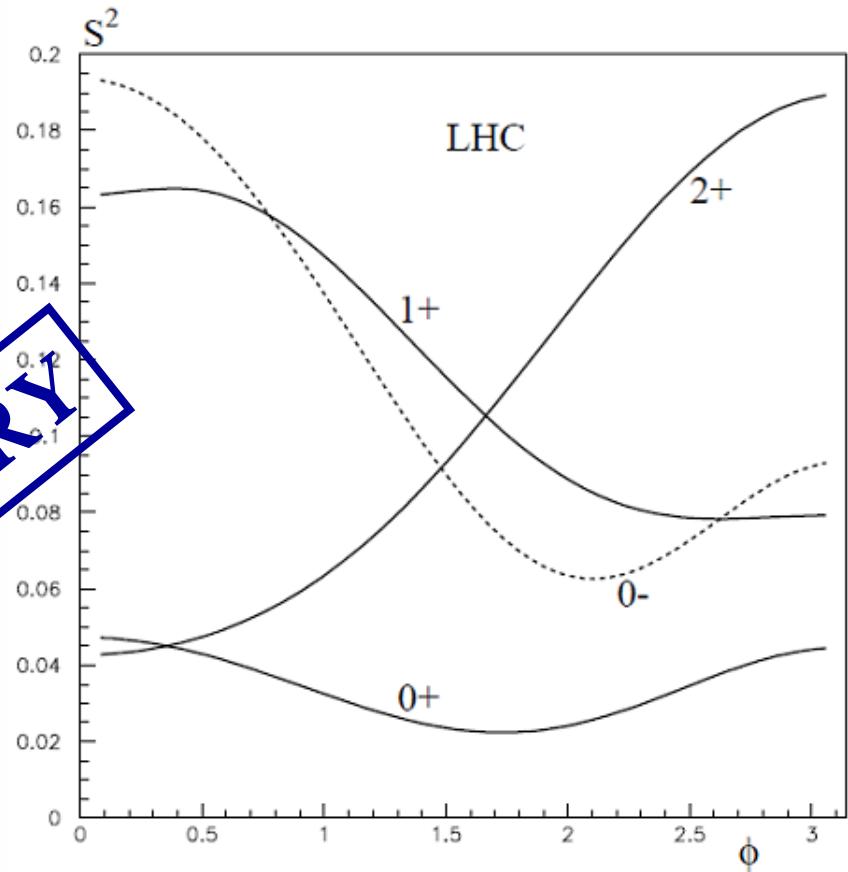
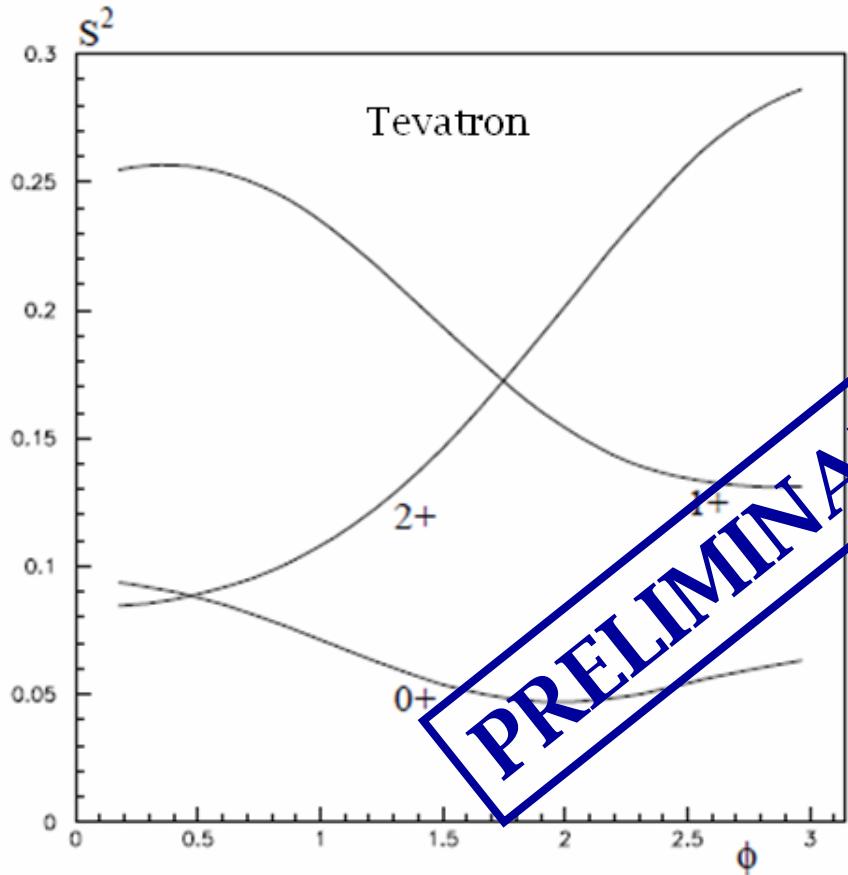
$$d\sigma(1^+)/d\phi \simeq (\vec{p}_3 - \vec{p}_4)^2_t;$$

$$d\sigma(2^+)/d\phi \simeq \text{const};$$

- Some dependence on the choice of unintegrated gluon densities ([KMR-02](#), [KKMR-03](#))
- Serious modification by the absorption effects  
(amount of suppression strongly depends on the impact parameter  $\bar{b}$ )



$\phi$ - dependence of  $S^2_{eik}$



- More peripheral nature of  $0^-, 1^+, 2^+$  production.
- Manifestations of the diffractive dip.
- Dependence on the  $p_t$  -cuts studied in [KMR02, KKMR-03](#) for the H-case.
- $S^2$  *enh* weakly depends on  $\phi$ .

for completeness

$$\frac{d\sigma(\eta_c)/dy}{|_{y=0}} = 0.5 \text{ nb at the Tevatron}$$

1.5 nb at the LHC

note: some decay modes are quite sizeable (  $\rho\rho$ ,  $KK\pi$ ,  $\eta\pi\pi..$  a few % level)

Spin-parity is not confirmed (yet) for some new mesons

- $X(3872)$ :  $1^{++}$  or  $2^{-+}$  ?
- $Y(4140)$  (decay  $J/\Psi\phi$ ) (molecule or tetra-quark or..?)
- $X(3940)$ ,  $Y(3940)$  and  $Z(3930)$ 
  - $\chi_{c1}(2P)$
  - $\chi_{c2}(2P)$
  - $J^{P+}$
- $\eta_c(3S)$  or ?
  - $X(3872)$  ,  $Y(3940)$   $Y(4260)$
  - $J/\psi$

# BOTTOMONIUM

Recall , the decay  $\chi_{b0}(1P) \rightarrow Y(1S)\gamma$  has not been seen (yet)

but  $Br(\chi_{b0}(2P) \rightarrow Y(2S)\gamma) = (4.6 \pm 2.1)\%$



According to evaluation by J.T. Laverty et al (2009)

$$\Gamma(\chi_{b0} \rightarrow gg) = 3.7 MeV ,$$

$$Br(\eta_b \rightarrow \gamma\gamma) = 3.4 * 10^{-5} \quad (\text{Exp. } Br(\eta_{c0} \rightarrow \gamma\gamma) \simeq 2 * 10^{-4} ).$$



Higher scale → better PT description

Smaller role of relativistic effects, better knowledge of gluon densities

- $1^+$ -is practically filtered out (strong  $M^2$  suppression)
- $2^+$  weak-  $\langle Q_t \rangle^2$  dependence
- Enhanced absorption- weaker ( $\sim 2$  times)

## Expectations for $\chi_b, \eta_b$

$d\sigma / dy |_{y=0}$  (in pb)

	$\chi_{b0}$	$\chi_{b1}$	$\chi_{b2}$	$\eta_b$
<b>Tevatron</b>	400	3	10	5
<b>LHC</b>	700	7	20	15

$\rightarrow D^0$  -modes (small background)  
 (significant production of  $D^0$  mesons from both the  $\chi_{b1}(1P)$  and  $\chi_{b1}(2P)$   
 NRQCD CLEO III- 2008)

## PROSPECTIVE MEASUREMENTS

- A clear way to resolve the issue of  $\chi_c$  spin-parity identification will be to search for the two-body decays:

$$Br(\chi_{c0} \rightarrow \pi\pi, K^+K^-) \simeq 1.3\% \quad \chi_{c1}, \eta_c \cancel{\rightarrow} \pi\pi, KK \quad Br(\chi_{c2} \rightarrow \pi\pi, K^+K^-) \simeq 0.3\%$$

$$Br(\chi_{c0} \rightarrow p\bar{p}) \simeq 2 * 10^{-4} \quad Br(\chi_{c1} \rightarrow p\bar{p}) \simeq 6.6 * 10^{-5} \quad Br(\chi_{c2} \rightarrow p\bar{p}) \simeq 6.7 * 10^{-5}$$

$$Br(\eta_c \rightarrow p\bar{p}) \simeq 0.13\%$$

- Tagged forward protons: **spin-parity ID** of old and new heavy meson states, detailed tests of absorption effects
- With sufficient statistics of  $\gamma\gamma$  CEP, the measurement of the ratio

$$\sigma(\chi_b)/\sigma(\gamma\gamma)$$

can be quite instructive (the same mass range, various uncertainties cancel).

# UNCERTAINTIES



## Known Unknowns

- N(N)LO- radiative effects (K-factors etc..)  
‘...possible inadequacy of PT theory in  $\alpha_s$  ...’, R.Barbieri et al-1980
- ‘Right’ choice of gluon densities, in particular at so low scales as in the  $\chi_c$  case ( potentiality of a factor of ~3 rise for the H-case ) .
- Complete model for calculation of enhanced absorption.
- $\chi_b$  -experimental widths, decays...

## Unknown Unknowns

- Non- pQCD effects in the meson characteristics.  
Currently no complete description of heavy quarkonium characteristics.  
‘Two gluon width does not tell the whole story.’
- Gluons at so low scales, surprises are not excluded at all.



Factor of 5 up or down  
(at best)



Far more theoretical papers than the expected number  
of the CED produced Higgs events

(Laslo)

'Well, it is a **possible** supposition  
'You think so, too ?'  
'I did not say a **probable** one'



# Selection Criteria for the Models of Soft Diffraction

(Aliosha, Asher, Uri,Sergey)



- We have to be **open-eyed** when the soft physics is involved.  
Theoretical models in the strong coupling regime contain various assumptions and parameters.
- Available data on soft diffraction **at high energies** are still fragmentary,  
especially concerning the (low mass) diffractive dissociation.

## A viable model should:

- incorporate the inelastic diffraction :SD, DD  
(for instance 2-3 channel eikonal **KMR** or **GLM(M)**)
- describe all the existing experimental data on elastic scattering and SD ,DD and **CED**  
at the Tevatron energies and below (**KMR**; **GLM(M)** )
- be able to explain the existing CDF data on the HERA-Tevatron factorization breaking  
and on the QED production of the di-jets, di-photons,  $\chi$ ,  $J/\psi$ ,  $Y$ .., lead. neutrons at HERA
- provide testable **pre-dictions** or at least **post-dictions** for the Tevatron and HERA

So far (only) **Durham** model has passed these tests.

Only a large enough data set would impose the **restriction order** on the theoretical models and to create a confidence in the determination of  $S^2$ .

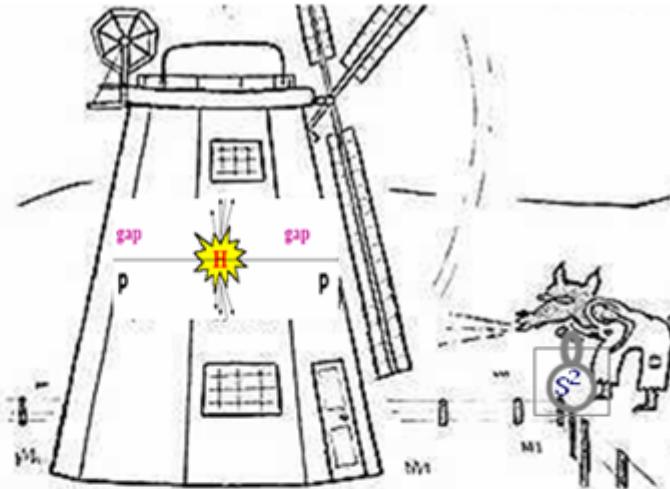


Program of Early LHC measurements **(KMR)**

LET THE DATA TALK !

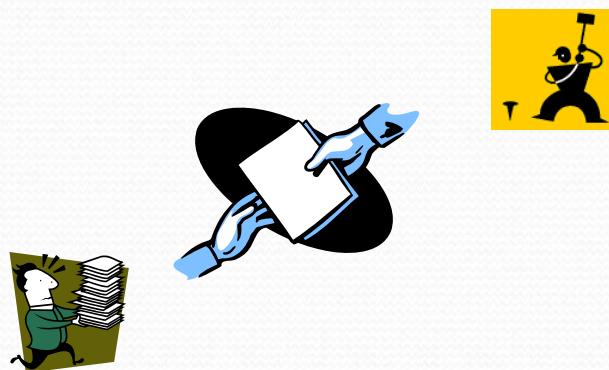
## Who's Afraid of the Big, Bad Wolf?

- $S^2$  does not affect the signal-to-background ratio- for all irreducible backgrounds (signal evidence is much less affected)
- Main reduction of the signal (factor of ~50) comes from the experimental requirements ( cuts and efficiencies...) which are currently known for the inclusive environment. Further progress with hard/soft -ware for the CEP processes can be expected.  
More experimental work needed
- Experimentally we have not seen (at least so far) any evidence in favour of large enhanced absorption (KKMR, KMR- 2001-2009).
- Durham selection of the UPDF is quite conservative. Due to the  $(fg)^4$  behaviour- rise up to a factor of 3 (Cox et al, KMR).  
New studies underway. Up to two orders of magnitude rise in the popular BSM Higgs models.
- We should be careful with relying on the NLO corrections (e.g. BBKM-06). Can be misleading when large parameters are involved.  
(textbook example: non-relativistic Coulomb correction)



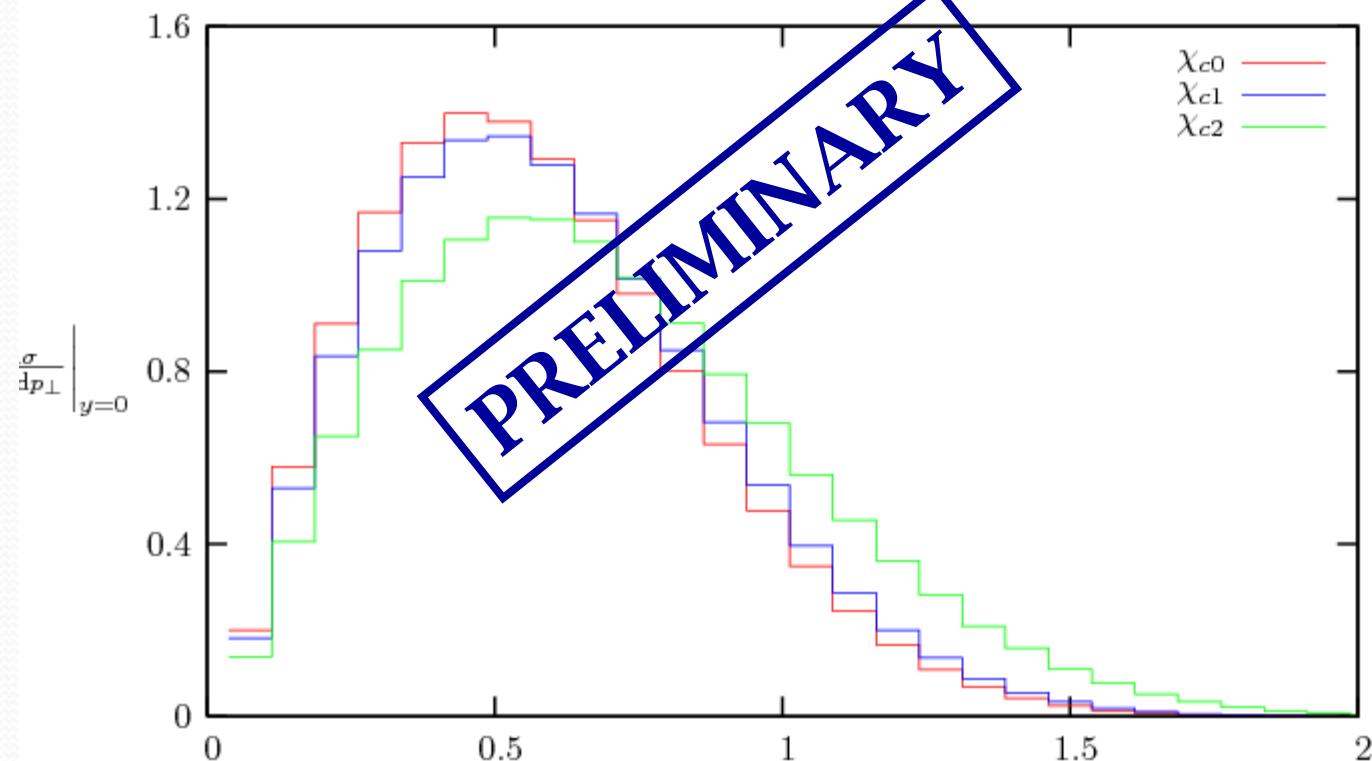
## CONCLUSION

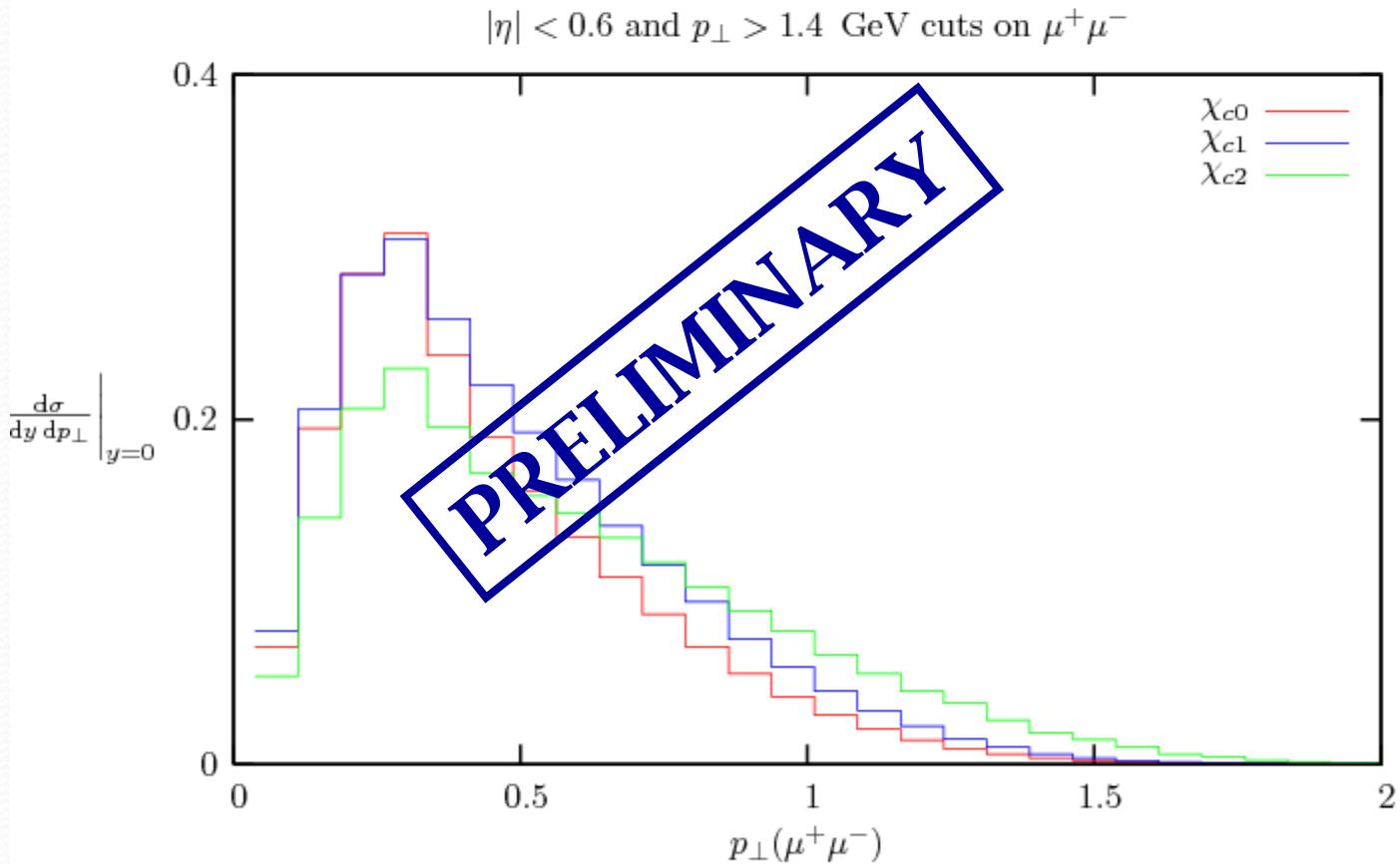
- CDF data on **CEP** of the  $\chi_c$  are in a broad agreement with the Durham results.
- **CEP** of heavy mesons - a new way to study quarkonium spectroscopy as well as to address the physics of absorption; can help to establish the nature of newly discovered heavy states.
- Promising prospects of studying heavy meson **CEP**, especially with tagged forward protons.
- Currently active studies are still in progress (both in theory and experiment).



# Backup

No cuts on  $\mu^+\mu^-$





**$\chi_{b1}(2P)$**  [d]

$I^G(J^{PC}) = 0^+(1^{++})$   
 $J$  needs confirmation.

Mass  $m = 10.25546 \pm 0.00022 \pm 0.00050$  GeV

$m_{\chi_{b1}(2P)} - m_{\chi_{b0}(2P)} = 23.5 \pm 1.0$  MeV

<b><math>\chi_{b1}(2P)</math> DECAY MODES</b>	Fraction ( $\Gamma_i/\Gamma$ )	Scale factor	$p$ (MeV/c)
$\omega \ U(1S)$	( $1.63^{+0.38}_{-0.34}$ ) %		135
$\gamma \ U(2S)$	( $21 \pm 4$ ) %	1.5	230
$\gamma \ U(1S)$	( $8.5 \pm 1.3$ ) %	1.3	764
$\pi \pi \chi_{b1}(1P)$	( $8.6 \pm 3.1$ ) $\times 10^{-3}$		238

**$\chi_{b2}(2P)$**  [d]

$I^G(J^{PC}) = 0^+(2^{++})$   
 $J$  needs confirmation.

Mass  $m = 10.26865 \pm 0.00022 \pm 0.00050$  GeV

$m_{\chi_{b2}(2P)} - m_{\chi_{b1}(2P)} = 13.5 \pm 0.6$  MeV

<b><math>\chi_{b2}(2P)</math> DECAY MODES</b>	Fraction ( $\Gamma_i/\Gamma$ )	$p$ (MeV/c)
$\omega \ U(1S)$	( $1.10^{+0.34}_{-0.30}$ ) %	194
$\gamma \ U(2S)$	( $16.2 \pm 2.4$ ) %	242
$\gamma \ U(1S)$	( $7.1 \pm 1.0$ ) %	777
$\pi \pi \chi_{b2}(1P)$	( $6.0 \pm 2.1$ ) $\times 10^{-3}$	229