

Quest for Precision in Simulations for the LHC

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- what the talk is about
- matching & merging with parton showers
- Electroweak corrections
- Revisit Parton Showers
- Revisit Hadronisation
- where we are and where we (should/could/would) go

motivation & introduction

motivation: the need for (more) accurate tools

- to date no survivors in searches for new physics & phenomena
(a pity, but that's what Nature hands to us)
- push into precision tests of the Standard Model
(find it or constrain it!)
- statistical uncertainties approach zero
(because of the fantastic work of accelerator, DAQ, etc.)
- systematic experimental uncertainties decrease
(because of ingenious experimental work)
- theoretical uncertainties are or become dominant
(it would be good to change this to fully exploit LHC's potential)

⇒ more accurate tools for more precise physics needed!

motivation: aim of the exercise

- review the state of the art in precision simulations
- highlight missing or ambiguous theoretical ingredients
- suggest some further studies – experiment and theory

(celebrate success)

(acknowledge failure)

(...)

not covered in this talk

(my apologies)

- Status of BSM simulations:
UFO routinely used in MADGRAPH, HERWIG7, & SHERPA
- Detailed status of individual generators:
will be discussed tomorrow
- Developments of alternative, specialised code:
DIPSY (HI), HEJ (FKL MEs), VINCIA (showering), etc.
- Status of analysis framework and long-term data preservation

matching @ (N)NLO
and
merging @ (N)LO

the aftermath of the NLO (QCD) revolution

- establishing a wide variety of automated tools for NLO calculations

BLACKHAT, GoSAM, MADGRAPH, NJET, OPENLOOPS, RECOLA + automated IR subtraction methods (MADGRAPH, SHERPA)

- first full NLO (EW) results with automated tools
- technical improvements still mandatory

(higher multis, higher speed, higher efficiency, easier handling, . . .)

- start discussing scale setting prescriptions

(simple central scales for complicated multi-scale processes? test smarter prescriptions?)

- steep learning curve still ahead: “NLO phenomenology”

(example: methods for uncertainty estimates beyond variation around central scale)

prequel: parton showers vs. resummation calculations

- various schemes for various logs in analytic resummation
- concentrate on parton shower instead \longleftrightarrow compare with Q_T resummation

(transverse momentum of Higgs boson etc.)

- parametric accuracy by comparing Sudakov form factors:

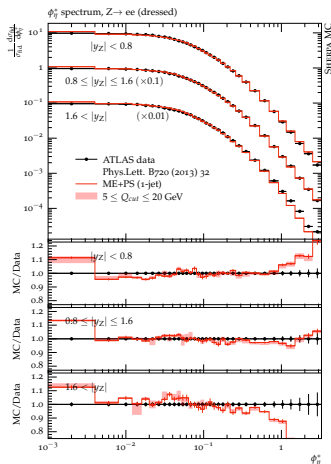
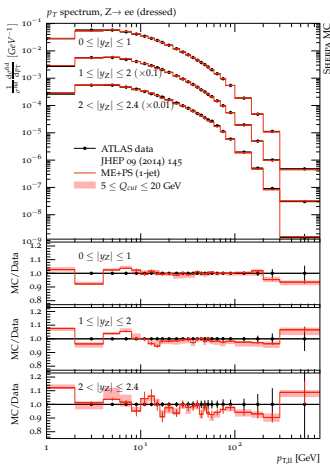
$$\Delta = \exp \left\{ - \int \frac{dk_{\perp}^2}{k_{\perp}^2} \left[A \log \frac{k_{\perp}^2}{Q^2} + B \right] \right\},$$

where A and B can be expanded in $\alpha_S(k_{\perp}^2)$

- showers usually include terms $A_{1,2}$ and B_1 (NLL)
- A_2 often realised by pre-factor multiplying scale $\mu_R \simeq k_{\perp}$

some parton shower fun with DY

(example of accuracy in description of standard precision observable)

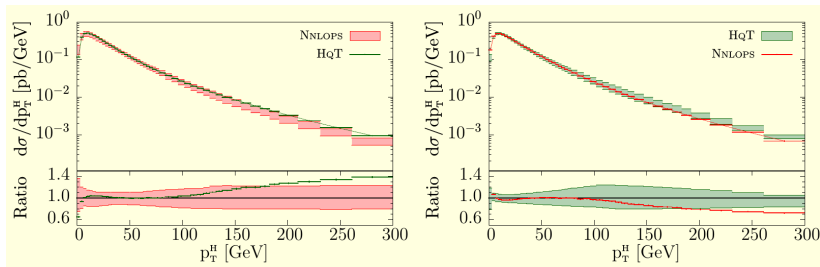


matching at NLO and NNLO

- avoid double-counting of emissions
 - two schemes at NLO: MC@NLO and POWHEG
 - mismatches of K factors in transition to hard jet region
 - MC@NLO: \rightarrow visible structures, especially in $gg \rightarrow H$
 - POWHEG: \rightarrow high tails, cured by h dampening factor
 - well-established and well-known methods
- (no need to discuss them any further)
- two schemes at NNLO: MINLO & UN²LOPS (singlets S only)
 - different basic ideas
 - MINLO: $S + j$ at NLO with $p_T^{(S)} \rightarrow 0$ and capture divergences by reweighting internal line with analytic Sudakov, NNLO accuracy ensured by reweighting with full NNLO calculation for S production
 - UN²LOPS identifies and subtracts and adds parton shower terms at FO from $S + j$ contributions, maintaining unitarity
 - available for two simple processes only: DY and $gg \rightarrow H$

NNLOs for H production: MINLO

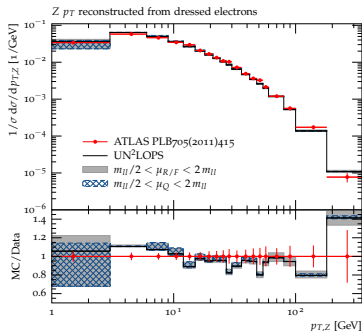
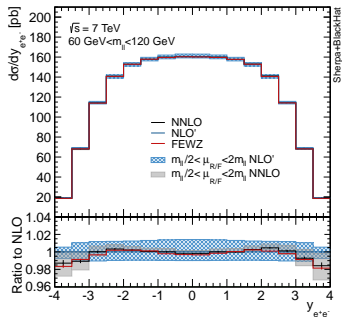
K. Hamilton, P. Nason, E. Re & G. Zanderighi, JHEP 1310



- also available for $Z/W/VH$ production

NNLOs for Z production: UN²LOs

S. Hoche, Y. Li, & S. Prestel, Phys.Rev.D90 & D91



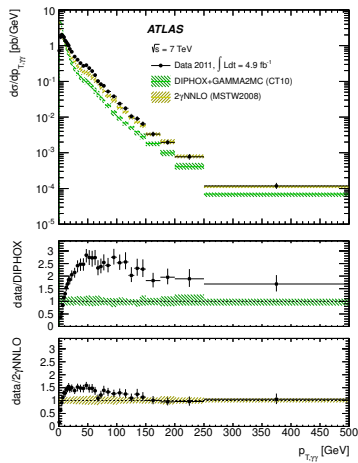
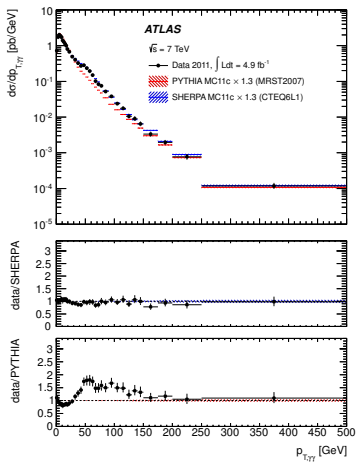
- also available for H production

NNLOs: shortcomings/limitations

- MINLO relies on knowledge of B_2 terms from analytic resummation
→ to date only known for colour singlet production
- MINLO relies on reweighting with full NNLO result
→ one parameter for H (y_H), more complicated for Z , ...
- UN²LOs relies on integrating single- and double emission to low scales and combination of unresolved with virtual emissions
→ potential efficiency issues, need NNLO subtraction
- UN²LOs puts unresolved & virtuals in “zero-emission” bin
→ no parton showering for virtuals (?)

merging example: $p_{\perp,\gamma\gamma}$ in MEPS@LO vs. NNLO

(arXiv:1211.1913 [hep-ex])



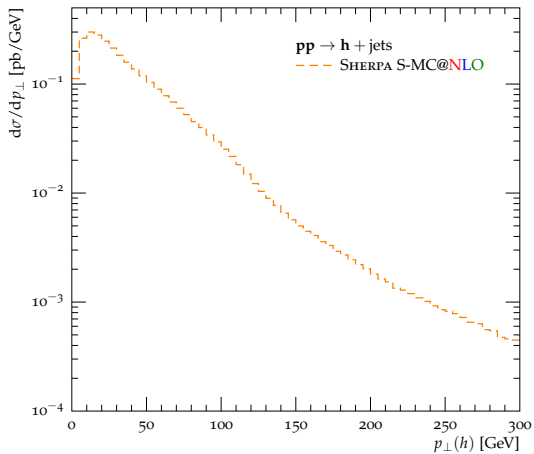
multijet-merging at NLO

- sometimes “more legs” wins over more loops
- basic idea like at LO: towers of MEs with increasing jet multi (but this time at NLO)
- combine them into one sample, remove overlap/double-counting
- maintain NLO and LL accuracy of ME and PS
- this effectively translates into a merging of MC@NLO simulations and can be further supplemented with LO simulations for even higher final state multiplicities
- different implementations, parametric accuracy not always clear
- starts being used, still lacks careful cross-validation

(MEPs@NLO, FxFx, UNLOPs)

illustration: p_{\perp}^H in MEPS@NLO

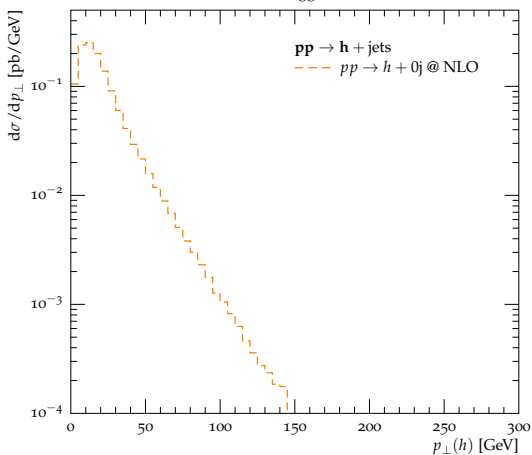
Transverse momentum of the Higgs boson



- first emission by MC@NLO

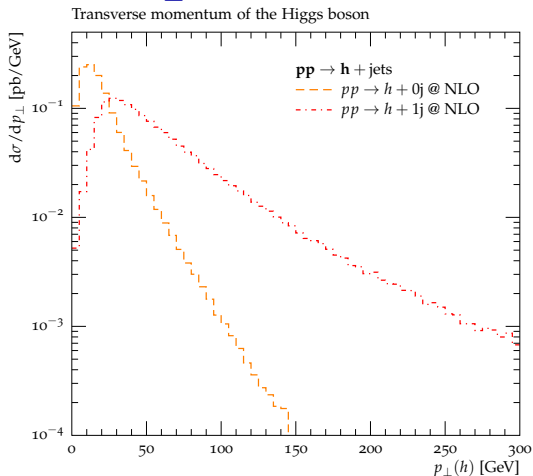
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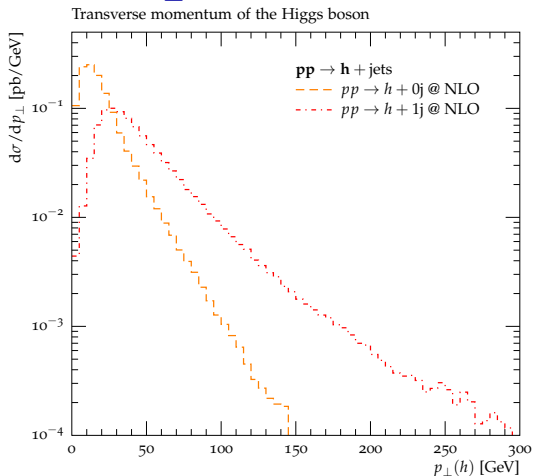
- first emission by MC@NLO, restrict to $Q_{n+1} < Q_{\text{cut}}$

illustration: p_{\perp}^H in MEPS@NLO



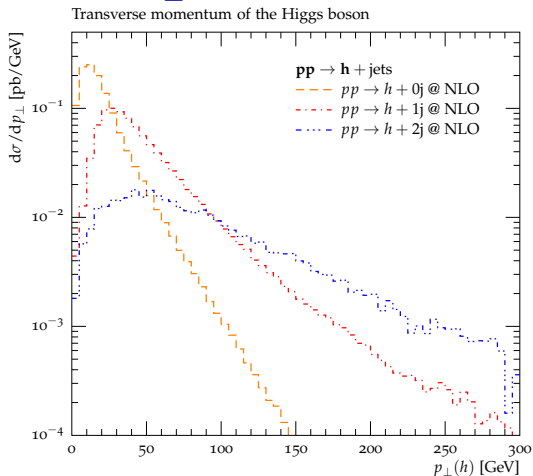
- first emission by MC@NLO, restrict to $Q_{n+1} < Q_{\text{cut}}$
- MC@NLO $pp \rightarrow h + \text{jet}$ for $Q_{n+1} > Q_{\text{cut}}$

illustration: p_{\perp}^H in MEPS@NLO



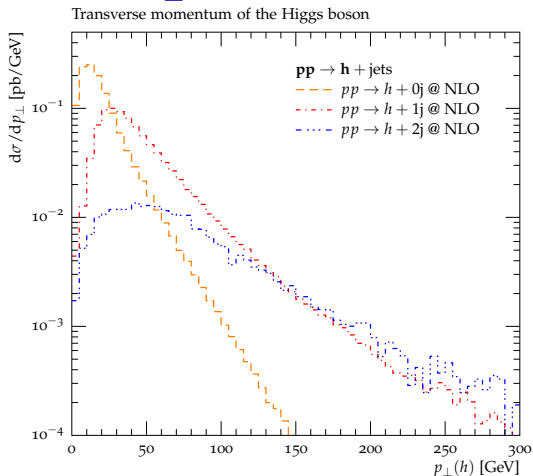
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illustration: p_{\perp}^H in MEPS@NLO



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- MC@NLO $pp \rightarrow h + 2\text{jets}$ for $Q_{n+2} > Q_{\text{cut}}$

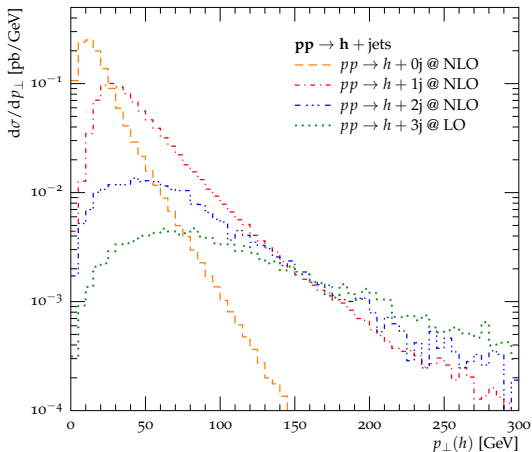
illustration: p_{\perp}^H in MEPS@NLO



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- iterate

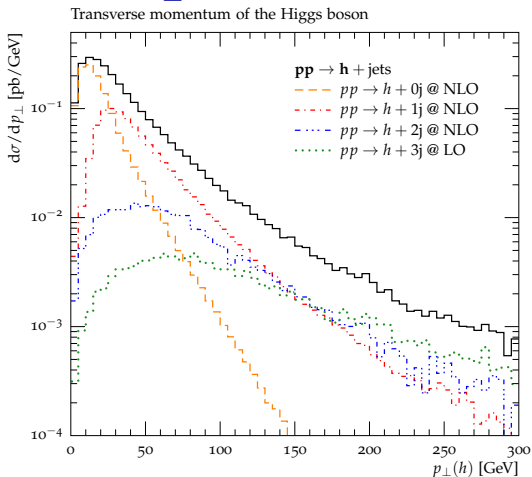
illustration: p_{\perp}^H in MEPS@NLO

Transverse momentum of the Higgs boson



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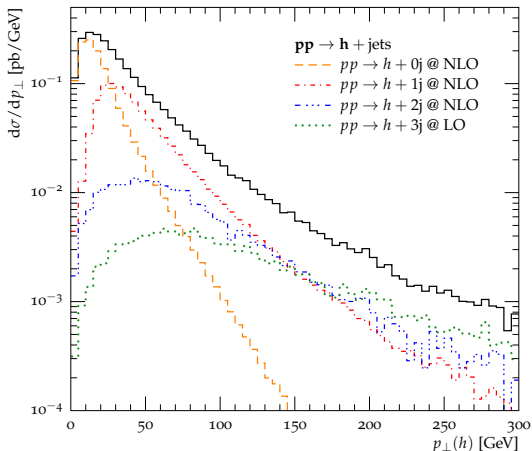
illustration: p_{\perp}^H in MEPS@NLO



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- sum all contributions

illustration: p_{\perp}^H in MEPS@NLO

Transverse momentum of the Higgs boson



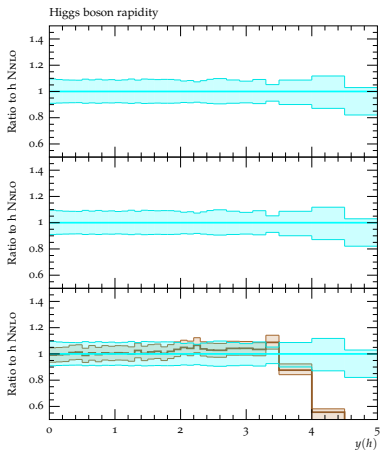
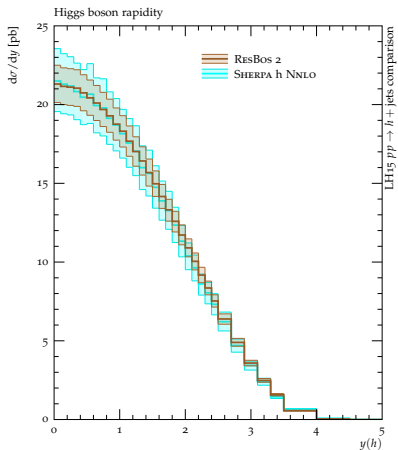
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- MC@NLO $pp \rightarrow h + 2\text{jets}$ for $Q_{n+2} > Q_{\text{cut}}$
- iterate
- sum all contributions
- eg. $p_{\perp}(h) > 200$ GeV has contributions fr. multiple topologies

detailed comparison of approaches

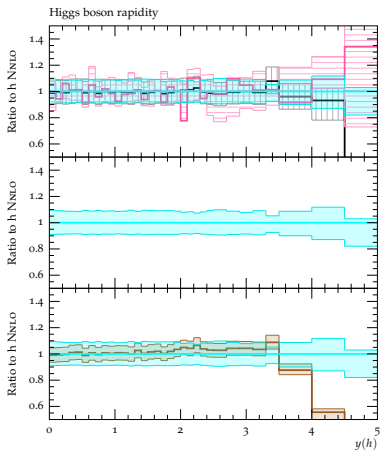
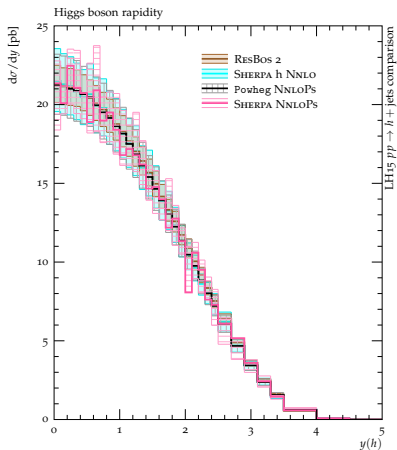
in

H +jets

inclusive Higgs boson rapidity

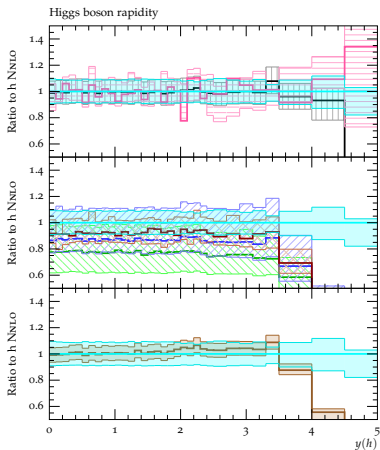
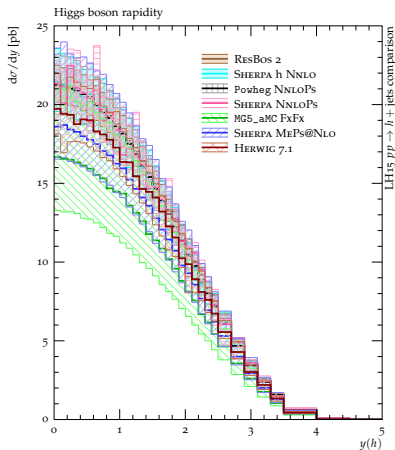


inclusive Higgs boson rapidity



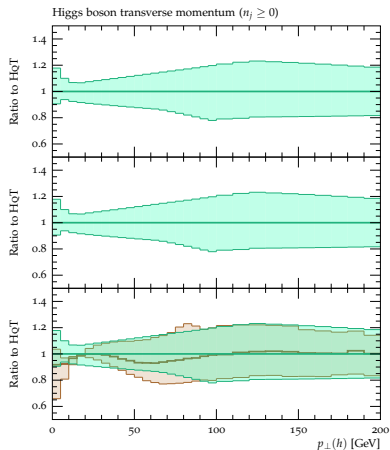
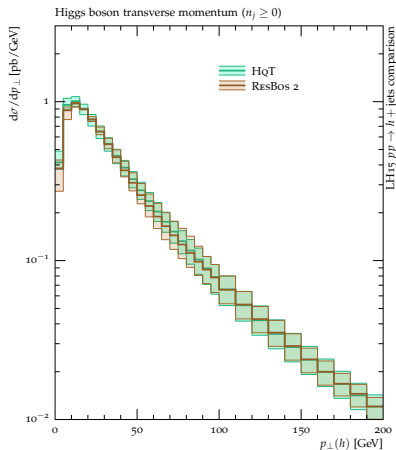
- excellent agreement between NNLO and NNLOPs

inclusive Higgs boson rapidity

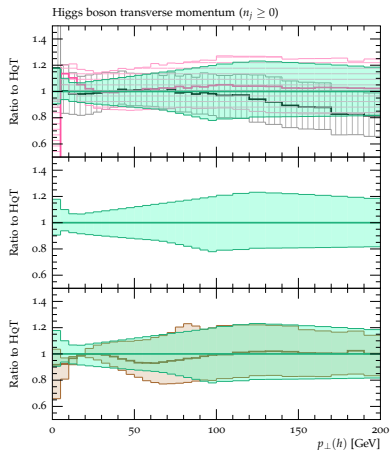
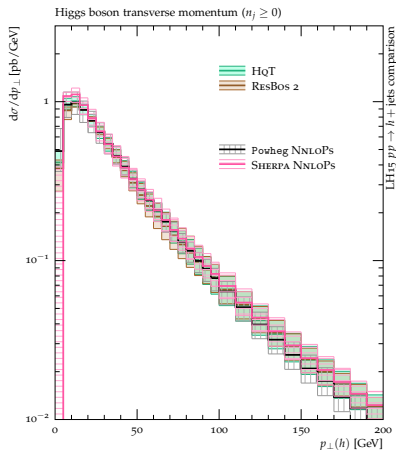


- excellent agreement between NNLO and NNLOPs
- multijet merged with NLO normalisation, PDF effects

Higgs boson transverse momentum ($n_j \geq 0$)

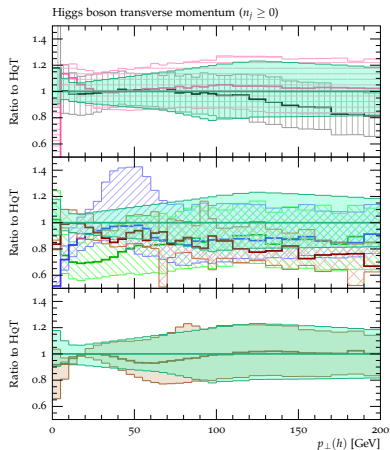
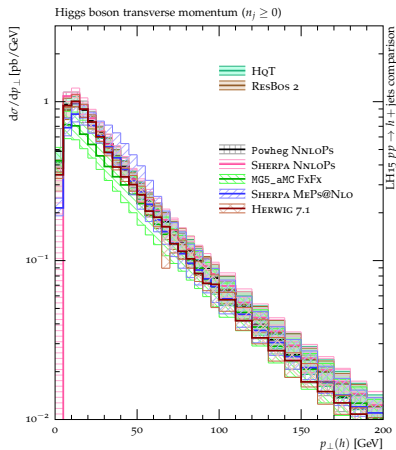


Higgs boson transverse momentum ($n_j \geq 0$)



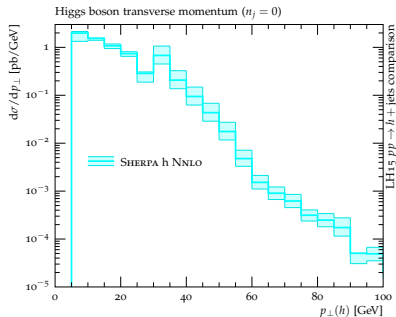
- good agreement between HqT and NNLOs, scale choice at high p_{\perp}

Higgs boson transverse momentum ($n_j \geq 0$)



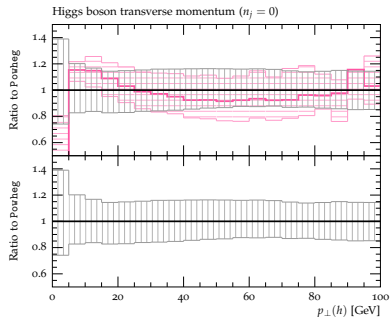
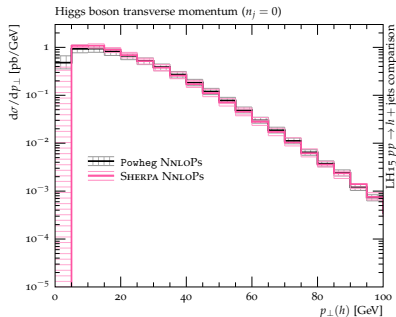
- good agreement between HqT and NNLOs, scale choice at high p_{\perp}
- multijet merged with NLO normalisation, very different at low p_{\perp}

Higgs boson transverse momentum ($n_j = 0$)



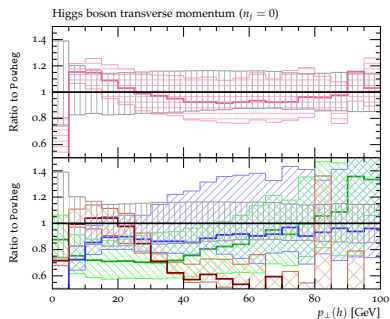
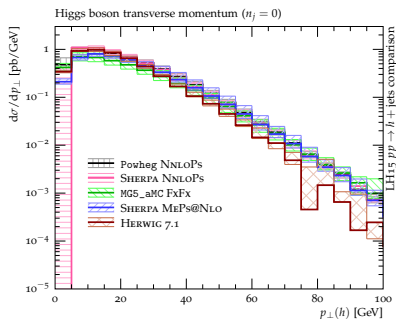
- fixed-order has various unphysical features

Higgs boson transverse momentum ($n_j = 0$)



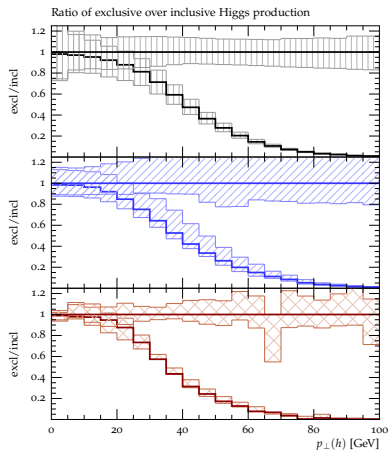
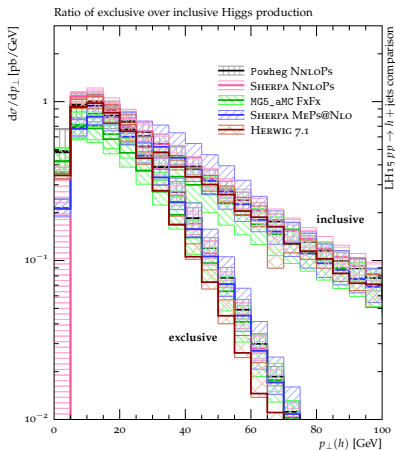
- fixed-order has various unphysical features
- good agreement between the NNLOPs

Higgs boson transverse momentum ($n_j = 0$)



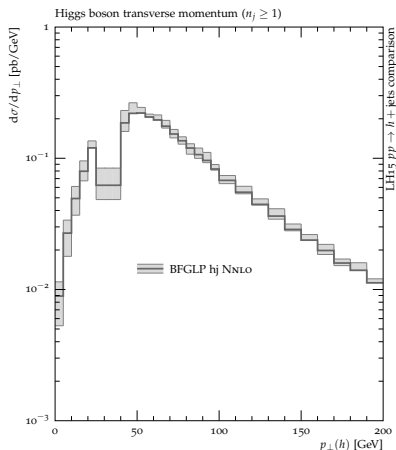
- fixed-order has various unphysical features
- good agreement between the NNLOPs
- multijet merged with NLO normalisation, HERWIG7 has much less soft radiation

exclusive over inclusive rate



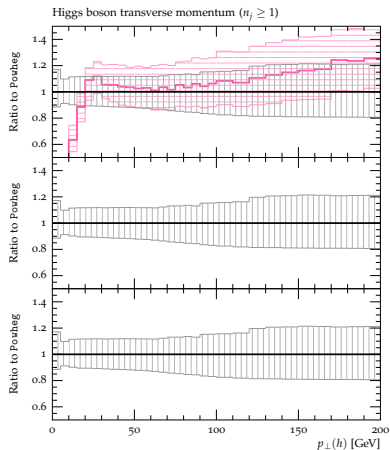
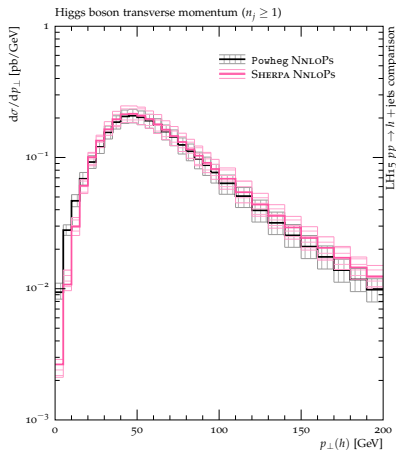
- $\approx 20\%$ of Higgs with $p_{\perp} = 60$ GeV are not accompanied by a jet

Higgs boson transverse momentum ($n_j \geq 1$)



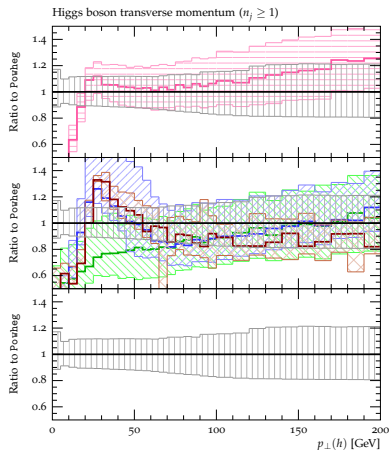
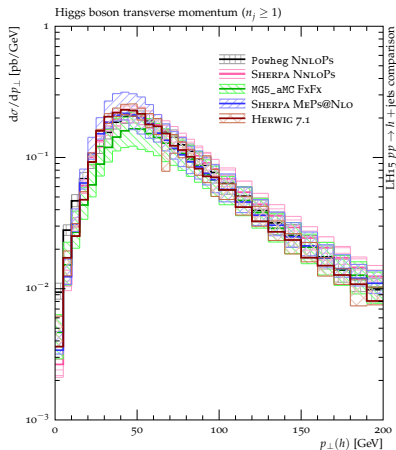
- fixed-order has Sudakov shoulder at $p_{\perp}^h = 30$ GeV due to jet cut
here: bins left and right set to average

Higgs boson transverse momentum ($n_j \geq 1$)



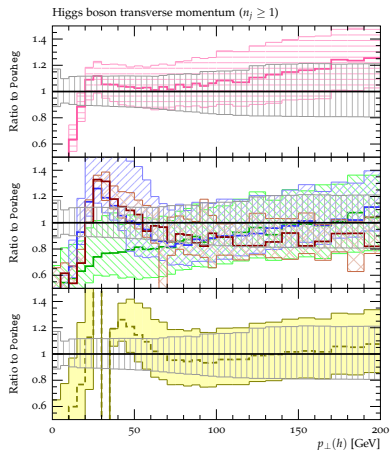
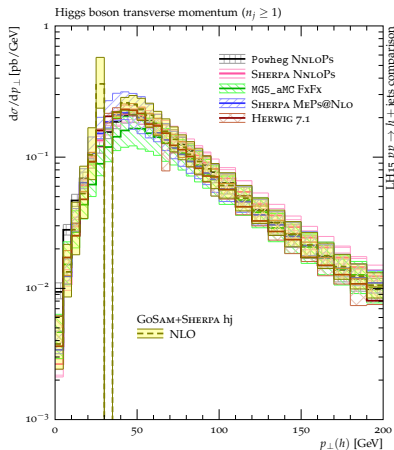
- good agreement between NNLOPs, different scales at large p_{\perp}
- excess of POWHEG as $p_{\perp} \rightarrow 0$ (Higgs strahlung off dijet)

Higgs boson transverse momentum ($n_j \geq 1$)



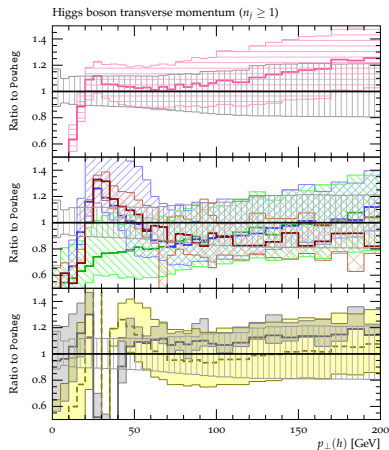
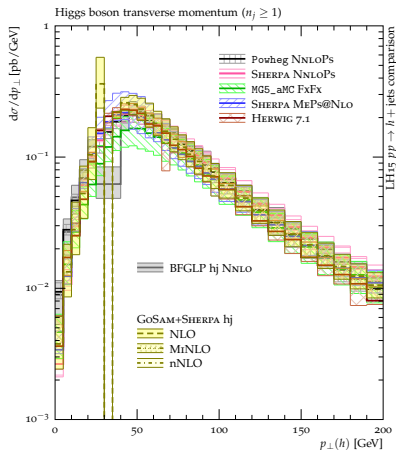
- multijet merged different shape at $p_{\perp} \lesssim 60$ GeV
- except aMC@NLO_MADGRAPH5 (due to different scales?)

Higgs boson transverse momentum ($n_j \geq 1$)



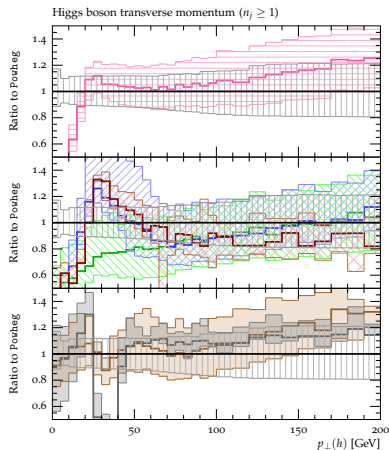
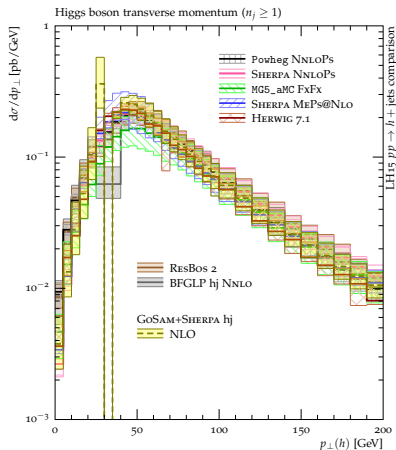
- multijet merged different shape at $p_{\perp} \lesssim 60$ GeV
- same as at fixed-order NLO

Higgs boson transverse momentum ($n_j \geq 1$)



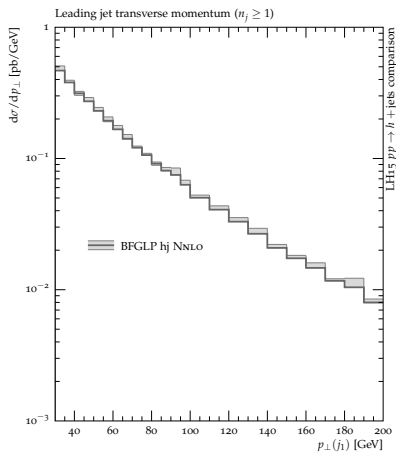
- NNLO impacts on shape at $p_{\perp} \lesssim 60$ GeV
- NLL+NLO resummation gets close to NNLO result

Higgs boson transverse momentum ($n_j \geq 1$)

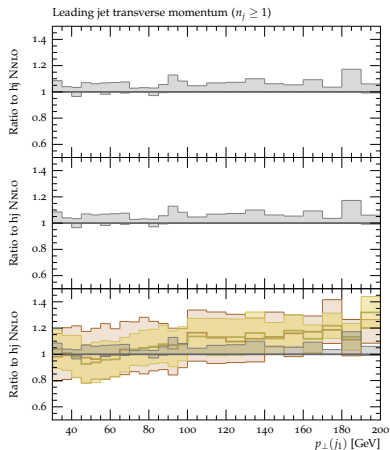
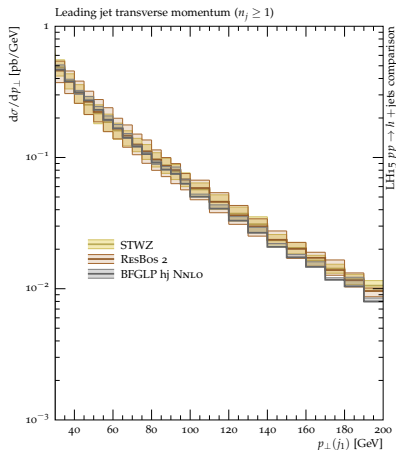


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leading jet transverse momentum ($n_j \geq 1$)

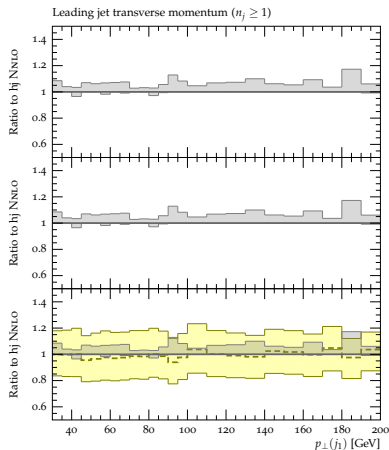
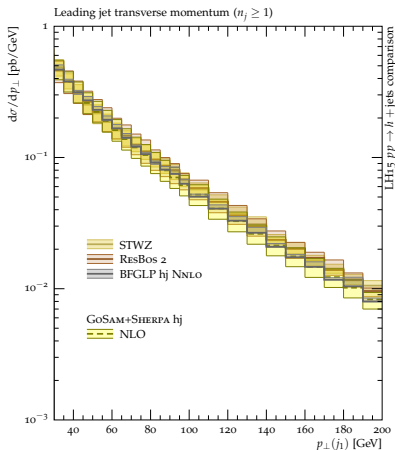


leading jet transverse momentum ($n_j \geq 1$)



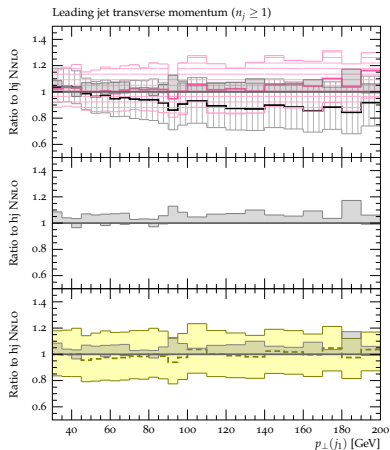
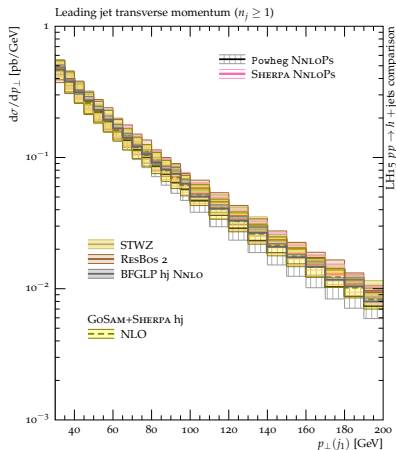
- NNLO and NLO show very good convergence for this scale choice

leading jet transverse momentum ($n_j \geq 1$)



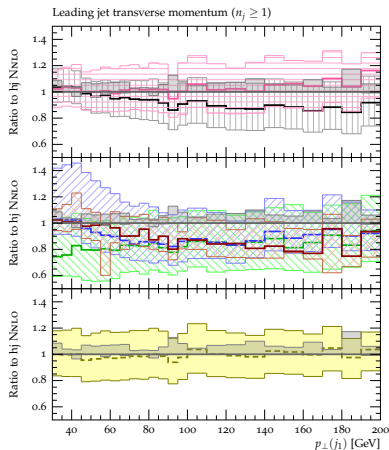
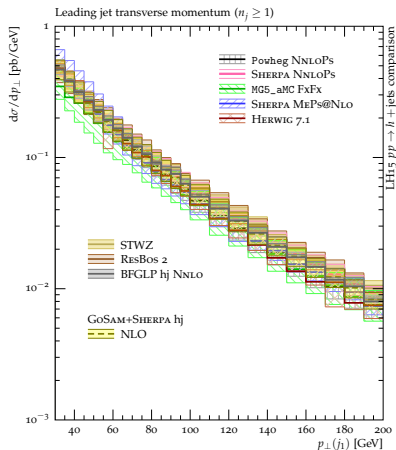
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leading jet transverse momentum ($n_j \geq 1$)



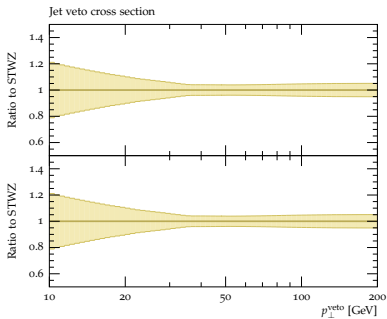
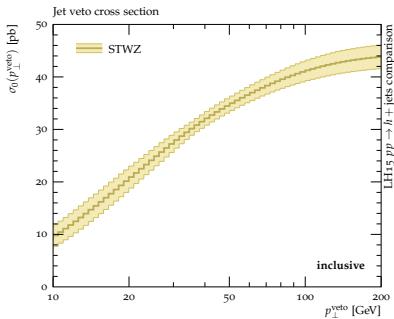
- NNLO and NLO show very good convergence for this scale choice
- multijet merged $\approx 20\%$ lower in high- p_{\perp} (due to showering)

leading jet transverse momentum ($n_j \geq 1$)

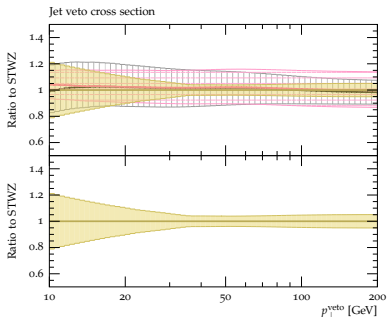
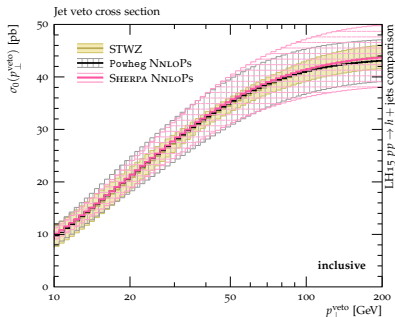


- NNLO and NLO show very good convergence for this scale choice
- multijet merged $\approx 20\%$ lower in high- p_{\perp} (due to showering)

jet vetoed cross sections – inclusive

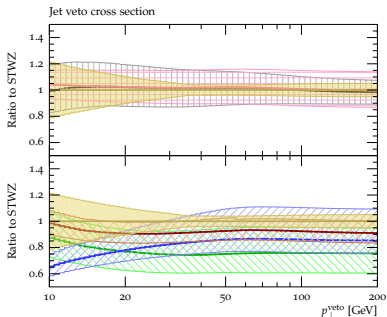
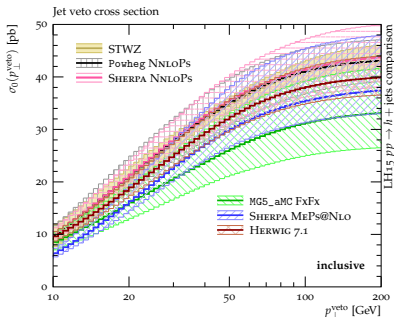


jet vetoed cross sections – inclusive



- very good agreement between NNLOPs and STWZ

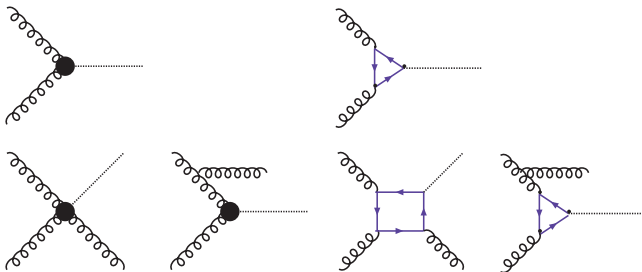
jet vetoed cross sections – inclusive



- very good agreement between NNLOPs and STWZ
- multijet merged with larger spread in shape, but within uncertainties once NLO normalisation accounted for
- PS resummation uncertainties nowhere fully assessed

aside: quark mass effects in GGF

- include effects of quark masses

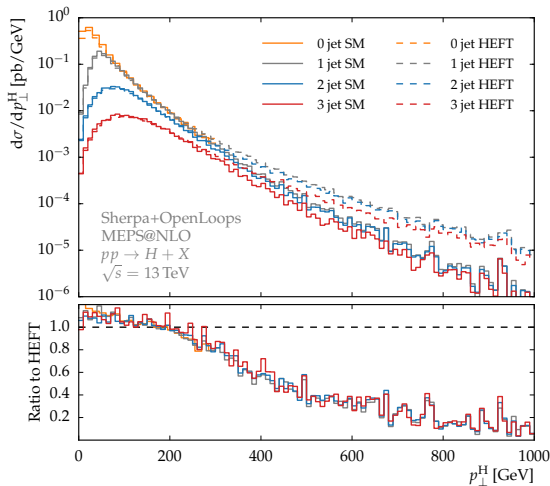
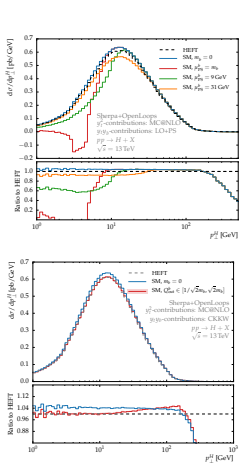


- reweight NLO HEFT with LO ratio:

(reweight virtual with Born ratio, real with real ratio)

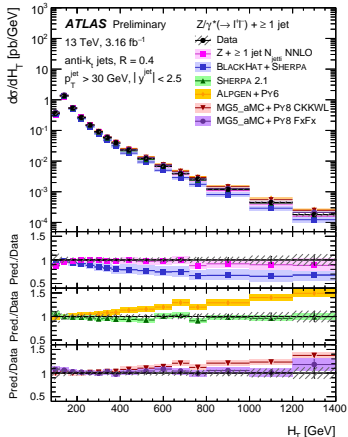
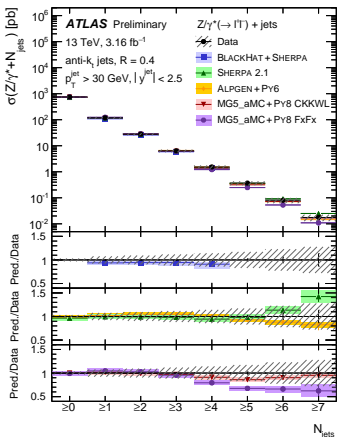
$$d\sigma_{\text{mass}}^{(\text{NLO})} \approx d\sigma_{\text{HEFT}}^{(\text{NLO})} \times \frac{d\sigma_{\text{mass}}^{(\text{LO})}}{d\sigma_{\text{HEFT}}^{(\text{LO})}}$$

- example: mass effects in $gg \rightarrow H$ (LO merging for b contribution)



Z+jets at 13 TeV: comparison with ATLAS data

- various merging codes at LO and NLO



including EW corrections

EW corrections

- EW corrections sizeable $\mathcal{O}(10\%)$ at large scales: **must include them!**
- but: more painful to calculate
- need EW showering & possibly corresponding PDFs

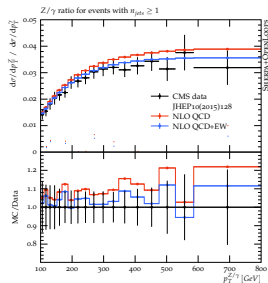
(somewhat in its infancy: chiral couplings)

- example: $Z\gamma$ vs. p_T (right plot)

(handle on p_{\perp}^Z in $Z \rightarrow \nu\bar{\nu}$)

(Kallweit, Lindert, Pozzorini, Schoenherr for LH'15)

- difference due to EW charge of Z
- no real correction (real V emission)
- improved description of $Z \rightarrow \ell\ell$



inclusion of electroweak corrections in simulation

- incorporate approximate electroweak corrections in MEPS@NLO
 - 1 using electroweak Sudakov factors

$$\tilde{B}_n(\Phi_n) \approx \tilde{B}_n(\Phi_n) \Delta_{EW}(\Phi_n)$$

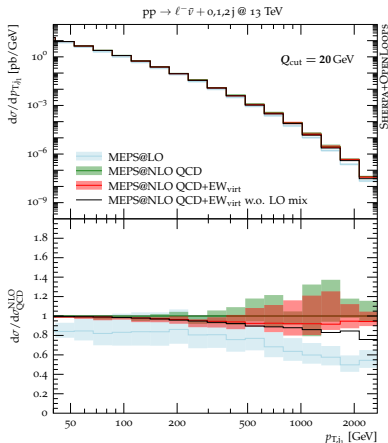
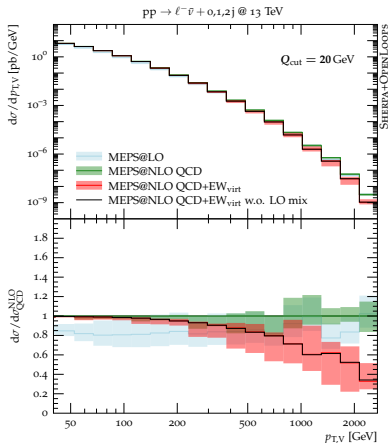
- 2 using virtual corrections and approx. integrated real corrections

$$\tilde{B}_n(\Phi_n) \approx \tilde{B}_n(\Phi_n) + V_{n,EW}(\Phi_n) + I_{n,EW}(\Phi_n) + B_{n,mix}(\Phi_n)$$

- real QED radiation can be recovered through standard tools (parton shower, YFS resummation)
- simple stand-in for proper QCD \oplus EW matching and merging
→ validated at fixed order, found to be reliable,
difference $\lesssim 5\%$ for observables not driven by real radiation

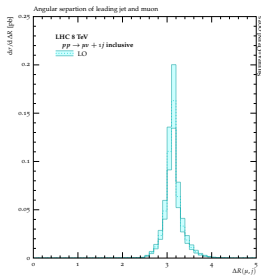
results: $pp \rightarrow \ell^- \bar{\nu} + \text{jets}$

(Kallweit, Lindert, Maierhöfer, Pozzorini, Schoenherr JHEP04(2016)021)



⇒ particle level events including dominant EW corrections

NLO EW predictions for $\Delta R(\mu, j_1)$

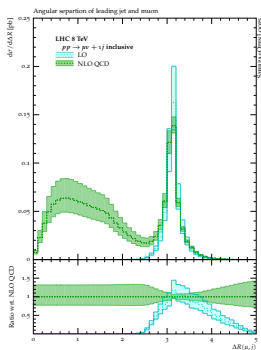


measure collinear W emission?

LHC@8TeV, $p_{\perp}^{j_1} > 500$ GeV, central μ and jet

- LO $pp \rightarrow Wj$ with $\Delta\phi(\mu, j) \approx \pi$

NLO EW predictions for $\Delta R(\mu, j_1)$

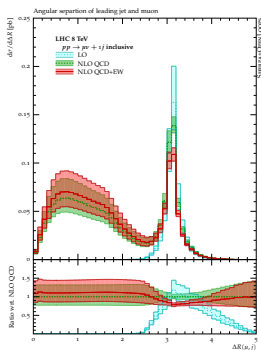


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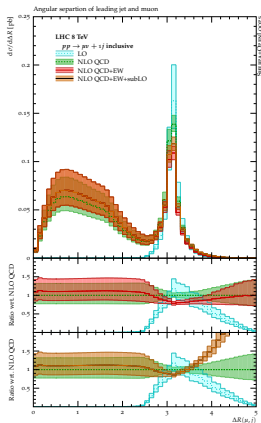
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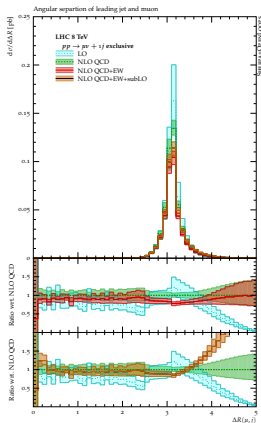


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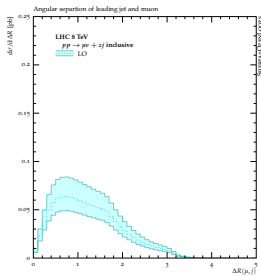
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NLO EW predictions for $\Delta R(\mu, j_1)$

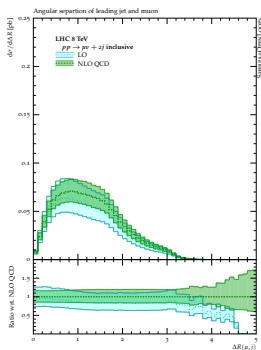


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NLO EW predictions for $\Delta R(\mu, j_1)$

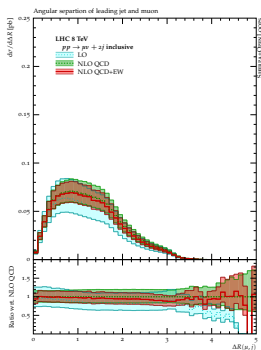


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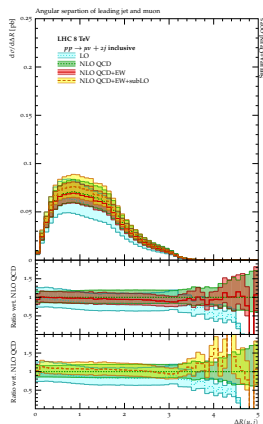


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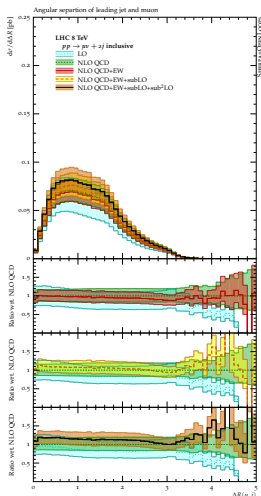
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- sub-leading Born contribs positive

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→ possible double counting with BG

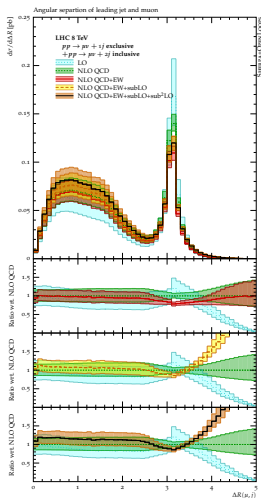


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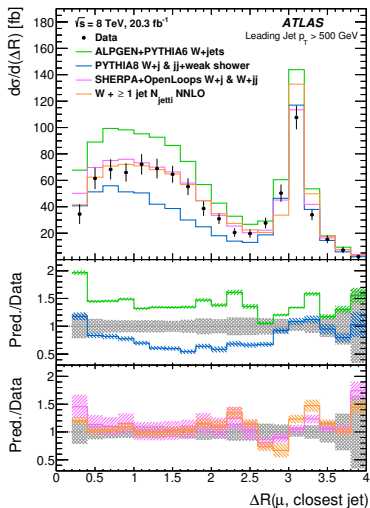
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- pos. NLO QCD, neg. NLO EW, \sim flat
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→ possible double counting with BG
- merge using exclusive sums



NLO EW predictions for $\Delta R(\mu, j_1)$

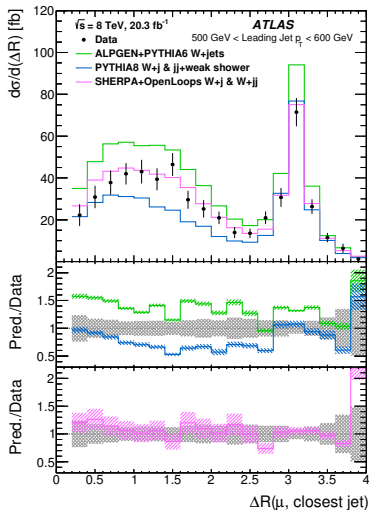


Data comparison

(M. Wu ICHEP'16, ATLAS arXiv:1609.07045)

- ALPGEN+PYTHIA
 $pp \rightarrow W + \text{jets}$ MLM merged
 (Mangano et al., JHEP07(2003)001)
- PYTHIA 8
 $pp \rightarrow Wj + \text{QCD shower}$
 $pp \rightarrow jj + \text{QCD+EW shower}$
 (Christiansen, Prestel, EPJC76(2016)39)
- SHERPA+OPENLOOPS
 NLO QCD+EW+subLO
 $pp \rightarrow Wj/Wjj$ excl. sum
 (Kallweit, Lindert, Maierhöfer,
 (Pozzorini, Schoenherr, JHEP04(2016)021)
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NLO EW predictions for $\Delta R(\mu, j_1)$

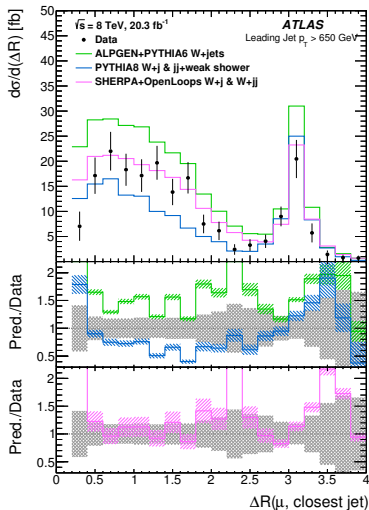


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improving parton showers

another systematic uncertainty

- parton showers are approximations, based on leading colour, leading logarithmic accuracy, spin-averaged
- parametric accuracy by comparing Sudakov form factors:

$$\Delta = \exp \left\{ - \int \frac{dk_{\perp}^2}{k_{\perp}^2} \left[A \log \frac{k_{\perp}^2}{Q^2} + B \right] \right\},$$

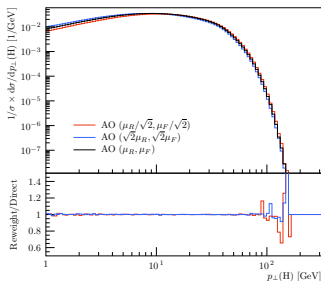
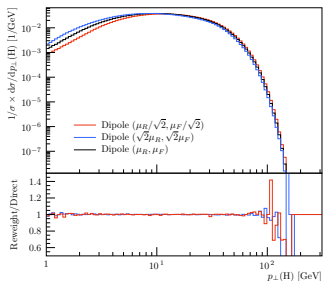
where A and B can be expanded in $\alpha_S(k_{\perp}^2)$

- showers usually include terms $A_{1,2}$ and B_1 (NLL)
- A_2 often realised by pre-factor multiplying scale $\mu_R \simeq k_{\perp}$
(CMW rescaling: Catani, Marchesini, Webber, Nucl Phys B,349 635)
- fixed-order precision necessitates to consistently assess uncertainties from parton showers
(quite often just used as black box)
- maybe improve by including higher orders?

event generation (on-the-fly scale variations)

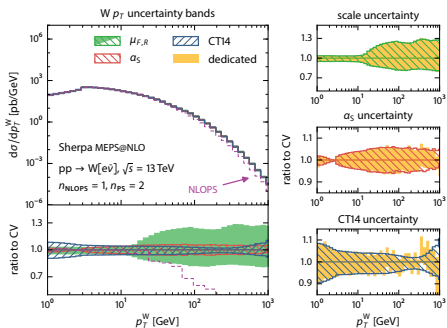
- basic idea: want to vary scales to assess uncertainties
- simple reweighting in matrix elements straightforward
- reweighting in parton shower more cumbersome
 - shower is probabilistic, concept of weight somewhat alien
 - introduce relative weight
 - evaluate (trial-)emission by (trial-)emission

implementation in HERWIG7

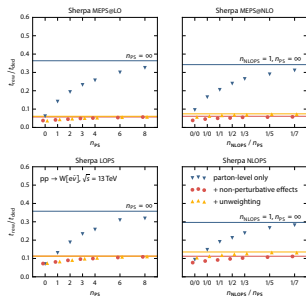


weight variation for $W+\text{jets}$ with MEPS@NLO

- uncertainties in p_{\perp}^W



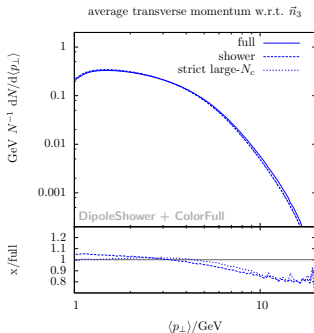
- CPU budget



going beyond leading colour

- start including next-to leading colour

(first attempts by Platzer & Sjodahl; Nagy & Soper)



- also included in 1st emission in SHERPA's MC@NLO

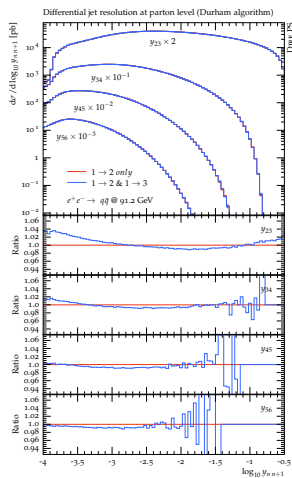
towards higher logarithmic accuracy

- reproduce DGLAP evolution at NLO
include all NLO splitting kernels
- corrections to standard $1 \rightarrow 2$ trivial
 - 2-loop cusp term subtracted & combined with LO soft contribution
 - use weighting algorithms

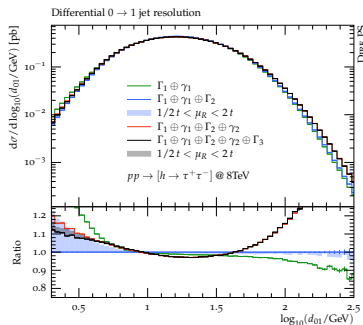
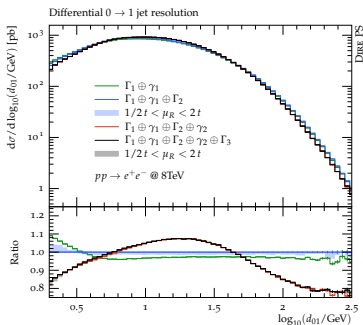
(Hoeche, Schumann, Siegert, 0912.3501)

- new topology at NLO from $q \rightarrow \bar{q}$ and $q \rightarrow q'$ splittings
- generic $1 \rightarrow 3$ process in parton shower
- first branching treated as soft gluon radiation, second as collinear splitting (to match diagrammatic structure)
- implementation complete and cross-checked (PYTHIA vs. SHERPA)

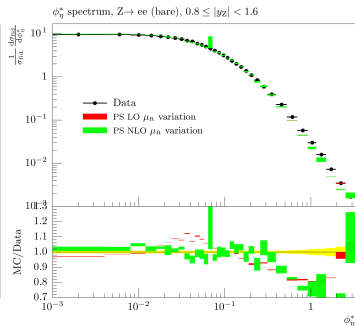
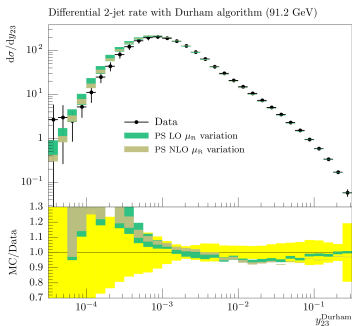
(Catani, Hoeche, FK, Prestel, in prep.)



- some first results:
DY ($pp \rightarrow e^+e^-$, left) ggF ($pp \rightarrow H \rightarrow \tau^+\tau^-$, right)
- added scale uncertainties in parton shower by varying μ_R



- some first results - comparison with data:
 y_{23} at LEP (left) and ϕ^* distribution in Drell-Yan production (right)



hadronisation

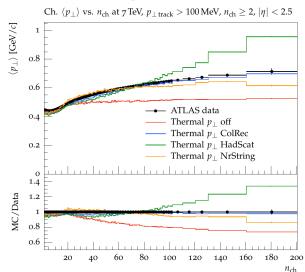
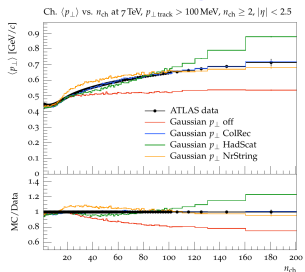
colour reconnections and friends

(Fischer, Sjostrand, 1610:09818)

Collective flow observed in pp at LHC. Partly unexpected.
New mechanisms required; could also (partly) replace CR.

Active field, e.g. N. Fischer & TS, arXiv:1610:09818 [hep-ph]:

- Thermal $\exp(-p_{\perp}/T) \rightarrow \exp(-m_{\perp}/T)$ hadronic spectrum.
- Close-packed strings \Rightarrow increased string κ or T .
- Dense hadronic gas \Rightarrow hadronic rescattering.

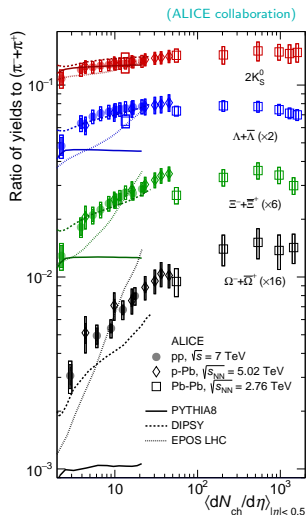


(slide stolen from Torbjorn Sjostrand)



strange strangeness

- universality of hadronisation assumed
- parameters tuned to LEP data in particular: strangeness suppression
- for strangeness: flat ratios but data do not reproduce this
- looks like $SU(3)$ restoration not observed for protons
- needs to be investigated

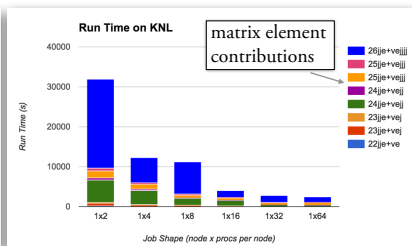
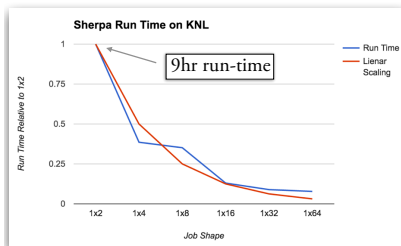


limitations
and
future challenges

limitation: computing short-distance cross sections – LO

(Childers, Uram, LeCompte, Benjamin, Hoeche, CHEP 2016)

- challenge of efficiency on tomorrow's (& today's) computers
- 2000's paradigm: memory free, flops expensive
(example: 16-core Xeon, 20MB L2 Cache, 64GB RAM)
- 2020's paradigm: flops free, memory expensive & must be managed
(example: 68-core Xeon KNL, 34MB L2 Cache, 16GB HBM, 96GB RAM)
- may **trigger rewrites of code to account for changing paradigm**



(figures stolen from Taylor Childers' talk at CHEP)



theory limitations/questions

- we have constructed lots of tools for precision physics at LHC
 - **but** we did not cross-validate them careful enough (yet)
 - **but** we did not compare their theoretical foundations (yet)
- we also need unglamorous improvements on existing tools:
 - account for new computer architectures and HPC paradigms
 - systematically check advanced scale-setting schemes (MINLO)
 - automatic (re-)weighting for PDFs & scales
 - scale compensation in PS is simple (implement and check)
- 4 vs. 5 flavour scheme → **really?**
- how about α_S : range from 0.113 to 0.118

(yes, I know, but still - it bugs me)

→ is there any way to settle this once and for all (measurements?)

achievable goals (I believe we know how to do this)

- NLO for loop-induced processes:
 - fixed-order starting, MC@NLO tedious but straightforward
- EW NLO corrections with tricky/time-consuming calculation setup
 - but important at large scales: effect often \sim QCD, but opposite sign
 - need maybe faster approximation for high-scales (EW Sudakovs)
 - work out full matching/merging instead of approximations
- improve parton shower:
 - beyond (next-to) leading log, leading colour, spin-averaged
 - HO effects in shower and scale uncertainties
 - start including next-to leading colour
 - include spin-correlations \rightarrow important for EW emissions

more theory uncertainties/issues?

- with NNLOPS approaching 5% accuracy or better:
 - non-perturbative uncertainties start to matter:
 - PDFs, MPIs, hadronization, etc.
 - question (example): with hadronization tuned to quark jets (LEP)
 - how important is the “chemistry” of jets for JES?
 - can we fix this with measurements?
 - example PDFs: to date based on FO vs. data
 - will we have to move to resummed/parton showered?
- (reminder: LO* was not a big hit, though)
- $g \rightarrow q\bar{q}$ at accuracy limit of current parton showers:
 - how bad are $\sim 25\%$ uncertainty on $g \rightarrow b\bar{b}$?
 - can we fix this with measurements?

the looming revolution: going beyond NLO

- H in ggF at N³LO (Anastasiou, Duhr and others)
- explosive growth in NNLO (QCD) $2 \rightarrow 2$ results

(apologies for any unintended omissions)

- $t\bar{t}$ (1303.6254; 1508.03585; 1511.00549)
 - single- t (1404.7116)
 - VV (1507.06257; 1605.02716; 1604.08576; 1605.02716)
 - HH (1606.09519)
 - VH (1407.4747; 1601.00658; 1605.08011)
 - $V\gamma$ (1504.01330)
 - $\gamma\gamma$ (1110.2375; 1603.02663)
 - Vj (1507.02850; 1512.01291; 1602.06965; 1605.04295; 1610.07922)
 - Hj (1408.5325; 1504.07922; 1505.03893; 1508.02684; 1607.08817)
 - jj (1310.3993; 1611.01460)
- NLO corrections to $gg \rightarrow VV$ (1605.04610)
 - WBF at NNLO (1506.02660) and N³LO (1606.00840)
 - different IR subtraction schemes:
N-jettiness slicing, antenna subtraction, sector decomposition,

living with the revolution

- we will include them into full simulations

(I am willing to place a bet: 5 years at most!)

- practical limitations/questions to be overcome:

- dealing with IR divergences at NNLO: slicing vs. subtracting

(I'm not sure we have THE solution yet)

- how far can we push NNLO? are NLO automated results stable enough for NNLO at higher multiplicity?
- matching for generic processes at NNLO?

(MINLO or UN²LOs or something new?)

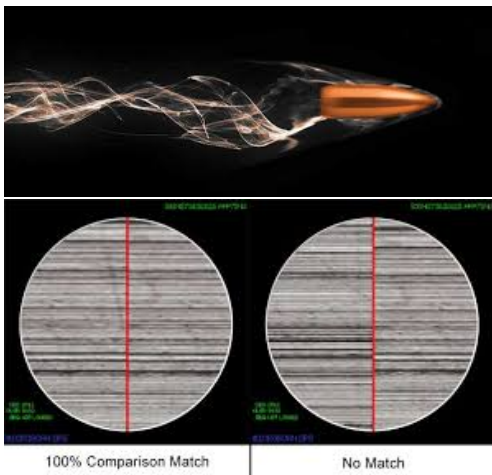
- more scales (internal or external) complicated – need integrals

- philosophical questions:

- going to higher power of N often driven by need to include larger FS multiplicity – maybe not the most efficient method
- limitations of perturbative expansion:
 - breakdown of factorisation at HO (Seymour et al.)
 - higher-twist: compare $(\alpha_S/\pi)^n$ with Λ_{QCD}/M_Z

outlook

- will need precision for ballistics of smoking guns



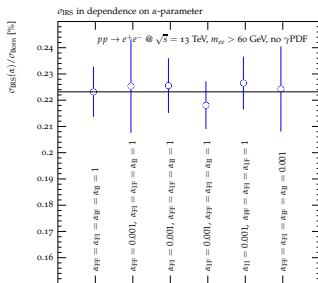
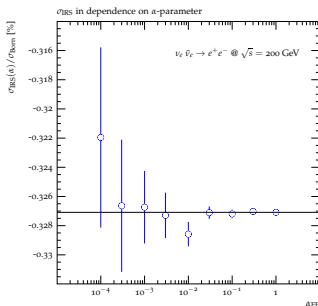
extra: NLO EW subtraction in SHERPA

(M. Schoenherr in preparation)

- adapt QCD subtraction (spl. fns. and colour-/spin-correlated MEs)

(Catani, Dittmaier, Seymour, Trocsanyi Nucl.Phys.B627(2002)189-265)

- replacements: $\alpha_s \rightarrow \alpha$, $C_F \rightarrow Q_f^2$, $C_A \rightarrow 0$,
 $T_R \rightarrow N_{c,f} Q_f^2$, $n_f T_R \rightarrow \sum_f N_{c,f} Q_f^2$,
 $\frac{\mathbf{T}_{ij} \cdot \mathbf{T}_k}{\mathbf{T}_{ij}^2} \rightarrow \frac{Q_{ij} Q_k}{Q_{ij}^2}$



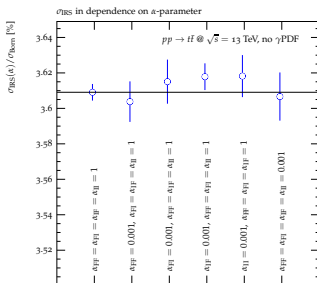
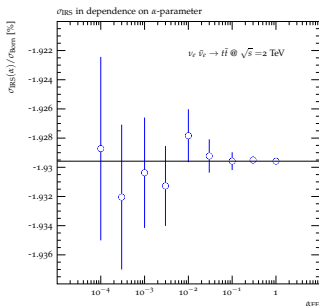
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