



## Exercise 1

1. (a) Convince yourself that the definition of the Lagrange density

$$\mathcal{L} = -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} - j_\kappa A^\kappa$$

is equivalent to

$$\mathcal{L} = \frac{1}{2}(\vec{E}^2 - \vec{B}^2) - \rho\phi + \vec{j} \cdot \vec{A}.$$

Here  $A^\mu = (\phi, \vec{A})$  and  $j^\mu = (\rho, \vec{j})$  denote the contravariant vector potential and current, respectively. The field strength tensor is given by  $F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu$ . Derive the equations of motion of the electromagnetic fields  $\vec{E}$  and  $\vec{B}$ .

- (b) The dual field strength tensor can be obtained through

$$\tilde{F}^{\mu\nu} = \frac{1}{2}\epsilon^{\mu\nu\rho\sigma}F_{\rho\sigma},$$

with  $\epsilon^{\mu\nu\rho\sigma}$  the Levi-Civita tensor. Prove that the homogeneous Maxwell equations can be written as  $\partial_\mu \tilde{F}^{\mu\nu} = 0$  and that  $F_{\mu\nu}\tilde{F}^{\mu\nu} = -4\vec{E} \cdot \vec{B}$ .

2. (a) Consider the Hamilton operator

$$\hat{H}_F = i\omega\hat{\psi}\hat{\pi},$$

where the Hermitian operators  $\hat{\psi}$  and  $\hat{\pi}$  fulfill the anti-commutation relations

$$\{\hat{\psi}, \hat{\pi}\} = 0, \quad \{\hat{\psi}, \hat{\psi}\} = \hat{1}, \quad \{\hat{\pi}, \hat{\pi}\} = \hat{1}.$$

In analogy to the bosonic harmonic oscillator, introduce the creation and annihilation operators

$$\hat{f}^\dagger = \frac{1}{\sqrt{2}}(\hat{\psi} - i\hat{\pi}) \quad \text{and} \quad \hat{f} = \frac{1}{\sqrt{2}}(\hat{\psi} + i\hat{\pi})$$

and solve the eigenvalue problem of  $\hat{H}_F$ .

- (b) Determine the spectrum of the supersymmetric Hamilton operator

$$\hat{H}_{\text{SUSY}} = \{\hat{Q}^\dagger, \hat{Q}\},$$

where  $\hat{Q}^\dagger = \sqrt{\omega}\hat{b}\hat{f}^\dagger$  and  $\hat{Q} = \sqrt{\omega}\hat{b}^\dagger\hat{f}$  and  $\hat{b}^{(\dagger)}$  the creation and annihilation operators of the bosonic oscillator.

(turn over)

3. The (anti-)symmetrization-operator is defined by its action on an  $n$ -particle wave function

$$\begin{aligned}\hat{\mathcal{S}}\Psi(x_1, \dots, x_N) &\equiv \frac{1}{N!} \sum_{\sigma \in \mathcal{P}_N} \hat{P}_\sigma \Psi(x_1, \dots, x_N), \\ \hat{\mathcal{A}}\Psi(x_1, \dots, x_N) &\equiv \frac{1}{N!} \sum_{\sigma \in \mathcal{P}_N} \text{sign}(\sigma) \hat{P}_\sigma \Psi(x_1, \dots, x_N),\end{aligned}$$

where  $\sigma$  is an element of the permutation group  $\mathcal{P}_N$  and

$$\hat{P}_\sigma \Psi(x_1, \dots, x_N) \equiv \Psi(x_{\sigma(1)}, \dots, x_{\sigma(N)}).$$

- (a) Prove that the relations  $\hat{\mathcal{S}}^2 = \hat{\mathcal{S}}$ ,  $\hat{\mathcal{A}}^2 = \hat{\mathcal{A}}$  and  $\hat{\mathcal{S}}\hat{\mathcal{A}} = \hat{\mathcal{A}}\hat{\mathcal{S}} = 0$  hold true.
- (b) Prove the **symmetry postulate**, stating that the Hilbert space of a system of identical particles contains either symmetric or anti-symmetric states only.
4. Compute the integrals

$$I_{ee} = \frac{e^2}{2} \frac{1}{V^2} \int_V \int_V d^3r d^3r' \frac{e^{-i(\vec{k}_1 - \vec{k}_3) \cdot \vec{r}} e^{-i(\vec{k}_2 - \vec{k}_4) \cdot \vec{r}'}}{|\vec{r} - \vec{r}'|} e^{-\alpha|\vec{r} - \vec{r}'|},$$

and

$$I_I = \frac{e^2}{2} \frac{N^2}{V^2} \int_V \int_V d^3R d^3R' \frac{e^{-\alpha|\vec{R} - \vec{R}'|}}{|\vec{R} - \vec{R}'|},$$

where  $\alpha > 0$  and

$$\delta_{\vec{k}_1, \vec{k}_2} \equiv \frac{1}{V} \int_V d^3r e^{-i(\vec{k}_1 - \vec{k}_2) \cdot \vec{r}}.$$