

a_μ and its implications for physics beyond the SM, particularly in view of the questions of grand unification, origin of flavour and CP violation!

G.Ross, Glasgow, October 2007

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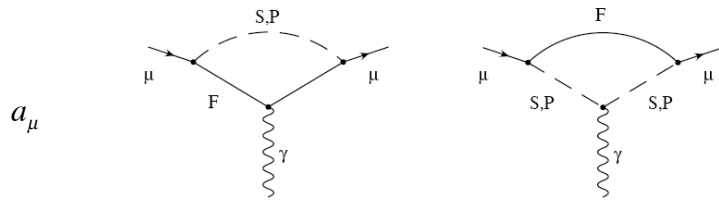
- Supersymmetry
- Extra dimensions
- Additional gauge bosons
- Additional (vectorlike) fermions
- Additional scalars
- Exotic flavour-changing interactions
- Nonperturbative effects
- Muon substructure
- Anomalous gauge boson couplings
-

a_μ and its implications for physics beyond the SM, particularly in view of the questions of grand unification, origin of flavour and CP violation!

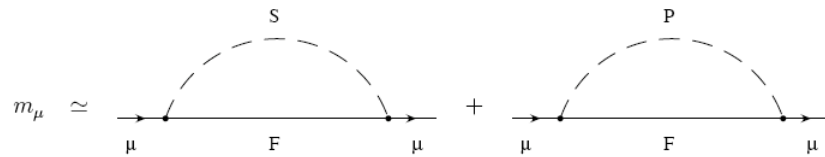
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- Supersymmetry Hierarchy problem, GUT
- Extra dimensions Hierarchy problem, GUT
- Additional gauge bosons L-R symmetry, Family group, ..?
- Additional (vectorlike) fermions Froggatt-Nielsen mass generation, string compact...
- Additional scalars L-R symmetry, Familons, little Higgs...
- Exotic flavour-changing interactions Family symmetries, technicolour...
- Nonperturbative effects Composite structure...
- Muon substructure $\times?$ $1.7\text{TeV} < \Lambda_\mu < 2.3\text{TeV}$ *c.f.* $\Lambda \geq 4 - 5\text{TeV}$
- Anomalous gauge boson couplings $\times?$
-

$$\underline{a_\mu \Leftrightarrow m_\mu}$$



$$\frac{g^2}{8\pi^2} \frac{m_\mu M_F}{M_P^2} \left(1 - \frac{M_P^2}{M_S^2} \right)$$



$$\frac{g^2}{16\pi^2} M_F \ln \left(\frac{M_S^2}{M_P^2} \right)$$

$$a_\mu(\text{New Physics}) \simeq C \frac{m_\mu^2}{M_P^2},$$

$$C = \left(1 - \frac{M_P^2}{M_S^2} \right) / \ln \frac{M_S^2}{M_P^2}$$

In general

$$a_\mu = C \frac{m_\mu^2}{M^2}$$

$$C \leq O(1)$$

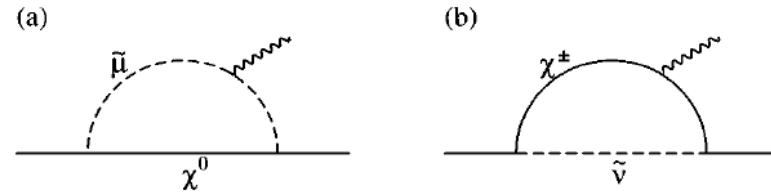
Application to SUSY

$$\Delta a_\mu \propto g_{1,2}^2 m_\mu^2 \frac{m_{G_{1,2}} \mu_H}{\widetilde{M}^2} \tan \beta$$

$$\Delta m_\mu \propto g_{1,2}^2 m_\mu \frac{m_{G_{1,2}} \mu_H}{\widetilde{M}^2} \tan \beta$$

$$a_\mu = C \frac{m_\mu^2}{\widetilde{M}^2}$$

Δa_μ large \Rightarrow Δm_μ large



Test of large Δm_μ ?

Yukawa Unification

$$\psi_\alpha = \begin{pmatrix} d \\ d \\ d \\ l \end{pmatrix}$$

$$|V_{us}| = \left| \sqrt{\frac{m_d}{m_s}} - e^{i\delta} \sqrt{\frac{m_u}{m_c}} \right|$$

$$\frac{m_s}{m_\mu}(M_X) = \frac{1}{3}$$

$$\frac{m_d}{m_e}(M_X) = 3$$

Georgi Jarlskog

$$\frac{M^{d,l}}{m_3} = \begin{pmatrix} < \epsilon^4 & \epsilon^3 \\ \epsilon^3 & a^{d,l} \epsilon^2 \\ & & 1 \end{pmatrix}$$

$\epsilon \approx 0.15$

ψ ψ_α

$$\frac{m_b}{m_\tau}(M_X) = 1$$

$$a^d = 1 \quad a^l = -3$$

$$\psi \begin{pmatrix} 1 & & & \\ & 1 & & \\ & & 1 & \\ & & & -3 \end{pmatrix} \psi_\alpha$$

$$\text{Det}(M^l) = \text{Det}(M^d) |_{M_X}$$

Effect of $\Delta m_\mu^{\text{SUSY}}$, Δm_b^{SUSY}

Parameters	Input SUSY Parameters					
$\tan \beta$	1.3	10	38	50	38	38
γ_b	0	0	0	0	-0.22	+0.22
γ_d	0	0	0	0	-0.21	+0.21
γ_t	0	0	0	0	0	-0.44
Parameters	Comparison with GUT Mass Ratios					
$(m_b/m_\tau)(M_X)$	$1.00^{+0.04}_{-0.4}$	0.73(3)	0.73(3)	0.73(4)	1.00(4)	1.00(4)
$(3m_s/m_\mu)(M_X)$	$0.70^{+0.8}_{-0.05}$	0.69(8)	0.69(8)	0.69(8)	0.9(1)	0.6(1)
$(m_d/3m_e)(M_X)$	0.82(7)	0.83(7)	0.83(7)	0.83(7)	1.05(8)	0.68(6)
$(\frac{\det Y^d}{\det Y^e})(M_X)$	$0.57^{+0.08}_{-0.26}$	0.42(7)	0.42(7)	0.42(7)	0.92(14)	0.39(7)

Serna, GGR

$$\begin{aligned} \gamma_t &\approx y_t^2 \mu A^t \frac{\tan \beta}{16\pi^2} I_3(m_{\tilde{t}_1}^2, m_{\tilde{t}_2}^2, \mu^2) \sim y_t^2 \frac{\tan \beta \mu A^t}{32\pi^2 m_{\tilde{t}}^2} \\ \gamma_u &\approx -g_2^2 M_2 \mu \frac{\tan \beta}{16\pi^2} I_3(m_{\chi_1}^2, m_{\chi_2}^2, m_u^2) \sim 0 \\ \gamma_b &\approx \frac{8}{3} g_3^2 \frac{\tan \beta}{16\pi^2} M_3 \mu I_3(m_{\tilde{b}_1}^2, m_{\tilde{b}_2}^2, M_3^2) \sim \frac{4}{3} g_3^2 \frac{\tan \beta \mu M_3}{16\pi^2 m_{\tilde{b}}^2} \\ \gamma_d &\approx \frac{8}{3} g_3^2 \frac{\tan \beta}{16\pi^2} M_3 \mu I_3(m_{\tilde{d}_1}^2, m_{\tilde{d}_2}^2, M_3^2) \sim \frac{4}{3} g_3^2 \frac{\tan \beta \mu M_3}{16\pi^2 m_{\tilde{d}}^2} \end{aligned}$$

$$\frac{\mu M_3}{m_{\tilde{b}}^2} \sim -0.5, \quad \frac{m_{\tilde{b}}^2}{m_{\tilde{d}}^2} \sim 1.0$$

$$c.f. \quad \Delta a_\mu \propto g_{1,2}^2 m_\mu^2 \frac{m_{G_{1,2}} \mu_H}{\widetilde{M}^2} \tan \beta$$

$$\Rightarrow M_3 < 0, M_{1,2} > 0 ?$$

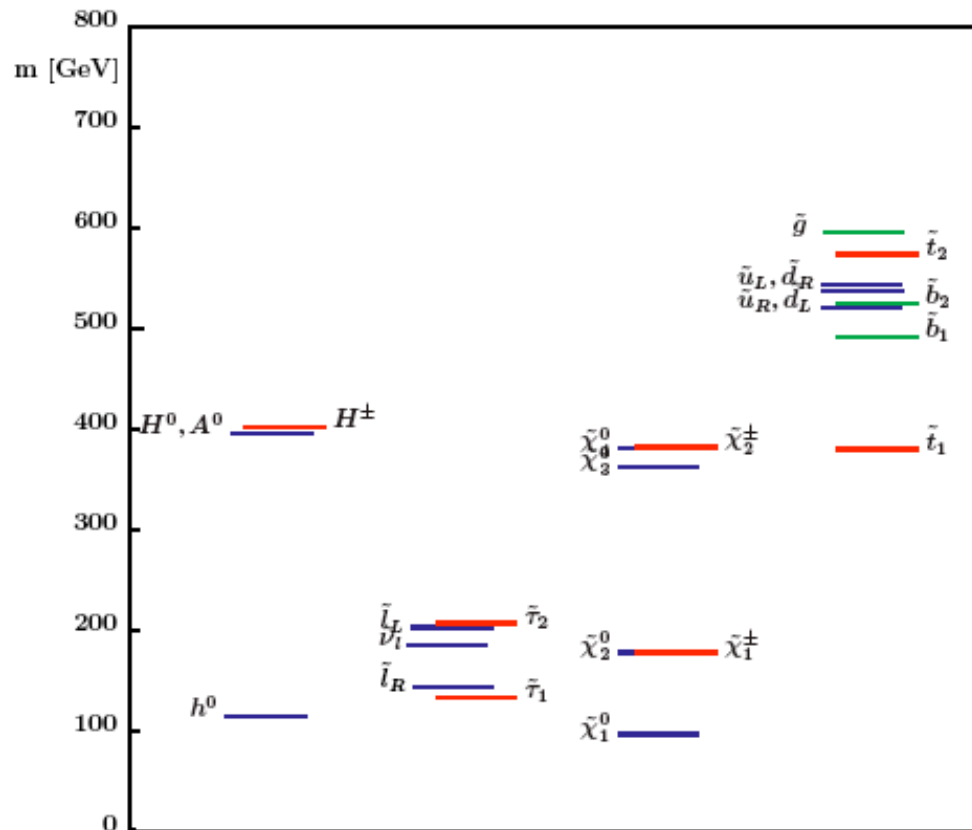
a_μ and Grand Unification

- GUT predictions for m improved
- $\widetilde{M}_{i=1,2,3}$ not universal?

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SPS 1



e.g. Anomaly mediation

RG invariant relations – fixed point

$$M_i = m_0 \beta_{g_i} / g_i \quad M_3 < 0$$

$$h_{t,b,\tau} = -m_0 \beta_{Y_{t,b,\tau}}$$

$$(m^2)^i_j = \frac{1}{2} m_0^2 \mu \frac{d}{d\mu} \gamma^i_j$$

$$m_3^2 = \kappa m_0 \mu - m_0 \beta_\mu$$

$$a_\mu^{\text{SUSY}} = 29.8$$

m_0	100 GeV
$m_{1/2}$	250 GeV
A_0	-100 GeV
$\tan \beta$	10
$\text{sign } \mu$	+

a_μ and Grand Unification

- $\widetilde{M}_{i=1,2,3}$ not universal.
- Δa_μ large \Rightarrow sneutrino and charginos relatively light
- $\Delta(b \rightarrow s\gamma)$ small \Rightarrow heavy charged Higgs and heavy stop

$$\mathcal{BR}(b \rightarrow s\gamma)|_{\chi^\pm} \propto \mu A_t \tan \beta f(m_{\tilde{t}_1}, m_{\tilde{t}_2}, m_{\tilde{\chi}^\pm}) \frac{m_b}{v(1 + \Delta m_b)}$$

Effect of \cancel{CP}

New SUSY sources of \cancel{CP} :

$$M_C = \begin{pmatrix} |\tilde{m}_2| e^{i\xi_2} & \sqrt{2} m_W \sin \beta \\ \sqrt{2} m_W \cos \beta & |\mu| e^{i\theta_\mu} \end{pmatrix}$$

To suppress neutron, electron EDM, $\xi_i, \theta_\mu = O(10^{-2})$ unless there are cancellations

For large phases significant corrections to a_μ

Table 1: Cases where the EDM and the g-2 experiments are satisfied

(case) ξ_2, θ_μ, ξ_3 (radian)	d_e, d_n (ecm)	a_μ^{SUSY}
(a) $-.63, .3, .37$	$-4.2 \times 10^{-27}, -5.3 \times 10^{-26}$	47.0×10^{-10}
(b) $-.85, .4, .37$	$4.2 \times 10^{-27}, 4.8 \times 10^{-26}$	10.8×10^{-10}
(c) $-.8, .2, 1.3$	$4.0 \times 10^{-27}, 5.4 \times 10^{-26}$	12.2×10^{-10}
(d) $-.32, .3, -.28$	$-1.2 \times 10^{-27}, 3.3 \times 10^{-26}$	20.1×10^{-10}
(e) $-.5, .49, -.5$	$1.8 \times 10^{-27}, -6.6 \times 10^{-27}$	12.7×10^{-10}

Ibrahim, Chattopadhyay, Nath

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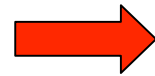
Ibrahim, Chattopadhyay, Nath

But spontaneously broken family symmetries can naturally explain why $\xi_i, \theta_\mu = O(10^{-2})$

Vives, GGR

Origin of flavour

$$\begin{pmatrix} m_e & 0 & 0 \\ 0 & m_\mu & 0 \\ 0 & 0 & m_\tau \end{pmatrix} \simeq \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & m_\tau \end{pmatrix}$$



FAMILY SYMMETRY?

Abelian, Non-Abelian $\subset (U(3))^6$

spontaneously broken:

$$\begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\xrightarrow{\langle \theta \rangle \neq 0}$$

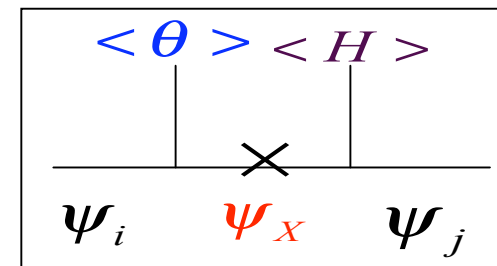
$$\begin{pmatrix} 0 & 0 & 0 \\ 0 & a\varepsilon^2 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$a = O(1), \quad \varepsilon^2 = \frac{\langle \theta \rangle}{M}$$



Messenger sector (heavy fermions, heavy Higgs)

$$m_{ij} = \frac{\langle \theta \rangle \langle H \rangle}{M_x}$$



+family gauge group or familons

- Additional gauge bosons L-R symmetry, Family group, ..?

$$a_\mu = C \frac{m_\mu^2}{M^2} \quad C \leq O(1)$$

$C \ll 1$ family diagonal couplings (m_μ small)

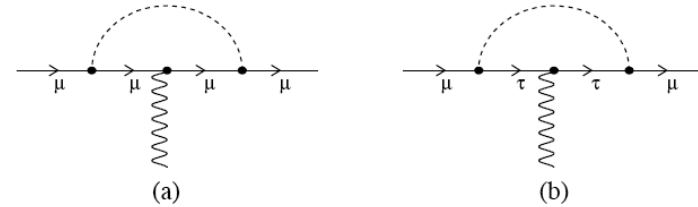
$C \sim 1$ coupling to heavy fermions... $M=1-2\text{TeV}$

Strong correlation with FCNC

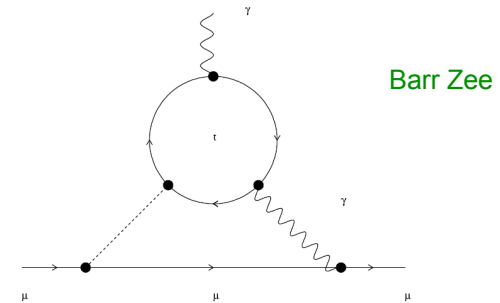
e.g. Z' with large $\tau - \mu$ mixing: $\Delta a_\mu = 221 \times 10^{-11} \Rightarrow BR(\tau \rightarrow \mu\gamma) = 1 \times 10^{-6}$

- Additional scalars Extended Higgs structure, familons...

$$\Delta a_\mu = \pm \frac{1}{8\pi^2} \frac{m_\mu^2}{m_\phi^2} \ln \left(\frac{m_\phi^2}{m_\mu^2} \right) Y_{ii}^2$$



$$\Delta a_\mu^h = \mp \frac{N_c q_t^2}{\pi^2} \frac{m_\mu m_t}{m_\phi^2} F \left(\frac{m_t^2}{m_\phi^2} \right) Y_{tt} Y_{\mu\mu}$$



For $m_A \simeq 150 GeV$ two loop term can dominate.

Significant contribution if large μ - τ mixing

$$Y_{\mu\tau} = 0.04 \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}$$

$\tau \rightarrow \mu\gamma$ close to present bound

Summary

- Many ways to generate

Δa_μ

$$a_\mu = C \frac{m_\mu^2}{M^2} \quad C \leq O(1)$$

Summary

- Many ways to generate Δa_μ $a_\mu = C \frac{m_\mu^2}{M^2}$ $C \leq O(1)$
- Provides strong constraint on details of new physics

SUSY: M_i not universal, spectrum constrained

Family symmetry: Associated FCNC e.g. $\tau \rightarrow \mu\gamma$.

Constrained Yukawa structure

