

Models for Leading Twist DAs

Patricia Ball

IPPP, Durham



DA06, Durham Sep 28 06

What do we know about T2 DAs?

Standard procedure: construct model for ϕ as **truncated conformal expansion**:

$$\phi(u, \mu^2) \approx 6u(1-u) \left(1 + \sum_1^{n_{\max}} a_n(\mu^2) C_n^{3/2}(2u-1) \right)$$

- for $\pi, \eta_8, \rho, \omega, \phi$: **odd** a_i vanish due to **G-parity**
- n_{\max} typically **2** or **4**
- model closed under renormalisation, obtain correct limit for $\mu \rightarrow \infty$, i.e. asymptotic DA $6u(1-u)$
- constrain $a_{2,4}^\pi$ from **experimental data** for π EM and $\gamma\gamma^* \rightarrow \pi$
(see talk by N. Stefanis)
- or **calculate** $a_{1,2}^{\pi,K,\dots}$ from non-perturbative methods (**QCD sum rules/lattice**)
(talks Thu afternoon)

Why can one truncate?

Actual observable is **convolution integral**

$$\int_0^1 du \phi(u) T(u)$$

with perturbatively calculable amplitude $T(u)$:

contributions from large-order a_n “washed out” for “smooth” T due to oscillatory behaviour of $C_n^{3/2}$:

$$(a) \int du \phi(u) \sqrt{u} \sim \sum \frac{1}{n^3} a_n : \quad \text{strong suppression of } a_n$$

Why can one truncate?

Actual observable is **convolution integral**

$$\int_0^1 du \phi(u) T(u)$$

with perturbatively calculable amplitude $T(u)$:

contributions from large-order a_n “washed out” for “smooth” T due to oscillatory behaviour of $C_n^{3/2}$:

$$(b) \int du \phi(u) \ln u \sim \sum \frac{1}{n^2} a_n : \quad \text{strong suppression of } a_n$$

Why can one truncate?

Actual observable is **convolution integral**

$$\int_0^1 du \phi(u) T(u)$$

with perturbatively calculable amplitude $T(u)$:

contributions from large-order a_n “washed out” for “smooth” T due to oscillatory behaviour of $C_n^{3/2}$:

$$(c) \int du \phi(u) \frac{1}{u} \sim \sum (-1)^n a_n : \quad \text{no suppression!}$$

Why can one truncate?

Actual observable is **convolution integral**

$$\int_0^1 du \phi(u) T(u)$$

with perturbatively calculable amplitude $T(u)$:

contributions from large-order a_n “washed out” for “smooth” T due to oscillatory behaviour of $C_n^{3/2}$:

(d) $\int du \phi(u) \frac{1}{u^2} \rightarrow \infty$: diverges independently of a_n

Why can one truncate?

Actual observable is **convolution integral**

$$\int_0^1 du \phi(u) T(u)$$

with perturbatively calculable amplitude $T(u)$:

contributions from large-order a_n “washed out” for “smooth” T due to oscillatory behaviour of $C_n^{3/2}$:

$$(d) \int du \phi(u) \frac{1}{u^2} \rightarrow \infty : \quad \text{diverges independently of } a_n$$

Truncated conformal expansion OK for $T \sim (\sqrt{u}, \ln u)$, but not necessarily for $1/u$! (And certainly not for $1/u^2$.)

Models beyond Conformal Expansion

Define

$$\Delta \equiv \frac{1}{6} \int du \frac{\phi(u) + \phi(1-u)}{u} = \sum_n a_{2n};$$

“worst case scenario” with **no suppression** of higher a_n

• features in non-leptonic B decays and $\gamma\gamma^* \rightarrow \pi$

Idea (Ball/Talbot 05): define **model for ϕ** in terms of

1. $\Delta \equiv \sum_{\text{even}} a_n$ (constrained to be 1.2 ± 0.2 at $\mu = 1.2$ GeV, for π from $\gamma\gamma^* \rightarrow \pi$)
2. power-like **fall-off behaviour** of a_n in n ($a > 1, b > 0$ arbitrary):

$$a_n = \frac{1}{(n/b + 1)^a} \quad \text{or} \quad a_n = \frac{(-1)^{n/2}}{(n/b + 1)^a}$$

Models beyond Conformal Expansion

Series in Gegenbauers can be summed explicitly:

$$\phi_{a,b}^{\pm}(\Delta) = 6u\bar{u} + \frac{\Delta - 1}{\Delta_{a,b}^{\pm} - 1} \left(\tilde{\phi}_{a,b}^{\pm}(u) - 6u\bar{u} \right), \quad \text{valid for } a > 1 \text{ and } b > 0$$

with

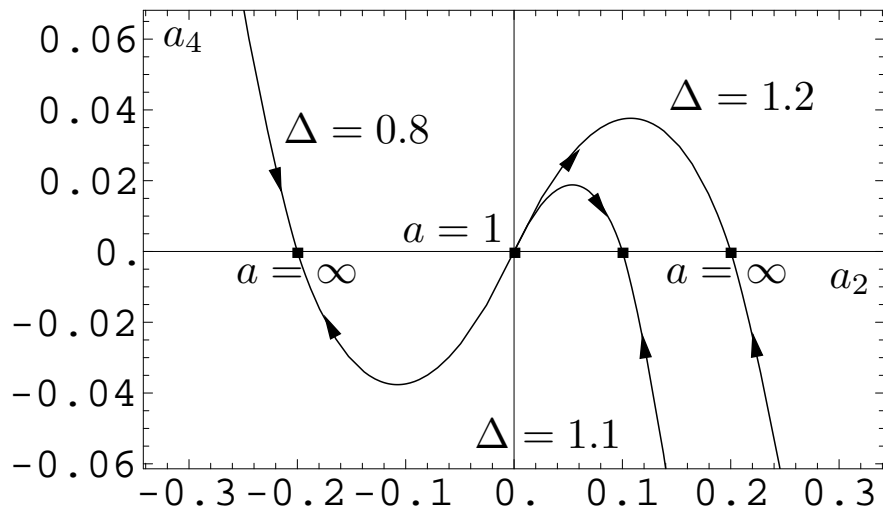
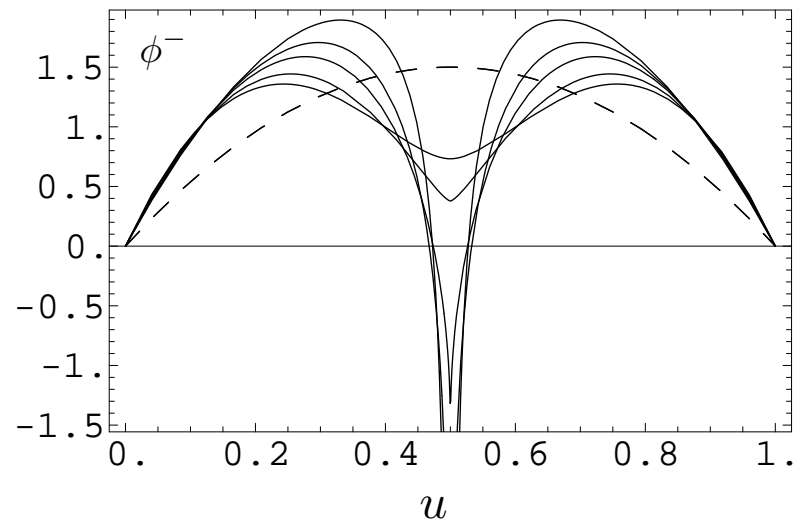
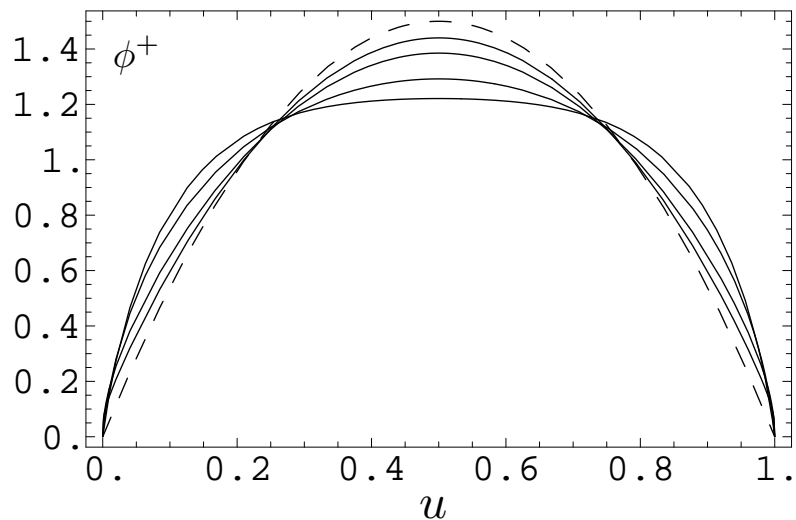
$$\tilde{\phi}_{a,b}^{+}(u) = \frac{3u\bar{u}}{\Gamma(a)} \int_0^1 dt (-\ln t)^{a-1} \left(f(2u - 1, t^{1/b}) + f(2u - 1, -t^{1/b}) \right)$$

f : generating function of Gegenbauer polynomials:

$$f(\xi, t) = \frac{1}{(1 - 2\xi t + t^2)^{3/2}} = \sum_{n=0}^{\infty} C_n^{3/2}(\xi) t^n,$$

$$\Delta_{a,b}^{+} = \left(\frac{b}{2} \right)^a \zeta(a, b/2), \quad \Delta_{a,b}^{-} = \left(\frac{b}{4} \right)^a \{ \zeta(a, b/4) - \zeta(a, 1/2 + b/4) \}$$

Models beyond Conformal Expansion



(Ball/Talbot 05)

How does it work?

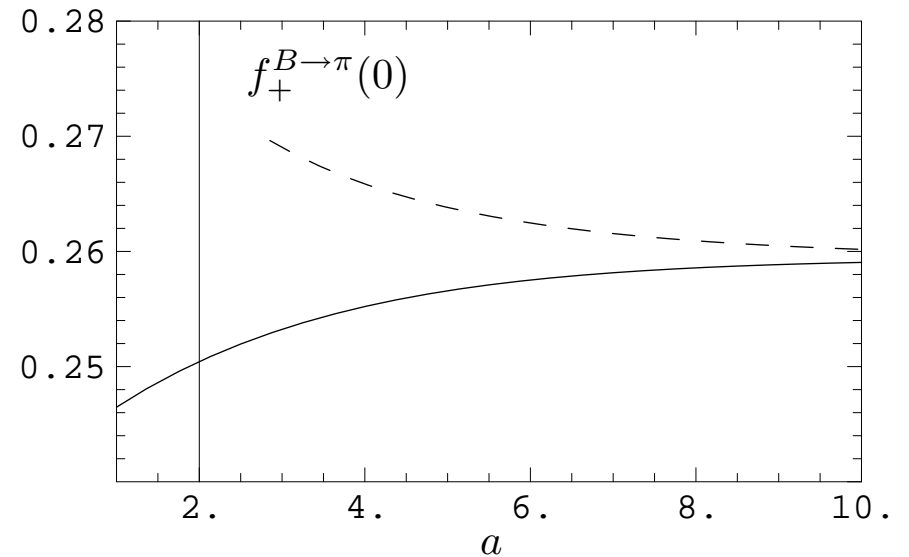
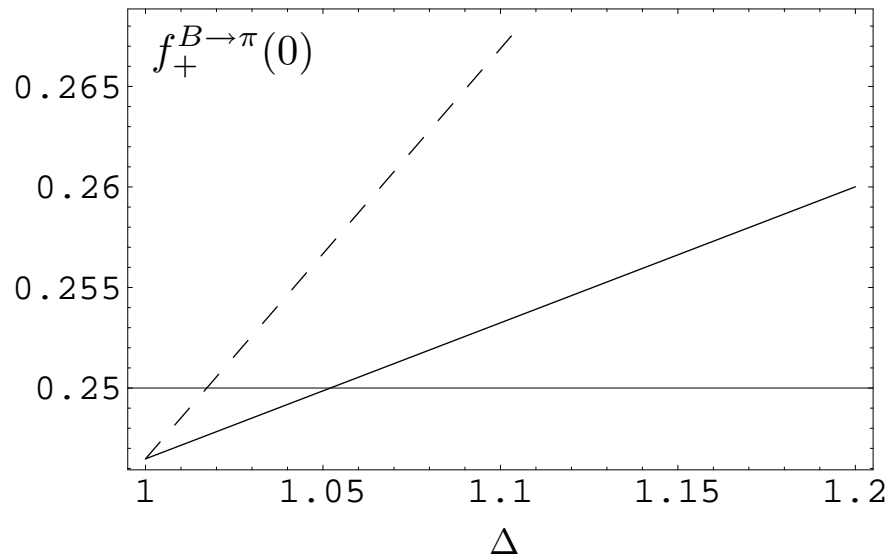
- similar to parametrisation of **parton distribution functions** in DIS:
 - choose suitable model defined at low scale in terms of a few parameters
 - use RG scaling to get PDF at higher scales
 - fit model parameters to experimental results

- **evolution equation**:
$$\mu^2 \frac{\partial}{\partial \mu^2} \phi(u, \mu^2) = \int_0^1 dv V(u, v, \mu^2) \phi(v, \mu^2);$$

kernel V known to 2-loop accuracy for π, ρ_{\parallel} (Mikhailov/Radyushkin 85) and 1-loop accuracy for ρ_{\perp} (Ball/Talbot 05)

- relation between (a_2, a_4) and (Δ, a) is one-to-one (at a fixed scale μ) (see p.5): difference in resulting observables/convolutions is, if nothing more, **estimate of truncation error** of conformal expansion
- **experimental data**: far fewer than in DIS...

Application: $f_+^{B \rightarrow \pi}(0)$ from LCSR



Solid: positive sign a_n

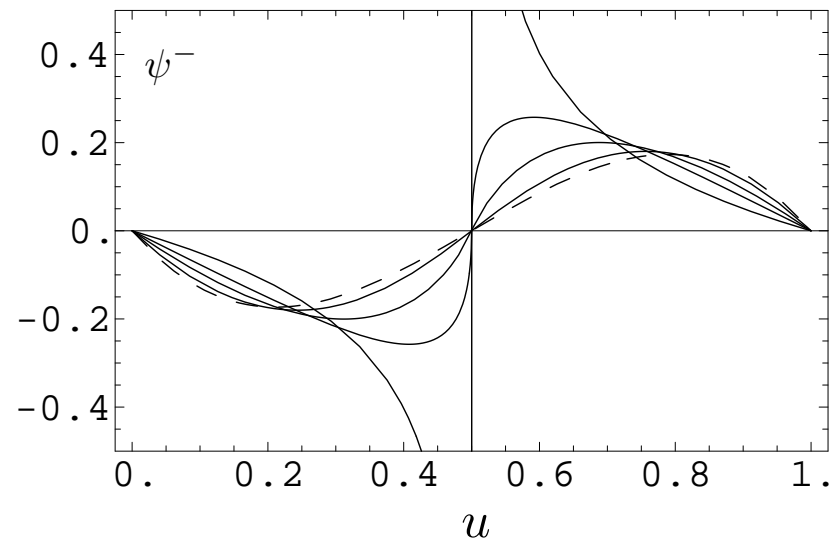
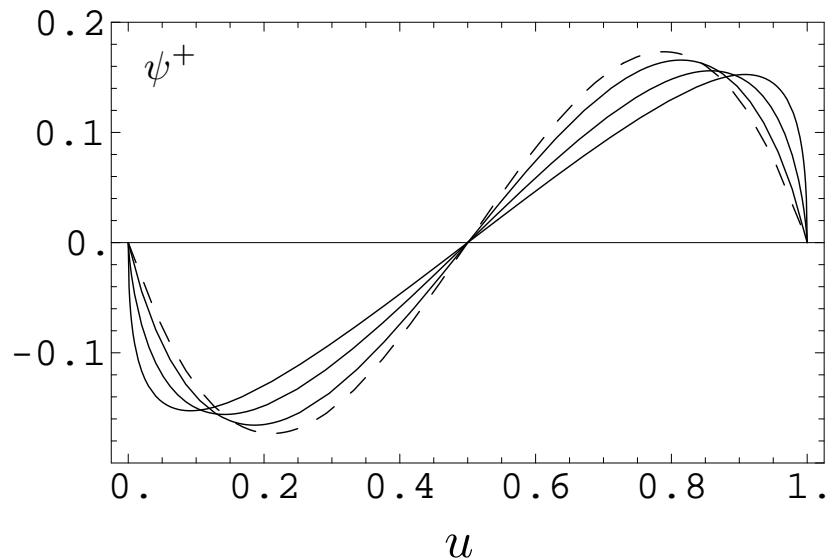
Dashed: alternating sign a_n

Models for odd a_n , normalised to a_1

$$\psi_c^\pm(u) = a_1 \left(\frac{3}{2}\right)^c \tilde{\psi}_c^\pm(u) \quad (c \text{ is power of fall-off})$$

with $\tilde{\psi}_c^+(u) = \frac{3u\bar{u}}{\Gamma(c)} \int_0^1 dt (-\ln t)^{c-1} \left(f(2u-1, \sqrt{t}) - f(2u-1, -\sqrt{t}) \right)$

$$\tilde{\psi}_c^-(u) = \frac{3u\bar{u}}{i\Gamma(c)} \int_0^1 dt (-\ln t)^{c-1} \left(f(2u-1, i\sqrt{t}) - f(2u-1, -i\sqrt{t}) \right)$$



Summary & Conclusions

- alternative to truncated conformal expansion based on 2 parameters for symmetric part of $\phi(u)$:
 - first inverse moment
 - power of fall-off of a_{2n}
- and two parameters for antisymmetric part of $\phi(u)$:
 - a_1
 - power of fall-off of a_{2n+1}
- closed form at fixed low-energy scale, use evolution equation to go to higher scales
- **deliverables**: fast and safe numerical code for solution of evolution equation
- **desirables 1**: two-loop evolution kernel for ρ_\perp
- **desirables 2**: fix parameters from fits to experimental data