

# Exploring the Standard Candle

## The many faces of $B \rightarrow X_s \gamma$ decay

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# References (2004-2005)

- MN: “Renormalization-Group Improved Calculation of the  $B \rightarrow X_s \gamma$  Branching Ratio” [hep-ph/0408179] → EPJC
- MN: “Advanced Predictions for Moments of the  $B \rightarrow X_s \gamma$  Photon Spectrum” [hep-ph/0506245] → PRD
- B. Lange, MN, G. Paz: “Theory of Charmless Inclusive  $B$  Decays and the Extraction of  $|V_{ub}|$ ” [hep-ph/0504071] → PRD
- B. Lange, MN, G. Paz: “A Two-Loop Relation between Inclusive Radiative and Semileptonic  $B$ -Decay Spectra” [hep-ph/0508178] → JHEP
- T. Becher, MN: “Toward a NNLO Prediction for the  $B \rightarrow X_s \gamma$  Decay Rate with a Cut on Photon Energy: I. Two-Loop Result for the Soft Function” [hep-ph/0512208]



# Exploring the Standard Candle

## The many faces of $B \rightarrow X_s \gamma$ decay

Introduction

Soft-collinear factorization

Scale separation

Results:

- $\Gamma(B \rightarrow X_s \gamma)$ , New Physics, b-quark mass
- Extraction of  $|V_{ub}|$

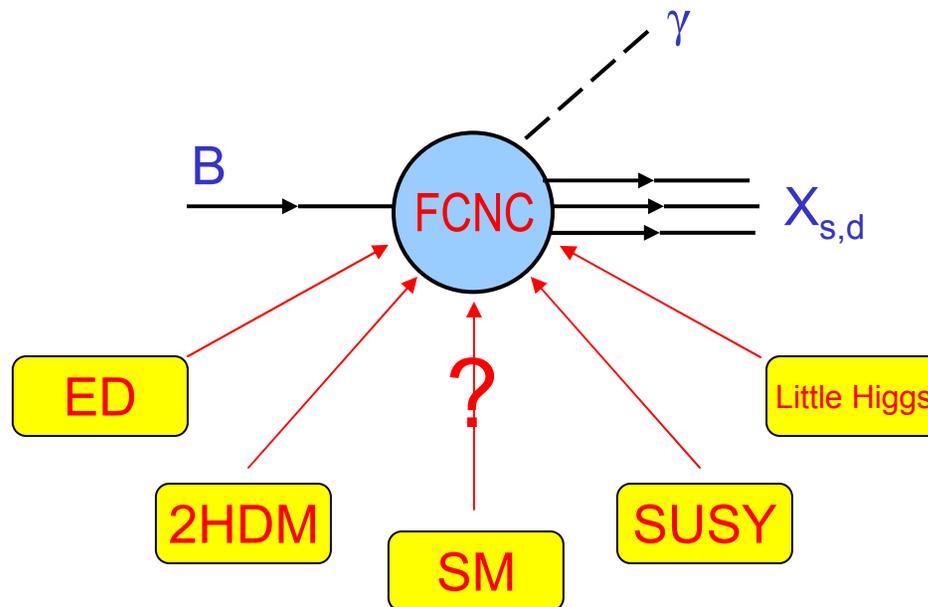
# Introduction

Inclusive and exclusive  
 $B \rightarrow X_s \gamma$  decays,  
Cut on photon energy



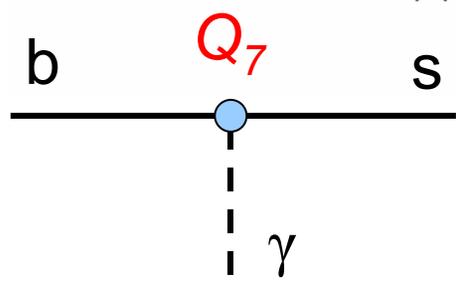
# Rare radiative decays $B \rightarrow X\gamma$

- Prototype FCNC processes



# Inclusive decay rate

$$\Gamma(B \rightarrow X_s \gamma) = \frac{G_F^2 \alpha}{32\pi^4} |V_{tb} V_{ts}^*|^2 m_b^5 C_{7\gamma}^2$$

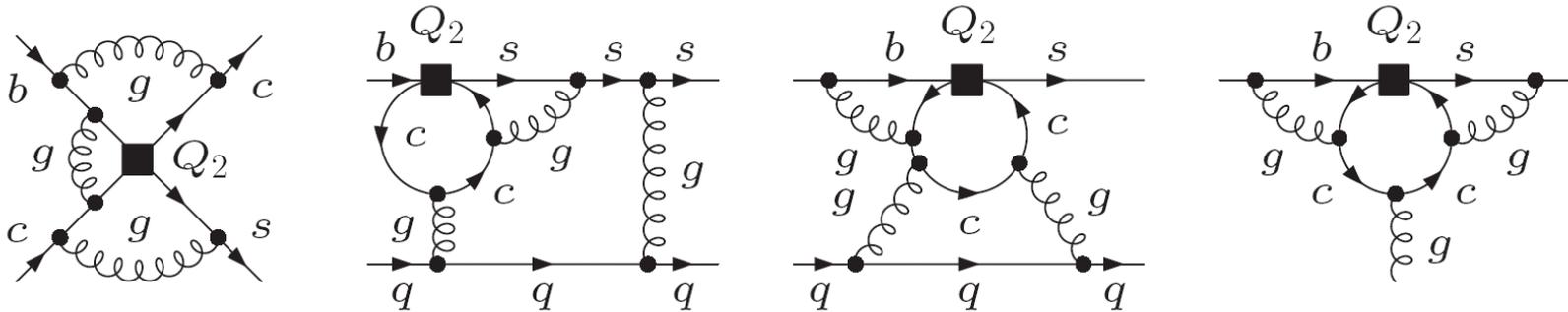
$$\times \left\{ 1 + O(\alpha_s) + O(\alpha_s^2) + O\left(\frac{\Lambda_{\text{QCD}}^2}{m_b^2}\right) + \dots \right\}$$


Calculable using OPE

- NLO calculation completed (1991-2001)
- NNLO calculation in progress (2003-...)

# Inclusive decay rate at NNLO

- Effective weak Hamiltonian at NNLO
  - 3-loop (☺) and 4-loop (⌚) anomalous dimensions
  - 2-loop (☺) and 3-loop (⌚) matching conditions

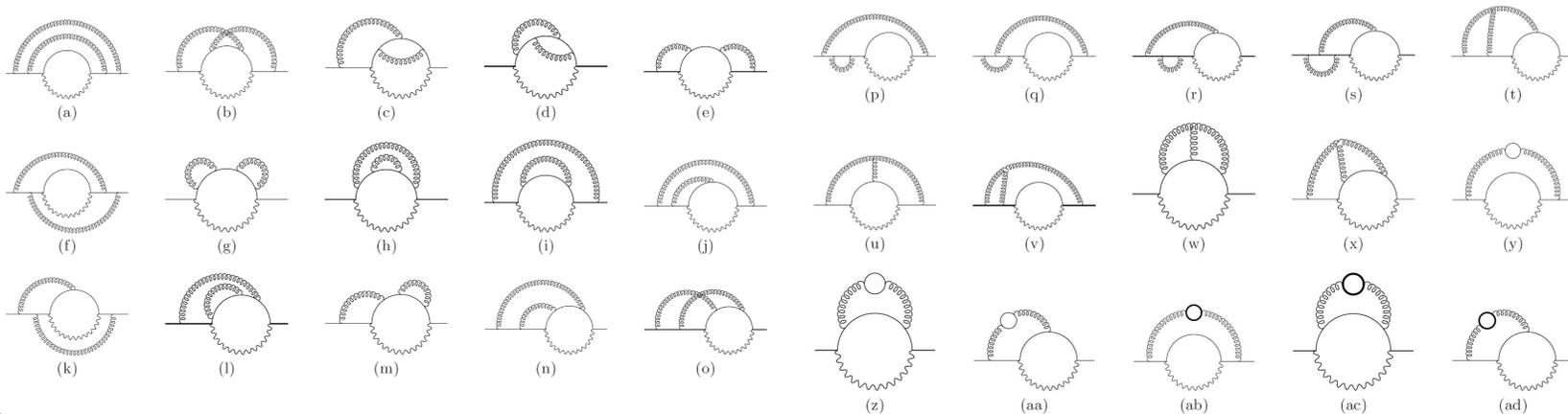


[Gorbahn, Haisch: hep-ph/0411071  
Gorbahn, Haisch, Misiak: hep-ph/0504194]

# Inclusive decay rate at NNLO

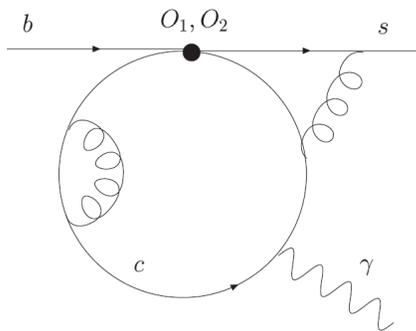
- Matrix elements at NNLO
  - 2-loop matrix elements of dipole operators  
(☺)

[Melnikov, Mitov: hep-ph/0505097  
Blokland, Czarnecki, Misiak, Slusarczyk,  
Tkachov: hep-ph/0506055]



# Inclusive decay rate at NNLO

- Matrix elements at NNLO
  - 3-loop (☹) penguin matrix elements



[Bieri, Greub, Steinhauser: hep-ph/0302051  
Asatrian, Greub, Hovhannisyan, Hurth,  
Poghosyan: hep-ph/0505068]

- One of the hardest calculations in QFT, with many people devoting significant resources

# Exclusive decay rates

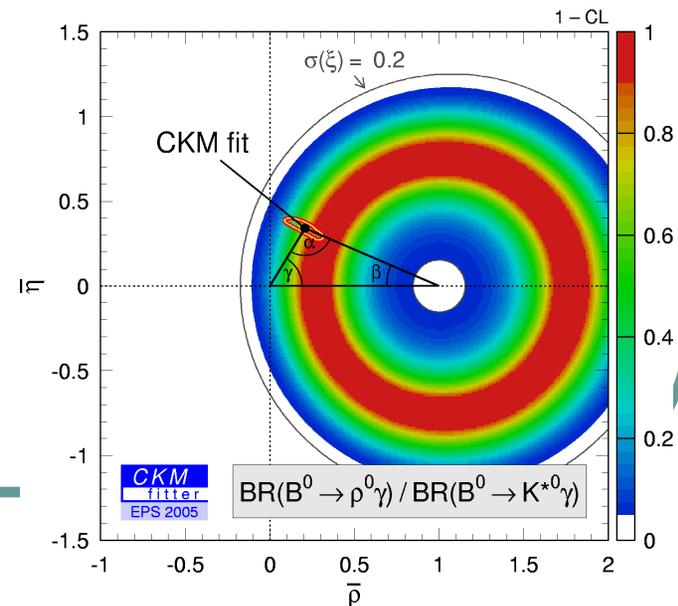
- QCD factorization formula

$$\langle V \gamma(\epsilon) | Q_i | \bar{B} \rangle = \left[ F^{B \rightarrow V}(0) T_i^I + \int_0^1 d\xi dv T_i^{II}(\xi, v) \Phi_B(\xi) \Phi_V(v) \right] \cdot \epsilon$$

[Beneke, Feldmann, Seidel: hep-ph/0106067  
 Bosch, Buchalla: hep-ph/0106081

**Proof:** Becher, Hill, MN: hep-ph/0503263]

- Hadronic uncertainties
- Useful constraint on UT  
 from  $\Gamma(B \rightarrow \rho \gamma) / \Gamma(B \rightarrow K^* \gamma)$   
 $\rightarrow |V_{td}/V_{ts}|$



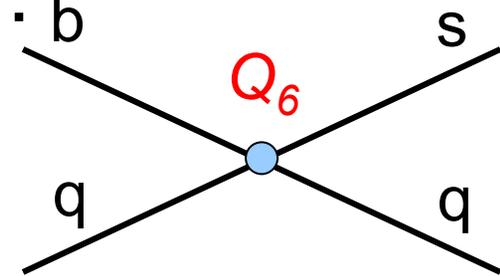
# Other interesting observables

- Exclusive isospin asymmetry:

$$\Delta_{0-} \equiv \frac{\Gamma(\bar{B}^0 \rightarrow \bar{K}^{*0}\gamma) - \Gamma(B^- \rightarrow K^{*-}\gamma)}{\Gamma(\bar{B}^0 \rightarrow \bar{K}^{*0}\gamma) + \Gamma(B^- \rightarrow K^{*-}\gamma)}$$

- Sensitive probe of new contributions to QCD penguin operators

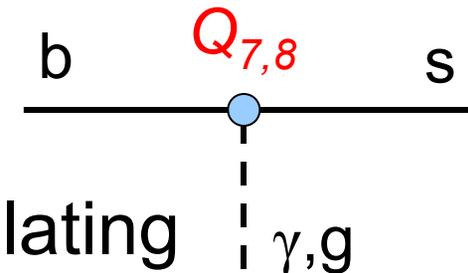
[Kagan, MN: hep-ph/0110078]



- Inclusive CP asymmetry:

- Sensitive probe of new, CP-violating contributions to dipole operators

[Kagan, MN: hep-ph/9803368]

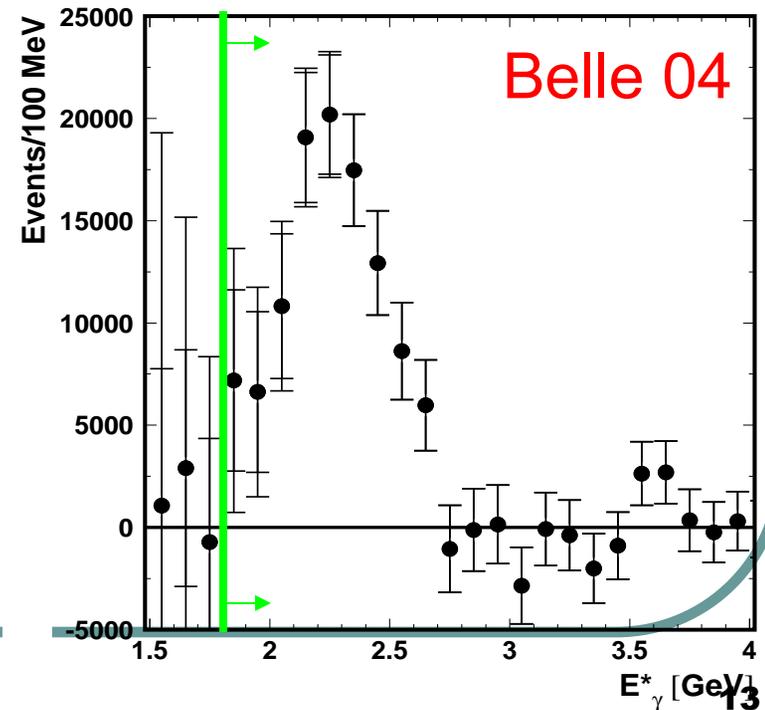
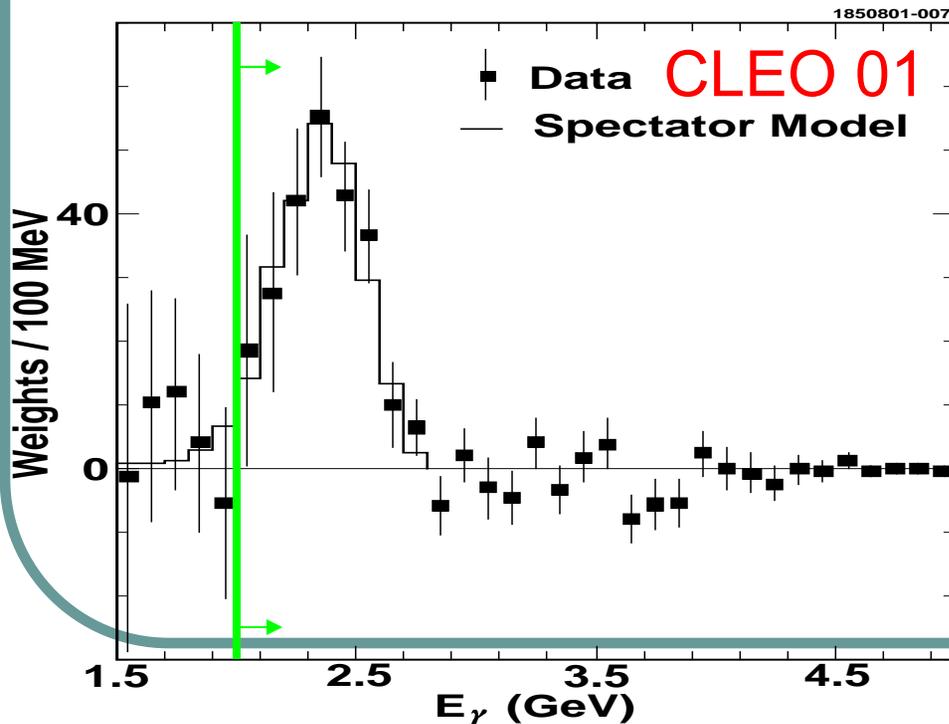


# Partially inclusive decay rate

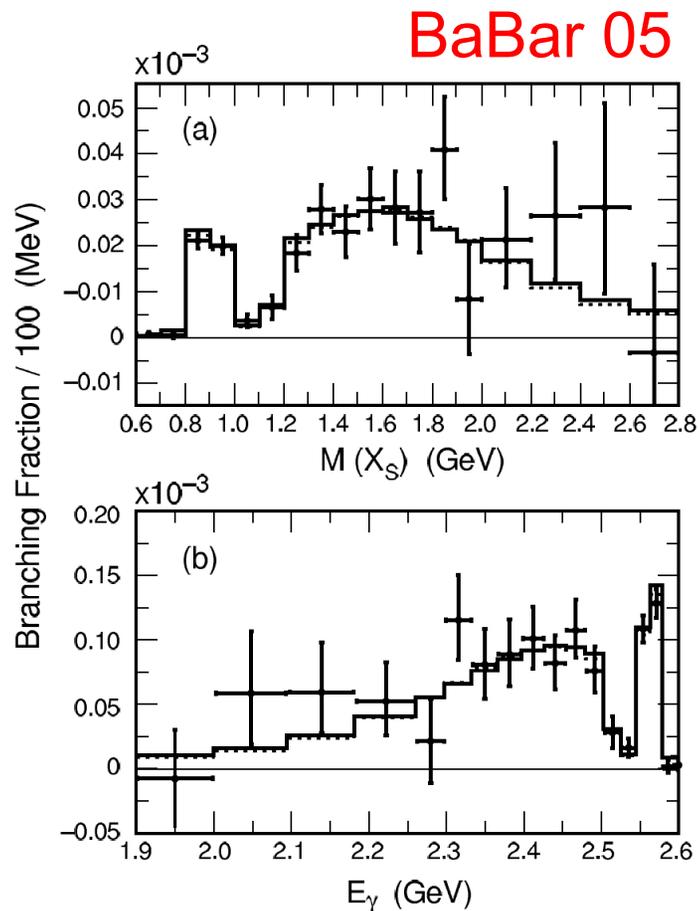
- Total decay rate cannot be measured for various reasons
  - Soft-photon divergence (need combination with  $\Gamma(B \rightarrow X_s g)$ )
  - Need to suppress background from  $B \rightarrow \psi X_s$  followed by  $\psi \rightarrow X \gamma$  [Misiak: hep-ph/0002007]
  - Experimental signature is high-energy photon (otherwise huge background)

# Partially inclusive decay rate

- Restriction to high-energy part of photon spectrum:  $E_\gamma > E_0 = 1.8 \text{ GeV}$  or larger (in B-meson rest frame)



# Partially inclusive decay rate



- Presence of the photon-energy cut leads to significant complications in the theoretical analysis!

[MN: [hep-ph/0408179](https://arxiv.org/abs/hep-ph/0408179)]

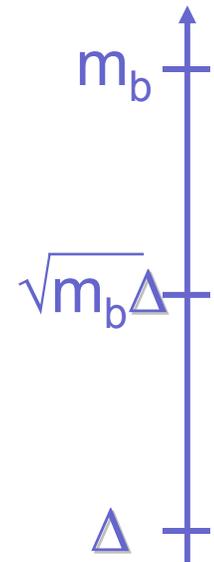
# Soft-collinear factorization

Different scales,  
Factorization formula,  
Non-local matrix elements

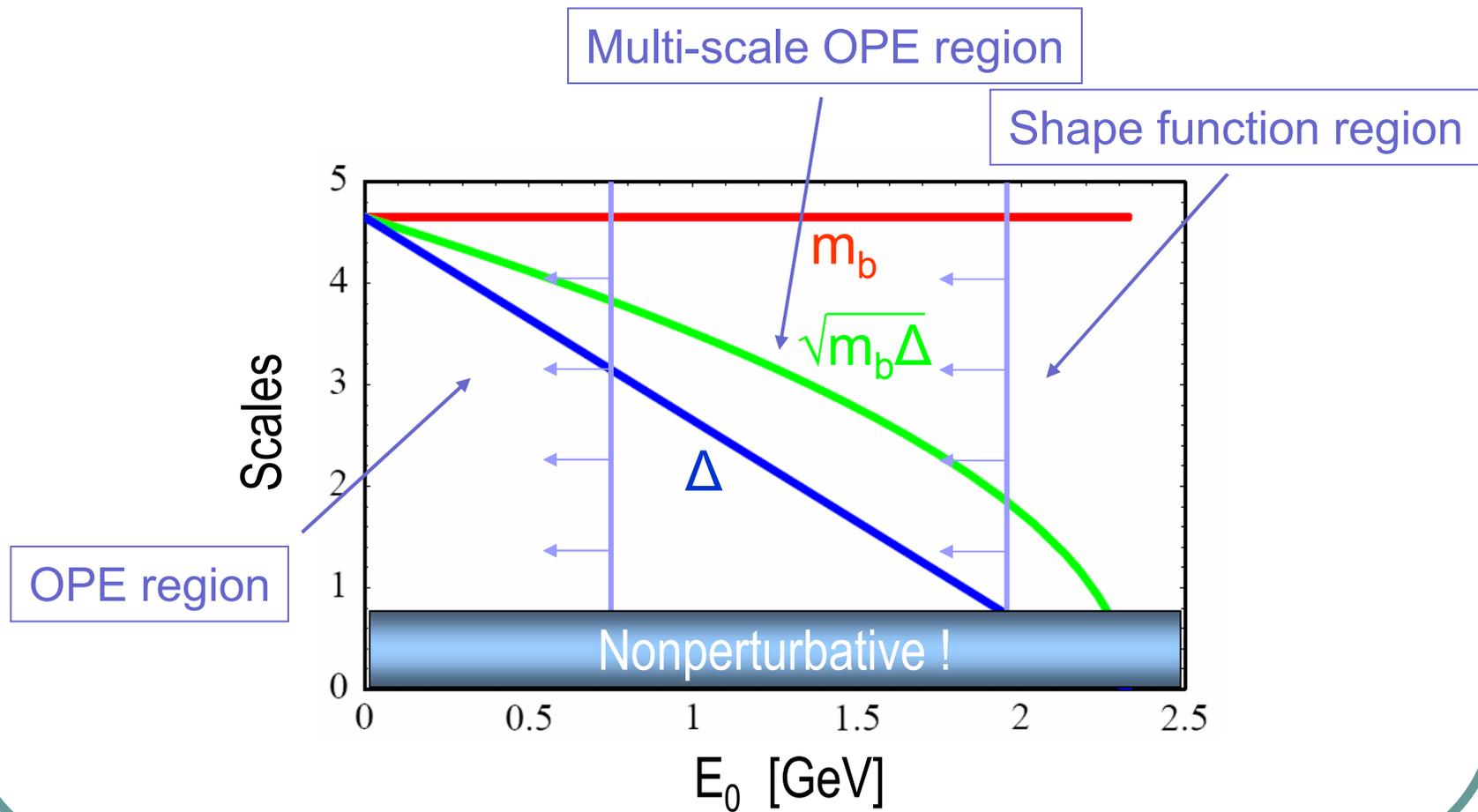


# Relevance of different scales

- Presence of photon-energy cut introduces new scale  $\Delta = m_b - 2E_0$
- Must distinguish:
  - Hard scale  $m_b$ : quantum corrections to effective weak-interaction vertices
  - Intermediate scale  $\sqrt{m_b\Delta}$ : invariant mass of final-state hadronic jet
  - Soft scale  $\Delta$ : scale at which internal structure of B-meson is probed



# Relevance of different scales



# Factorization theorem

- At leading power in  $\Lambda_{\text{QCD}}/m_b$ , the decay rate factorizes into a convolution of three objects:

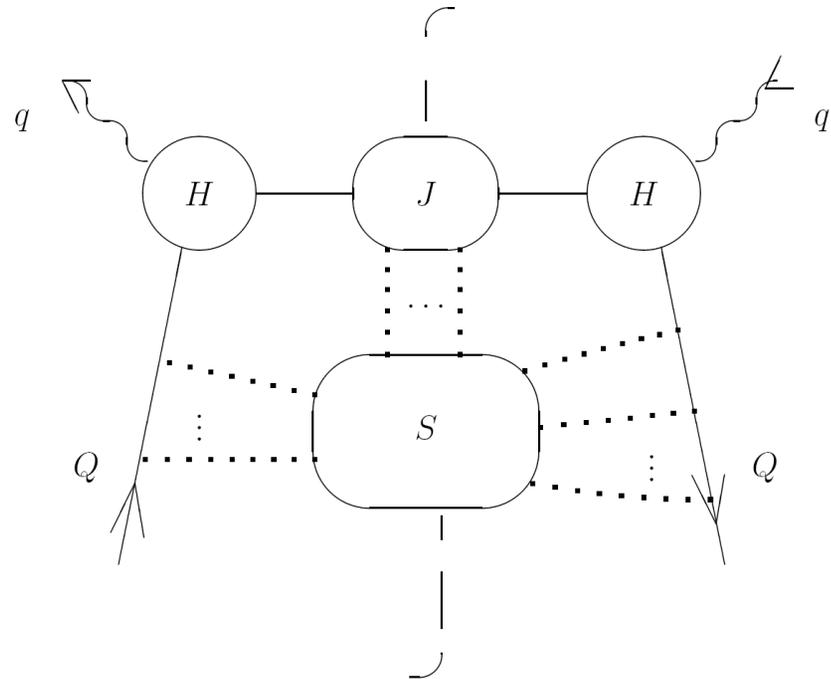
$$\frac{d\Gamma}{dE_\gamma} = \frac{G_F^2 \alpha}{2\pi^4} |V_{tb} V_{ts}^*|^2 m_b^2 E_\gamma^3$$

$$\times |H_\gamma(\mu)|^2 \int_0^{M_B - 2E_\gamma} d\hat{\omega} m_b J(m_b(M_B - 2E_\gamma - \hat{\omega}), \mu) S(\hat{\omega}, \mu) + \dots$$

[MN: hep-ph/9312311

Korchensky, Serman: hep-ph/9407344

**Proof:** Bauer, Pirjol, Stewart: hep-ph/0109045]



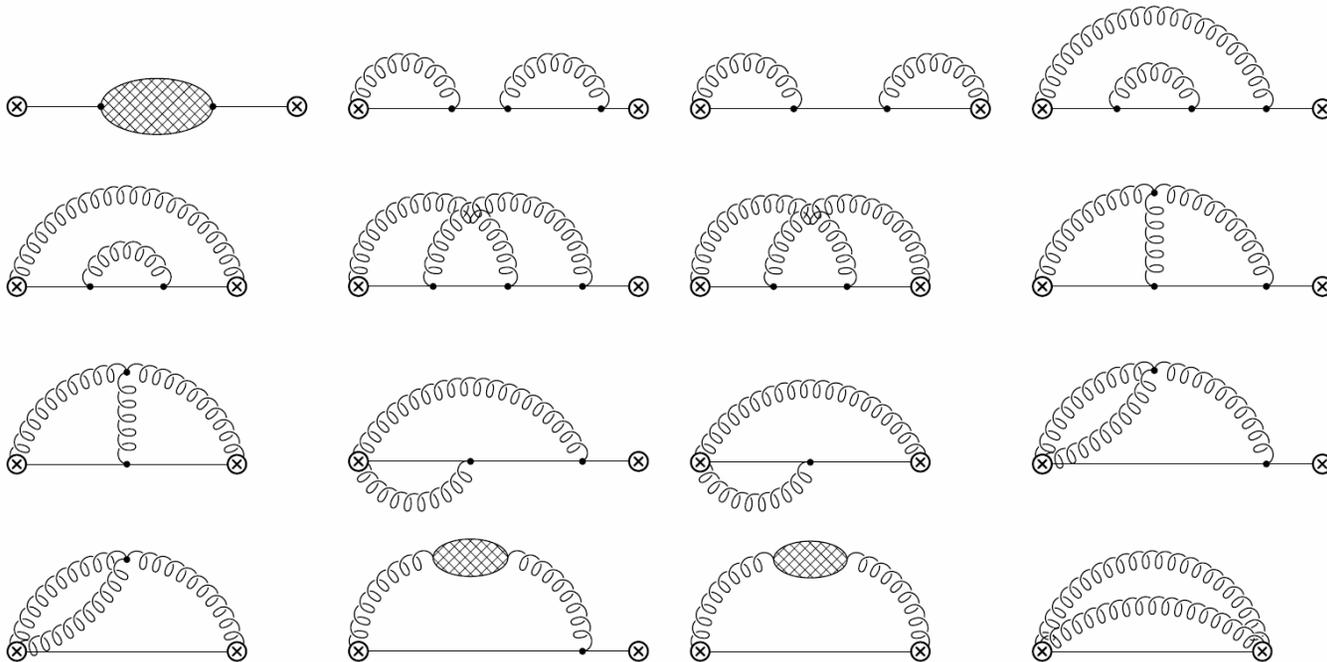
# Non-local matrix elements

- Jet function  $J(p^2, \mu)$ :
  - Fourier transform of hard-collinear quark propagator dressed by Wilson line  
*perturbative quantity*
- Shape function  $S(\hat{\omega}, \mu)$ :
  - Fourier transform of forward matrix element  $\langle B | \bar{h}_v(0) [0, x] h_v(x) | B \rangle$  at light-like separation  
*nonperturbative object* (parton distribution)

# Non-local matrix elements

$$J(p^2, \mu) = \delta(p^2) \left[ 1 + \frac{C_F \alpha_s(\mu)}{4\pi} (7 - \pi^2) \right] + \frac{C_F \alpha_s(\mu)}{4\pi} \left[ \frac{1}{p^2} \left( 4 \ln \frac{p^2}{\mu^2} - 3 \right) \right]_*^{[\mu^2]}$$

+  $O(\alpha_s^2)$  [Becher, MN: in preparation]



# Hadronic physics

- Shape function integrals can be written as contour integrals in the complex  $\omega$ -plane:

$$\int_{-\bar{\Lambda}}^{\Delta} d\omega S(\omega, \mu) f(\omega) \propto \oint_{|\omega|=\Delta} d\omega f(\omega) \langle \bar{B}(v) | \bar{h}_v \frac{1}{\omega + in \cdot D + i\epsilon} h_v | \bar{B}(v) \rangle$$

- For sufficiently large  $\Delta$ , right-hand side can expand in *local* operators
- Expansion in  $(\Lambda_{\text{QCD}}/\Delta)^n$  and  $\alpha_s(\Delta)$   
→ not a heavy-quark expansion

# Scale separation

Evolution equations,  
Exact solutions



# Evolution equations

- Use RGEs to disentangle contributions

associated with:

- Hard scale  $\mu_h \sim m_b$

- Intermediate scale  $\mu_i \sim \sqrt{m_b \Delta}$

- Soft scale  $\mu_0 \sim \Delta$

- Set  $\mu = \mu_i$  in factorization formula and evolve hard and soft functions to the hard and soft scales, respectively

$$\frac{d\Gamma}{dE_\gamma} = \frac{G_F^2 \alpha}{2\pi^4} |V_{tb} V_{ts}^*|^2 m_b^2 E_\gamma^3 \times |H_\gamma(\mu)|^2 \int_0^{M_B - 2E_\gamma} d\hat{\omega} m_b J(m_b(M_B - 2E_\gamma - \hat{\omega}), \mu) S(\hat{\omega}, \mu) + \dots$$

# Evolution equations

$$\frac{d}{d \ln \mu} H_\gamma(\mu) = \left[ -\Gamma_c(\alpha_s) \ln \frac{\mu}{m_b} + \gamma_J(\alpha_s) \right] H_\gamma(\mu)$$

$$\frac{d}{d \ln \mu} S(\hat{\omega}, \mu) = \left[ 2\Gamma_c(\alpha_s) \ln \frac{\mu}{\hat{\omega}} - 2\gamma(\alpha_s) \right] S(\hat{\omega}, \mu)$$

$$+ 2\Gamma_c \int_0^{\hat{\omega}} d\hat{\omega}' \frac{S(\hat{\omega}', \mu) - S(\hat{\omega}, \mu)}{\hat{\omega} - \hat{\omega}'}$$

Sudakov logarithms

- **Anomalous dimensions:**

$$\Gamma_c(\alpha_s) = 0.424\alpha_s + 0.271\alpha_s^2 + 0.216\alpha_s^3 + \dots$$

[Korchemsky, Radyushkin: 1987  
Moch, Vermaseren, Vogt:  
hep-ph/0403192]

$$\gamma_J(\alpha_s) = -0.531\alpha_s - 0.200\alpha_s^2 + \dots$$

[MN: hep-ph/0408179]

$$\gamma(\alpha_s) = -0.212\alpha_s - 0.389\alpha_s^2 + \dots$$

[Korchemsky, Marchesini: hep-ph/9210281  
Gardi: hep-ph/0501257 ]

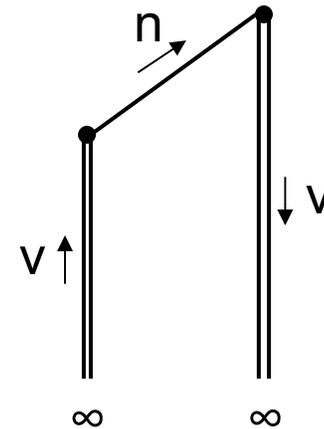
# Exact solutions

$$|H_\gamma(\mu_i)|^2 = U_1(\mu_h, \mu_i) |H_\gamma(\mu_h)|^2$$

$$S(\hat{\omega}, \mu_i) = U_2(\mu_i, \mu_0) \frac{e^{-\gamma E \eta}}{\Gamma(\eta)} \int_0^{\hat{\omega}} d\hat{\omega}' \frac{S(\hat{\omega}', \mu_0)}{\mu_0^\eta (\hat{\omega} - \hat{\omega}')^{1-\eta}}$$

- With:

$$\eta = 2 \int_{\mu_0}^{\mu_i} \frac{d\mu}{\mu} \Gamma_c[\alpha_s(\mu)]$$



# Wonderful formula

- Master formula for the partial rate:

$$\Gamma \sim H(\mu_h) * U(\mu_h, \mu_i) * J(\mu_i) * U(\mu_i, \mu_0) * M(\mu_0)$$

QCD  $\rightarrow$  SCET  $\rightarrow$  RG evolution  $\rightarrow$  HQET  $\rightarrow$  RG evolution  $\rightarrow$  Local OPE

$$\begin{aligned}\mu_h &\sim m_b \\ \mu_i &\sim \sqrt{m_b \Delta} \\ \mu_0 &\sim \Delta\end{aligned}$$

Perturbation theory

Hadronic physics

# Wonderful formula

- Perform triple integral exactly: [MN: hep-ph/0506245]

$$\Gamma_{\text{OPE}}(\Delta) = \frac{G_F^2 \alpha}{32\pi^4} |V_{tb} V_{ts}^*|^2 m_b^3 \overline{m}_b^2(\mu_h) |H_\gamma(\mu_h)|^2 U_1(\mu_h, \mu_i) U_2(\mu_i, \mu_0) \\ \times \tilde{j}\left(\ln \frac{m_b \mu_0}{\mu_i^2} + \partial_\eta, \mu_i\right) \tilde{s}(\partial_\eta, \mu_0) \frac{e^{-\gamma_E \eta}}{\Gamma(1+\eta)} \left(\frac{\Delta}{\mu_0}\right)^\eta \left[1 - \frac{\eta(1-\eta)}{6} \frac{\mu_\pi^2}{\Delta^2} + \dots\right]$$

Scales:  $\mu_h \sim m_b$   
 $\mu_i \sim \sqrt{m_b \Delta}$   
 $\mu_0 \sim \Delta$

Explicit dependence on  $\Delta$

Kinetic-energy parameter

# Wonderful formula

- The jet and soft functions,  $\tilde{j}$  and  $\tilde{s}$ , are polynomials of their arguments and are derived from integrals of the original jet and shape functions (J and S): [\[MN: hep-ph/0506245\]](#)

$$\tilde{s}(L, \mu) = 1 + \left(-0.873 - 0.424L - 0.424L^2\right) \alpha_s$$
$$+ \left(-0.603 + 0.750L + 0.471L^2 + 0.368L^3 + 0.090L^4\right) \alpha_s^2 + \dots$$

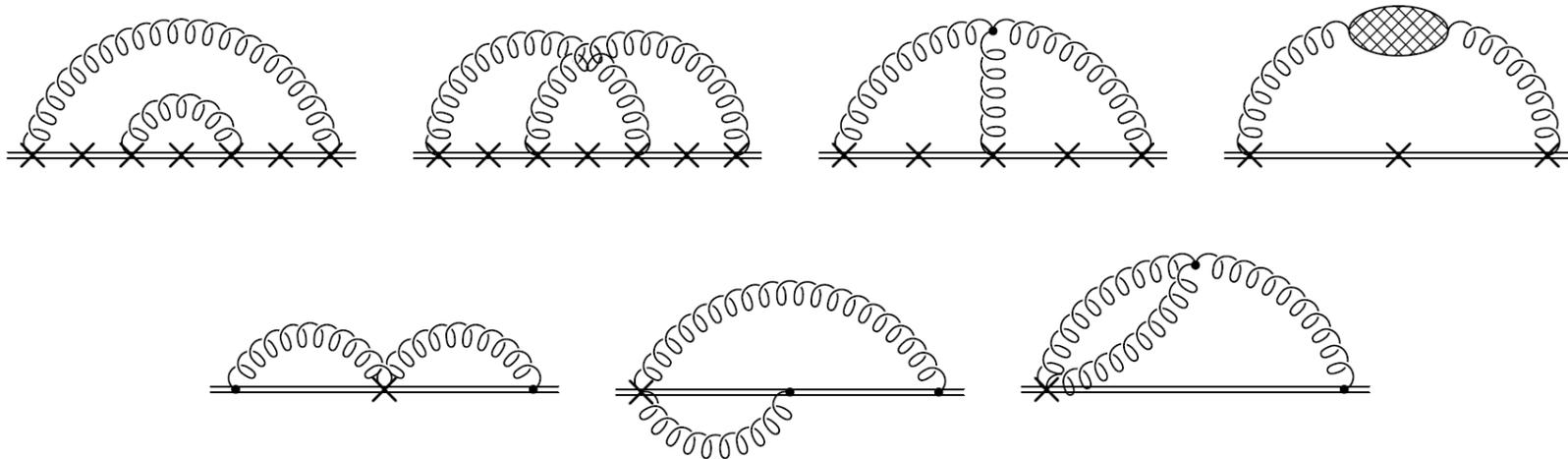
$$\tilde{j}(L, \mu) = 1 + \left(0.045 - 0.318L + 0.212L^2\right) \alpha_s$$
$$+ \left(\dots + 0.145L + 0.301L^2 - 0.114L^3 + 0.023L^4\right) \alpha_s^2 + \dots$$

[\[Becher, MN: hep-ph/0512208\]](#)

work in progress

# Wonderful formula

- Two-loop calculation of soft function:



[Becher, MN: hep-ph/0512208]

# Results

Partial decay rate,  
Implications for New Physics,  
Determination of  $m_b$



# Partial $B \rightarrow X_s \gamma$ branching ratio

- Theoretical calculation with a cut at  $E_0 = 1.8\text{GeV}$  (NLO):

$$\text{Br}(1.8\text{GeV}) = (3.30 \pm 0.33[\text{pert}] \pm 0.17[\text{pars}]) \cdot 10^{-4}$$

[MN: hep-ph/0408179]

- Significant reduction of perturbative error expected at NNLO

- Experiment (Belle 2004):

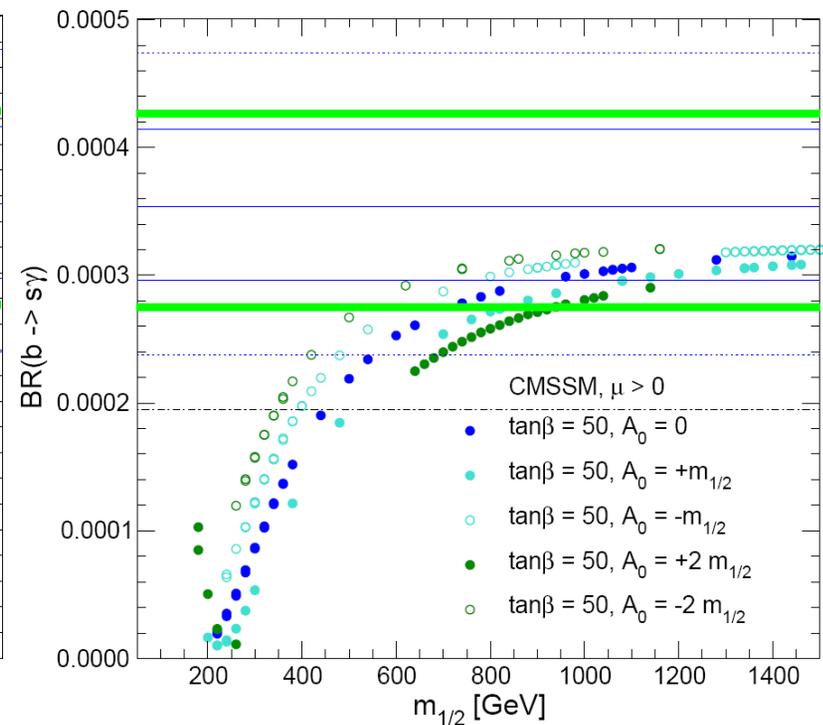
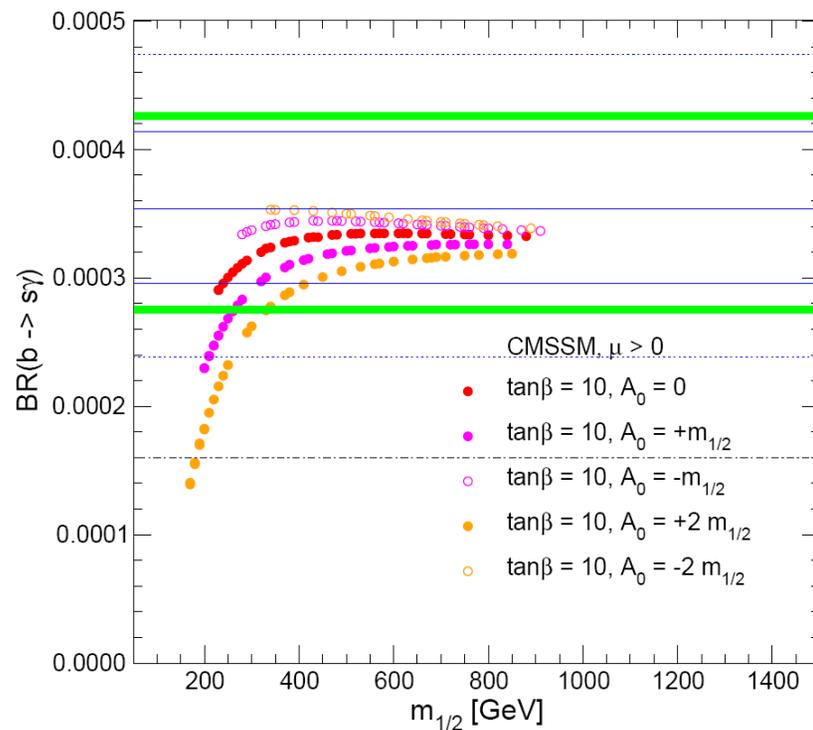
$$\text{Br}(1.8\text{GeV}) = (3.38 \pm 0.30[\text{stat}] \pm 0.28[\text{syst}]) \cdot 10^{-4}$$

# Implications for New Physics

- Increased theory errors and improved agreement with experiment weaken constraints on parameter space of New Physics models!
- E.g., type-II two-Higgs doublet model:
  - $m(H^+) > 200 \text{ GeV}$  (at 95% CL)  
(compared with previous bound of 500 GeV)

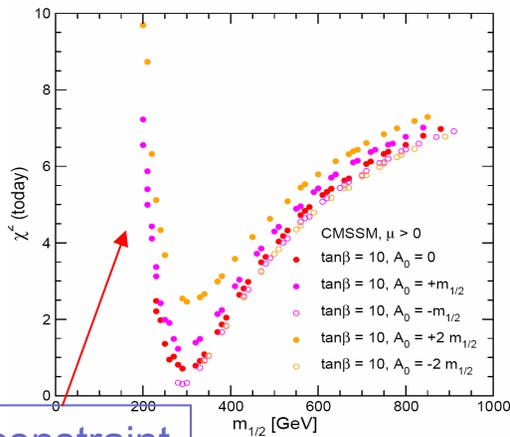
# Implications for New Physics

## ● Bounds on CMSSM parameters:

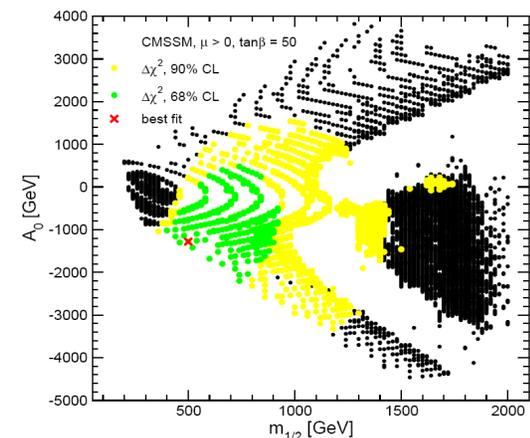
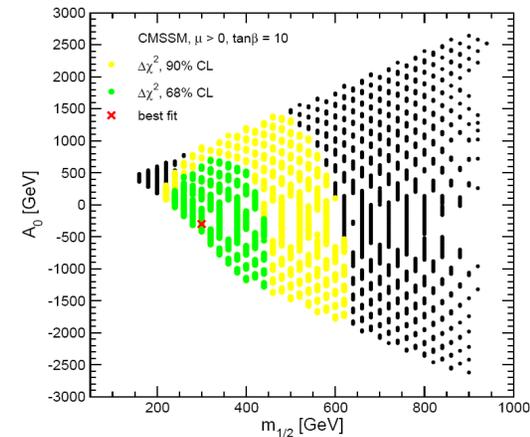
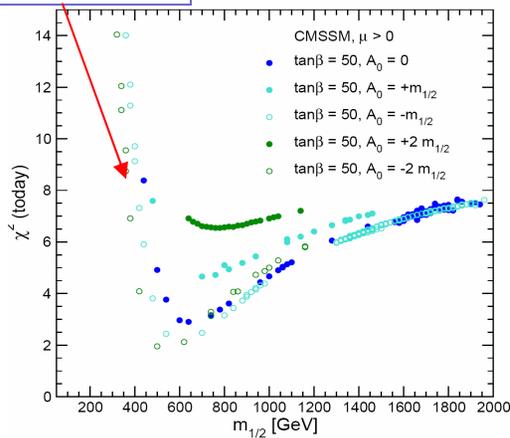


[Ellis, Heinemeyer, Olive, Weiglein: hep-ph/0411216]

# Implications for New Physics



$B \rightarrow X_s \gamma$  constraint



[Ellis, Heinemeyer, Olive, Weiglein: hep-ph/0411216]

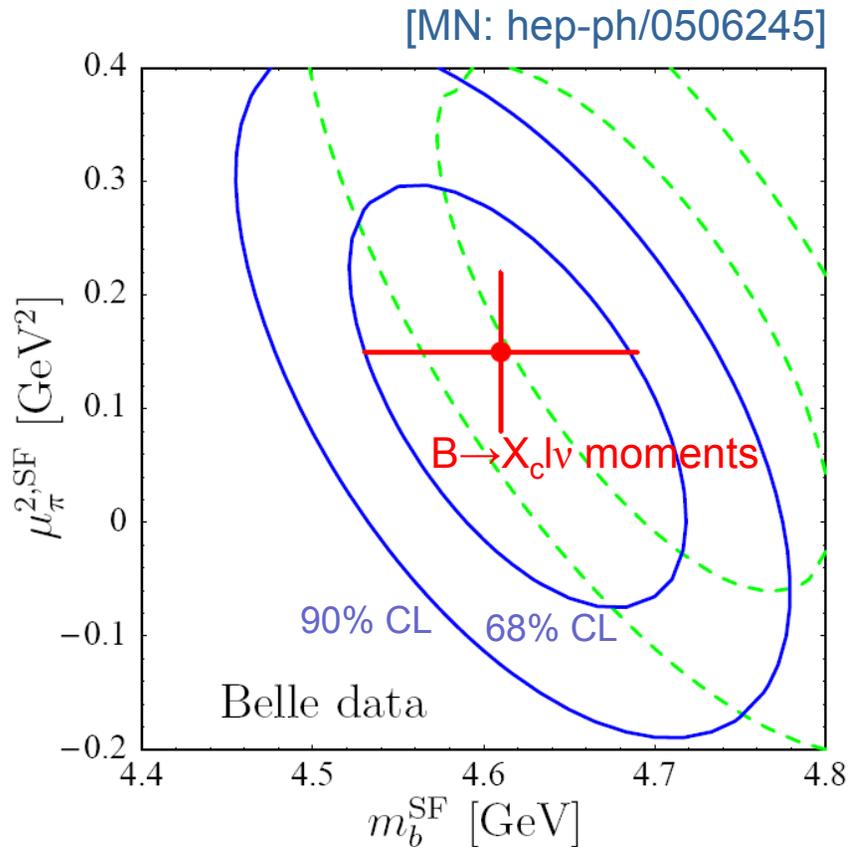
# Moments of the photon spectrum

- Marvelous QCD laboratory
- Extraction of heavy-quark parameters  $(m_b, \mu_\pi^2)$  with exquisite precision
- NNLO accuracy already achieved for  $\langle E_\gamma \rangle$  and  $\langle E_\gamma^2 \rangle - \langle E_\gamma \rangle^2$ :
  - Full two-loop corrections (+ 3-loop running)
  - Same accuracy for leading power corrections  $\sim (\Lambda_{\text{QCD}}/\Delta)^2$ ; fixed-order results for  $1/m_b$  terms

# Predictions for moments

	$O(1)$	$O(1/m_b)$	$O(1/m_b^2)$
Perturbation Theory	Complete resummation at NNLO	$\alpha_s^2$	$\alpha_s^2$
Hadronic Parameters	$m_b, \mu_\pi^2$	$\mu_\pi^2, \rho_D^3, \rho_{LS}^3$	$\rho_D^3, \rho_{LS}^3$

# Fit to Belle data ( $E_0 = 1.8$ GeV)



(parameters defined in shape-function scheme)

- Fit results:

$$m_b = (4.62 \pm 0.10_{\text{exp}} \pm 0.03_{\text{th}}) \text{ GeV}$$

$$\mu_{\pi}^2 = (0.11 \pm 0.19_{\text{exp}} \pm 0.08_{\text{th}}) \text{ GeV}^2$$

- Combined results ( $B \rightarrow X_s \gamma$  and  $B \rightarrow X_c l \nu$ ):

$$m_b = (4.61 \pm 0.06) \text{ GeV}$$

$$\mu_{\pi}^2 = (0.15 \pm 0.07) \text{ GeV}^2$$

# Fit to Belle data ( $E_0 = 1.8 \text{ GeV}$ )

- Theoretically, the most precise extraction of  $m_b$  to date (0.7% accuracy!)
- Translation to  $\overline{\text{MS}}$  scheme (2 loops):

$$\overline{m}_b(\overline{m}_b) = (4.23 \pm 0.05) \text{ GeV}$$

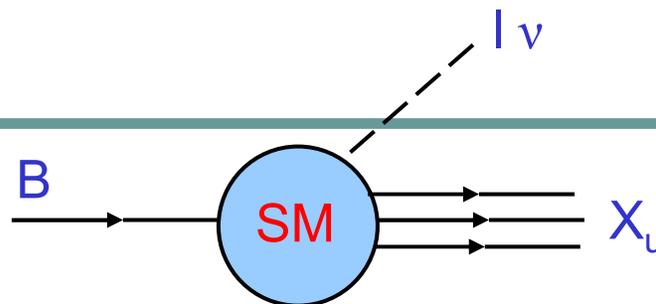
- Compare with extractions from  $\Upsilon$  spectrum (sum rules and lattice) :

$$\begin{aligned}\overline{m}_b(\overline{m}_b) &= (4.20 \pm 0.09) \text{ GeV} \\ &= (4.2 \pm 0.1 \pm 0.1) \text{ GeV}\end{aligned}$$

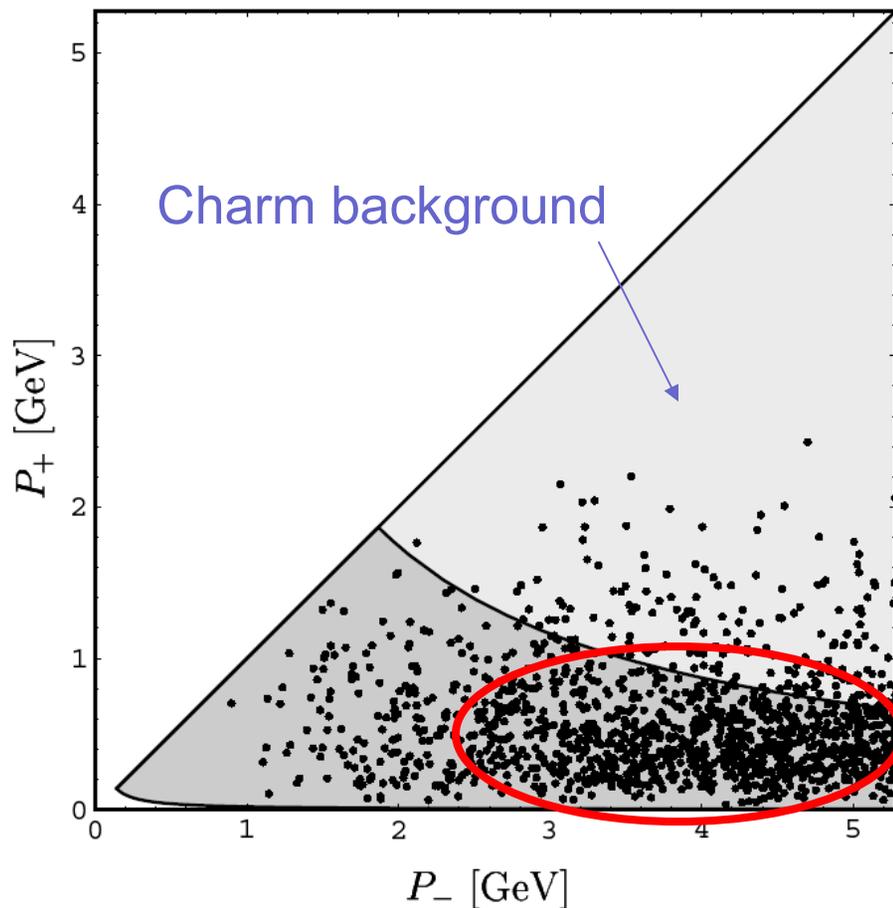
[Corcella, Hoang:  
hep-ph/0212297]  
[Lattice compilation:  
Sachrajda for PDG]

# $|V_{ub}|$ from $B \rightarrow X_u l \nu$ Decay

Factorization for inclusive semileptonic decay spectra



# Inclusive semileptonic decays



- Factorization theorem analogous to  $B \rightarrow X_s \gamma$
- Hadronic phase space is most transparent in the variables  $P_{\mp} = E_X \pm P_X$
- In practice,  $\Delta = P_+ - \bar{\Lambda}$  is always of order  $\Lambda_{\text{QCD}}$  for cuts eliminating the charm background

Shape-function region

# Strategy

- Exploit universality of shape function
- Extract shape function in  $B \rightarrow X_s \gamma$  (fit to photon spectrum), then predict arbitrary distributions in  $B \rightarrow X_u | \nu$  decay [Lange, MN, Paz: hep-ph/0504071]
- Functional form of fitting function constrained by model-independent moment relations
  - Knowledge of  $m_b$  and  $\mu_\pi^2$  helps!
- Variant: construct “shape-function independent relations” between spectra (equivalent) [Lange, MN, Paz: hep-ph/0508178]

# Results for various cuts

	$m_b$ [GeV]	4.50	4.55	4.60	4.65	4.70	Theory Error
$M_X \leq M_D$ <b>Eff = 84%</b>	$a$ Functional Form	9.5 1.4%	8.8 1.1%	8.2 0.8%	7.7 0.5%	7.3 0.4%	7%
$M_X \leq 1.7 \text{ GeV}$ <b>Eff = 75%</b>	$a$ Functional Form	12.5 2.9%	11.5 2.6%	10.5 2.2%	9.7 1.9%	8.9 1.6%	7%
$M_X \leq 1.7 \text{ GeV}$ $q^2 \geq 8 \text{ GeV}^2$ <b>35%</b>	$a$ Functional Form	10.3 2.0%	9.8 1.7%	9.3 1.5%	9.0 1.4%	8.7 1.4%	10%
$q^2 \geq (M_B - M_D)^2$ <b>Eff = 18%</b>	$a$ Functional Form	11.4 5.0%	11.1 4.4%	10.9 4.0%	10.8 3.6%	10.6 3.2%	15%
$P_+ \leq M_D^2/M_B$ <b>Eff = 65%</b>	$a$ Functional Form	16.7 5.3%	15.0 4.8%	13.6 4.4%	12.2 4.0%	11.1 3.6%	7%
$E_l \geq 2.2 \text{ GeV}$ <b>Eff = 11%</b>	$a$ Functional Form	22.6 16.2%	21.0 13.1%	19.7 11.0%	18.5 9.3%	17.4 7.9%	19%

Rate  $\Gamma \sim (m_b)^a$

[Lange, MN, Paz: hep-ph/0504071]

# Results for various cuts

- Different determinations now consistent:

	nominal $f_u$	$ V_{ub}  \times 10^3$
*CLEO [79] $E_e > 2.1 \text{ GeV}$	0.19	$4.02 \pm 0.47 \pm 0.35$
*BABAR [82] $E_e, s_h^{\max}$	0.19	$4.06 \pm 0.27 \pm 0.36$
*BABAR [81] $E_e > 2.0 \text{ GeV}$	0.26	$4.23 \pm 0.27 \pm 0.31$
*BELLE [80] $E_e > 1.9 \text{ GeV}$	0.34	$4.82 \pm 0.45 \pm 0.31$
*BABAR [86] $M_X/q^2$	0.34	$4.76 \pm 0.34 \pm 0.32$
*BELLE [87] $M_X/q^2$	0.34	$4.38 \pm 0.46 \pm 0.30$
BELLE [85] $M_X/q^2$	0.34	$4.68 \pm 0.37 \pm 0.32$
BELLE [85] $P_+ < 0.66 \text{ GeV}$	0.57	$4.14 \pm 0.35 \pm 0.29$
*BELLE [85] $M_X < 1.7 \text{ GeV}$	0.66	$4.08 \pm 0.27 \pm 0.25$
Average of *		$4.38 \pm 0.19 \pm 0.27$
$\chi^2 = 5.9/6, \text{ CL}=0.43$		

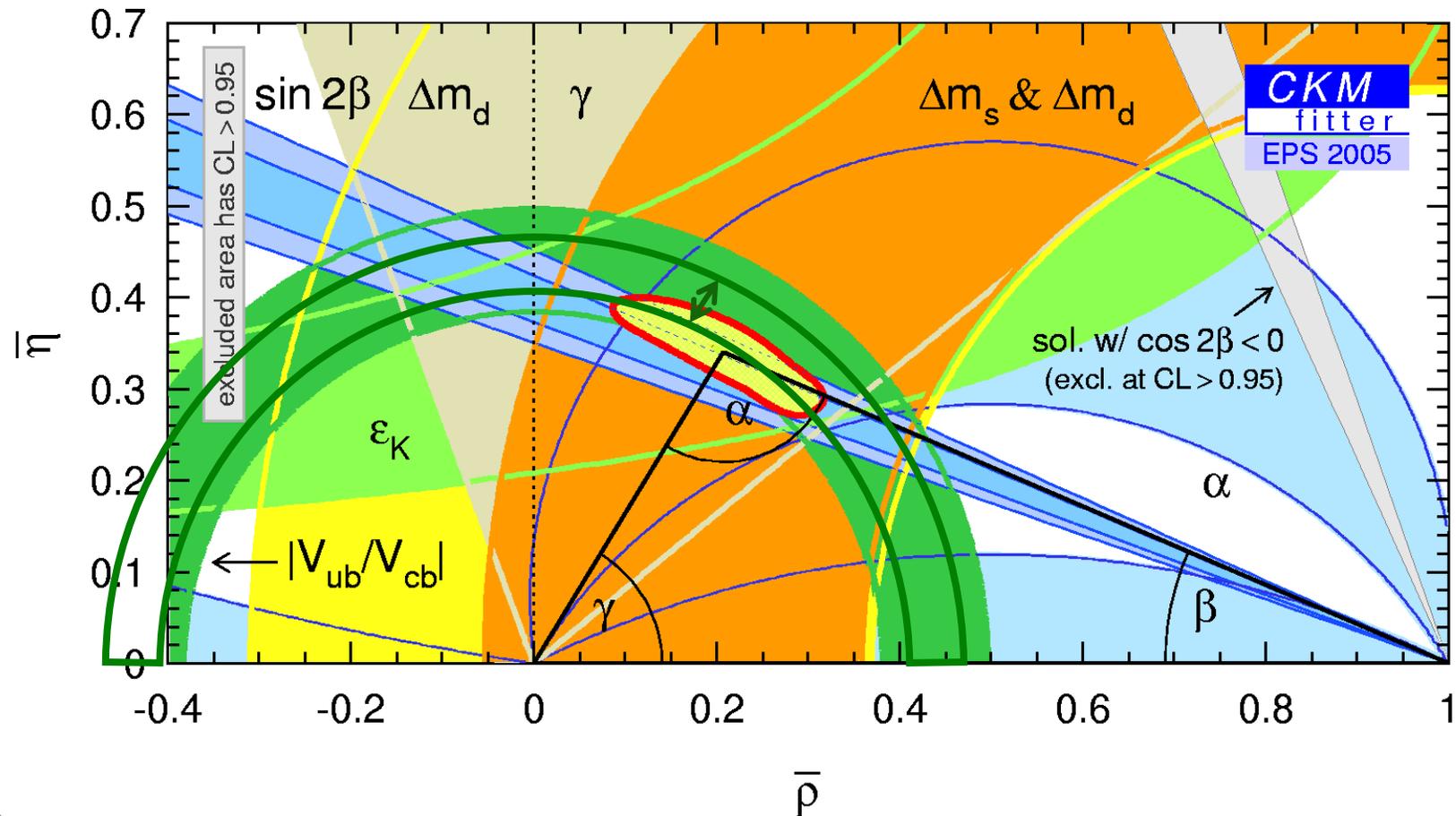
# Combined result

- Theory error on  $|V_{ub}|$  is 5-10% for several different cuts (10% now conservative – seemed unrealistic only a few years ago)
- Average of different extractions gives  $|V_{ub}|$  with a *total* error of 7%:

$$|V_{ub}| = (4.38 \pm 0.33) \cdot 10^{-3}$$

- Needed to match the precision of  $\sin 2\beta$

# Impact of precise $|V_{ub}|$



# Shape-function free relations

- Example for  $P_+ = E_X - P_X$  spectrum:

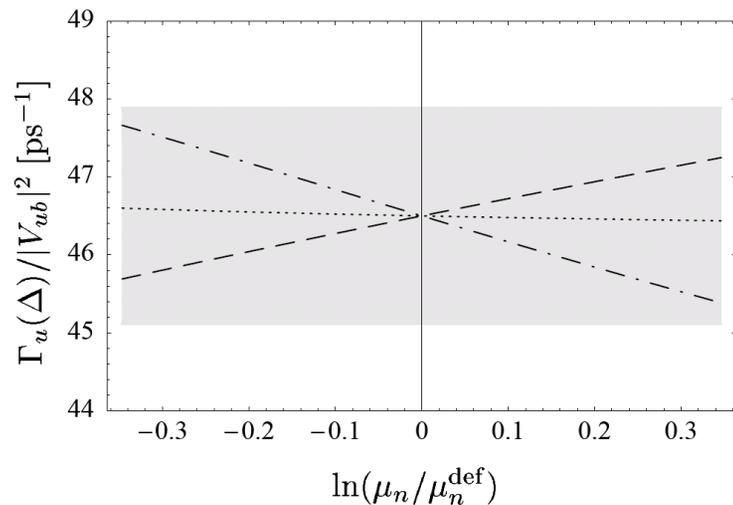
$$\Gamma_u(\Delta) = \underbrace{\int_0^\Delta dP_+ \frac{d\Gamma_u}{dP_+}}_{\text{exp. input}} = |V_{ub}|^2 \int_0^\Delta dP_+ \underbrace{W(\Delta, P_+)}_{\text{theory}} \underbrace{\frac{1}{\Gamma_s(E_*)} \frac{d\Gamma_s}{dP_+}}_{\text{exp. input}}$$

- Weight function is perturbatively calculable and *independent* of soft scale
- Small hadronic uncertainties enter at order  $1/m_b$  only

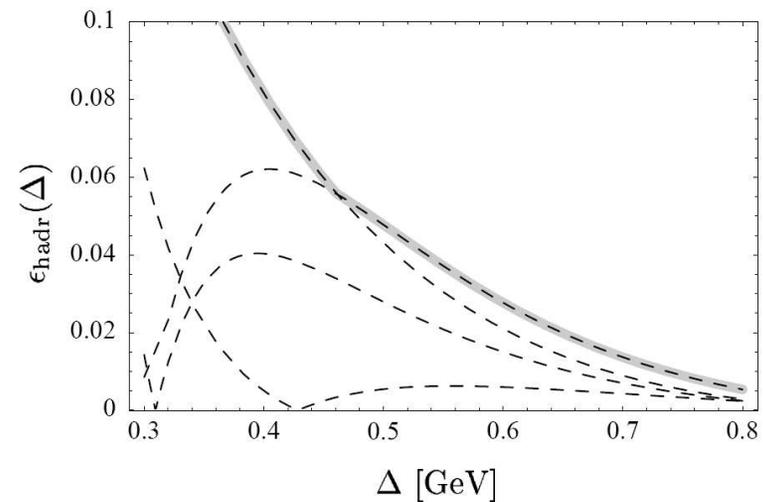
[Lange, MN, Paz: hep-ph/0508178]

# Shape-function free relations

- Perturbative stability:



- Hadronic uncertainty:



$$\begin{aligned} \Gamma_u(0.65 \text{ GeV}) &= (46.5 \pm 1.4 \text{ [pert]} \pm 1.8 \text{ [hadr]} \pm 1.8 \text{ [} m_b \text{]} \pm 0.8 \text{ [pars]} \pm 2.8 \text{ [norm]}) |V_{ub}|^2 \text{ ps}^{-1} \\ &= (46.5 \pm 4.1) |V_{ub}|^2 \text{ ps}^{-1}, \end{aligned}$$

Provides extraction of  $|V_{ub}|$  with 4.4% theoretical uncertainty!

# Conclusions



# Summary

- $B \rightarrow X_s \gamma$  decay remains one of the most sensitive probes of New Physics
- Strong motivation for NNLO calculation in Standard Model (many people involved!)
- Important to disentangle hard and soft effects using factorization and resummation

# Summary

- Many applications besides New Physics searches:
  - Most precise determination of b-quark mass
  - Most precise determination of  $|V_{ub}|$

$$\bar{m}_b(\bar{m}_b) = (4.23 \pm 0.05) \text{ GeV}$$

$$|V_{ub}| = (4.38 \pm 0.33) \cdot 10^{-3}$$

Substantial advances in a challenging field !